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Blow-up Rate Estimates and Blow-up Set for a System of Two Heat Equations with Coupled Nonlinear Neumann Boundary Conditions

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Abstract

This paper deals with the blow-up properties of positive solutions to a parabolic system of two heat equations, defined on a ball in R^n , associated with coupled Neumann boundary conditions of exponential type. The upper bounds of blow-up rate estimates are derived. Moreover, it is proved that the blow-up in this problem can only occur on the boundary.

Keywords: Heat equation; Neumann boundary conditions; Blow-up set; Blow-up rate estimate; Green function.

1. Introduction

In this paper, we consider the following parabolic system of two heat equations associated with Neumann boundary conditions:

$$\left. \begin{aligned} u_t &= \Delta u, & v_t &= \Delta v, & (x, t) &\in B_R \times (0, T), \\ \frac{\partial u}{\partial \eta} &= \lambda_1 e^{v^p}, & \frac{\partial v}{\partial \eta} &= \lambda_2 e^{u^q}, & (x, t) &\in \partial B_R \times (0, T), \\ u(x, 0) &= u_0(x), & v(x, 0) &= v_0(x), & x &\in B_R, \end{aligned} \right\} \quad (1)$$

where $p, q > 1; \lambda_1, \lambda_2 > 0; B_R$ is a ball in $R^n; \eta$ is the outward normal; u_0, v_0 are both smooth functions, radially symmetric, nonzero, nonnegative and satisfy the condition:

$$\begin{aligned} \Delta u_0, \Delta v_0 &\geq 0, & u_{0r}(|x|), v_{0r}(|x|) &\geq 0, & \text{for } x \in \bar{B}_R, \\ \text{and } \frac{\partial u_0}{\partial \eta} &= \lambda_1 e^{v_0^p}, & \frac{\partial v_0}{\partial \eta} &= \lambda_2 e^{u_0^q}, & x \in \partial B_R \end{aligned} \quad (2)$$

Since the last decades, many authors have studied the blow-up properties to solutions of parabolic problems, defined on bounded domains [see for instance 1, 2]. One of the studied problems is the system of two heat equations defined in a ball, associated with coupled Neumann boundary conditions:

$$\left. \begin{aligned} u_t &= \Delta u, & v_t &= \Delta v, & (x, t) &\in B_R \times (0, T), \\ \frac{\partial u}{\partial \eta} &= f(v), & \frac{\partial v}{\partial \eta} &= g(u), & (x, t) &\in \partial B_R \times (0, T), \\ u(x, 0) &= u_0(x), & v(x, 0) &= v_0(x), & x &\in B_R, \end{aligned} \right\} \quad (3)$$

This problem was previously studied [3-6]; in case of the nonlinear functions f and g take one of the two forms:

$$f(v) = v^p, \quad g(u) = u^q, \quad p, q > 1. \quad (4)$$

$$f(v) = e^{pv}, \quad g(u) = e^{qu}, \quad p, q > 0. \quad (5)$$

For both cases, it was shown that if the initial data (u_0, v_0) are nonzero and nonnegative, then the blow-up can only occur on the boundary.

In addition to that, with case 4, it was proved that the blow-up rate estimates take the form:

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$$c \leq \max_{x \in \Omega} u(x, t)(T - t)^{\frac{p+1}{2(pq-1)}} \leq C, \quad t \in (0, T),$$

$$c \leq \max_{x \in \Omega} v(x, t)(T - t)^{\frac{q+1}{2(pq-1)}} \leq C, \quad t \in (0, T)$$

where c and C are positive constants.

While, with case 5, it was proved that the blow-up rate estimates take the form:

$$C_1 \leq e^{qu(R,t)}(T - t)^{\frac{1}{2}} \leq C_2,$$

$$C_3 \leq e^{pv(R,t)}(T - t)^{1/2} \leq C_4$$

where C_1, C_2, C_3 and C_4 are positive constants.

In this paper, firstly, we show that the upper blow-up rate estimates for problem 1 are as follows

$$\max_{\bar{B}_R} u(x, t) \leq \log C_1 - \frac{\alpha}{2} \log(T - t), \quad 0 < t < T,$$

$$\max_{\bar{B}_R} v(x, t) \leq \log C_2 - \frac{\beta}{2} \log(T - t), \quad 0 < t < T,$$

where $\alpha = \frac{p+1}{pq-1}, \beta = \frac{q+1}{pq-1}$.

Secondly, we prove that the blow-up in problem 1 can only occur on the boundary.

2. Preliminaries

It is well known that with any smooth initial functions (u_0, v_0) , satisfying the compatibility condition 2, there exists a unique local classical solution to problem 1 [7]. On the other hand, it is easy to show that every nontrivial solution blows up simultaneously in finite time and that due to the known blow-up results of problem 3 with 4 and the comparison principle [2,3].

The next lemma, which was previously proved [2], states some properties of the classical solutions of problem1.

For simplicity, we denote $u(r, t) = u(x, t)$.

Lemma 2.1 Let (u, v) be a classical solution to problem 1. Then

1. u, v are positive, radial. Moreover, $u_r, v_r \geq 0$ in $[0, R] \times (0, T)$.
2. $u_t, v_t > 0$ in $\bar{B}_R \times (0, T)$.

3. Blow –up Upper Rate Estimates

The next Lemmas and theorem, proved in other articles [5,8], will be used in this section to derive the upper blow-up rate estimates for problem 1.

Lemma 3.1 [5]: Let A and B be positive and differentiable functions in $[0, T)$, such that they satisfy the two inequalities:

$$A'(t) \geq c \frac{B^p(t)}{\sqrt{T-t}}, \quad B'(t) \geq c \frac{A^q(t)}{\sqrt{T-t}}$$

for $t \in [0, T)$,

$$A(t) \rightarrow +\infty \quad \text{or} \quad B(t) \rightarrow +\infty \quad \text{as } t \rightarrow T^-,$$

where $p, q > 0, c > 0$ and $pq > 1$.

Then there exists $C > 0$ such that

$$A(t) \leq C(T - t)^{-\frac{\alpha}{2}}, \quad B(t) \leq C(T - t)^{-\frac{\beta}{2}}, \quad t \in [0, T),$$

where $\alpha = \frac{p+1}{pq-1}, \beta = \frac{q+1}{pq-1}$.

Lemma 3.2 [6]: Let $x \in \bar{B}_R$. If $0 \leq a < n - 1$. Then there exists $C > 0$ such that

$$\int_{S_R} \frac{ds_y}{|x-y|^a} \leq C.$$

Theorem 3.3 (Jump relation, [8]) Let $\Gamma(x, t)$ be the fundamental solution of heat equation, namely

$$\Gamma(x, t) = \frac{1}{(4\pi t)^{(n/2)}} \exp\left[-\frac{|x|^2}{4t}\right]. \tag{6}$$

Let φ be a continuous function on $S_R \times [0, T]$. Then for any $x \in B_R, x^0 \in S_R, 0 < t_1 < t_2 \leq T$, for some $T > 0$, the function

$$U(x, t) = \int_{t_1}^{t_2} \int_{S_R} \Gamma(x - y, t - z) \varphi(y, z) ds_y d\tau$$

satisfies the jump relation

$$\frac{\partial}{\partial \eta} U(x, t) \rightarrow -\frac{1}{2} \varphi(x^0, t) + \frac{\partial}{\partial \eta} U(x^0, t)$$

$$asx \rightarrow x^0.$$

Theorem 3.4 Let (u, v) be a blow-up solution to problem1, and $T > 0$ is the blow-up time. Then there exist two positive constants C_1, C_2 such that

$$\max_{\overline{B_R}} u(x, t) \leq \log C_1 - \frac{\alpha}{2} \log(T - t), \quad 0 < t < T,$$

$$\max_{\overline{B_R}} v(x, t) \leq \log C_2 - \frac{\beta}{2} \log(T - t), \quad 0 < t < T.$$

Proof: In order to prove this theorem, we follow the technique used in a previous work [5].

Define the functions M and M_b as follows:

$$M(t) = \max_{\overline{B_R}} u(x, t), \quad and M_b(t) = \max_{S_R} u(x, t).$$

Similarly,

$$N(t) = \max_{\overline{B_R}} v(x, t), \quad and N_b(t) = \max_{S_R} v(x, t).$$

Depending on Lemma 2.1, both of M, M_b are monotone increasing functions.

Since u is a solution of heat equation, it cannot attain interior maximum without being constant. Therefore,

$$M(t) = M_b(t). \quad Similarly \quad N(t) = N_b(t).$$

Moreover, since u, v blow up simultaneously, we have

$$M(t) \rightarrow +\infty, \quad N(t) \rightarrow +\infty \quad as t \rightarrow T^-(7)$$

According to the second Green's identity [5,7, 9], with considering the Green function:

$$G(x, y; z_1, t) = \Gamma(x - y, t - z_1), \text{ for } 0 < z_1 < t < T \text{ and } x \in B_R,$$

where Γ is defined in 6, the integral equation to problem 1, with respect to u , takes the form:

$$u(x, t) = \int_{B_R} \Gamma(x - y, t - z_1) u(y, z_1) dy + \int_{z_1}^t \int_{S_R} \lambda_1 e^{v^p(y, \tau)} \Gamma(x - y, t - \tau) ds_y d\tau$$

$$- \int_{z_1}^t \int_{S_R} u(y, \tau) \frac{\partial \Gamma}{\partial \eta_y}(x - y, t - \tau) ds_y d\tau,$$

By applying Theorem 3.3 on the third term in the right-hand side of the last equation and with letting $x \rightarrow S_R$, we obtain

$$\frac{1}{2} u(x, t) = \int_{B_R} \Gamma(x - y, t - z_1) u(y, z_1) dy + \int_{z_1}^t \int_{S_R} \lambda_1 e^{v^p(y, \tau)} \Gamma(x - y, t - \tau) ds_y d\tau$$

$$- \int_{z_1}^t \int_{S_R} u(y, \tau) \frac{\partial \Gamma}{\partial \eta_y}(x - y, t - \tau) ds_y d\tau,$$

for $x \in S_R, 0 < z_1 < t < T$.

Depending on Lemma 2.1, u, v are both radial and positive functions.

Therefore,

$$\int_{B_R} \Gamma(x - y, t - z_1) u(y, z_1) dy > 0,$$

$$\int_{z_1}^t \int_{S_R} \lambda_1 e^{v^p(y, \tau)} \Gamma(x - y, t - \tau) ds_y d\tau = \int_{z_1}^t \lambda_1 e^{v^p(R, \tau)} [\int_{S_R} \Gamma(x - y, t - \tau) ds_y] d\tau.$$

This leads to

$$\frac{1}{2} M(t) \geq \int_{z_1}^t \lambda_1 e^{N^p(\tau)} [\int_{S_R} \Gamma(x - y, t - \tau) ds_y] d\tau - \int_{z_1}^t M(\tau) \left[\int_{S_R} \left| \frac{\partial \Gamma}{\partial \eta_y}(x - y, t - \tau) \right| ds_y \right] d\tau,$$

$$x \in S_R, 0 < z_1 < t < T.$$

It is known that (see [8]) there exists $C_0 > 0$, such that Γ satisfies

$$\left| \frac{\partial \Gamma}{\partial \eta_y}(x - y, t - \tau) \right| \leq \frac{C_0}{(t - \tau)^\mu} \cdot \frac{1}{|x - y|^{(n+1-2\mu-\sigma)'}}$$

$$x, y \in S_R, \sigma \in (0, 1).$$

Choose $1 - \frac{\sigma}{2} < \mu < 1$, from Lemma 3.2, there exists $C^* > 0$ such that

$$\int_{S_R} \frac{ds_y}{|x - y|^{(n+1-2\mu-\sigma)}} < C^*.$$

Moreover, for $0 < t_1 < t_2$ and t_1 is closed to t_2 , there exists $c > 0$, such that

$$\int_{S_R} \Gamma(x - y, t_2 - t_1) ds_y \geq \frac{c}{\sqrt{t_2 - t_1}}$$

Thus

$$\frac{1}{2}M(t) \geq c \int_{z_1}^t \frac{\lambda_1 e^{N^p(\tau)}}{\sqrt{t - \tau}} d\tau - C \int_{z_1}^t \frac{M(\tau)}{|t - \tau|^\mu} d\tau.$$

Since for $0 < z_1 < t_0 < t < T$, it follows that $M(t_0) \leq M(t)$, thus the last equation becomes

$$\frac{1}{2}M(t) \geq c \int_{z_1}^t \frac{\lambda_1 e^{N^p(\tau)}}{\sqrt{T - \tau}} d\tau - C_1^* M(t) |T - z_1|^{1-\mu}.$$

Similarly, for $0 < z_2 < t < T$, we have

$$\frac{1}{2}N(t) \geq c \int_{z_2}^t \frac{\lambda_2 e^{M^q(\tau)}}{\sqrt{T - \tau}} d\tau - C_2^* N(t) |T - z_2|^{1-\mu}.$$

Taking z_1, z_2 so that

$$C_1^* |T - z_1|^{1-\mu} \leq 1/2, \quad C_2^* |T - z_2|^{1-\mu} \leq 1/2,$$

it follows

$$M(t) \geq c \int_{z_1}^t \frac{\lambda_1 e^{N^p(\tau)}}{\sqrt{T - \tau}} d\tau, \quad N(t) \geq c \int_{z_2}^t \frac{\lambda_2 e^{M^q(\tau)}}{\sqrt{T - \tau}} d\tau. \tag{8}$$

Since M, N are both increasing functions and by 7, we can find $T_1 \in (0, T)$, such that

$$M(t) \geq q^{\frac{1}{q-1}}, \quad N(t) \geq p^{\frac{1}{p-1}}, \quad \text{for } T_1 \leq t < T.$$

Thus

$$e^{M^q(t)} \geq e^{qM(t)}, \quad e^{N^p(t)} \geq e^{pN(t)}, \quad T^* \leq t < T.$$

Therefore, if we choose z_1, z_2 in (T^*, T) , then 8 becomes

$$e^{M(t)} \geq c \int_{z_1}^t \frac{\lambda_1 e^{pN(\tau)}}{\sqrt{T - \tau}} d\tau \equiv I_1(t),$$

$$e^{N(t)} \geq c \int_{z_2}^t \frac{\lambda_2 e^{qM(\tau)}}{\sqrt{T - \tau}} d\tau \equiv I_2(t).$$

Clearly,

$$I_1'(t) = c \frac{\lambda_1 e^{pN(t)}}{\sqrt{T - t}} \geq \frac{c\lambda_1 I_2^p}{\sqrt{T - t}}, \quad I_2'(t) = c \frac{\lambda_2 e^{qM(t)}}{\sqrt{T - t}} \geq \frac{c\lambda_2 I_1^q}{\sqrt{T - t}}$$

By Lemma 3.1, it follows that

$$I_1(t) \leq \frac{C\lambda_1}{(T-t)^{\frac{\alpha}{2}}}, \quad I_2(t) \leq \frac{C\lambda_2}{(T-t)^{\frac{\beta}{2}}}, \tag{9}$$

$t \in (\max\{z_1, z_2\}, T)$.

On the other hand, with assuming that t is close to T , we have

$$I_1(t) \geq c \int_{t^*}^t \frac{\lambda_1 e^{pN(\tau)}}{\sqrt{T - \tau}} d\tau \geq c\lambda_1 e^{pN(t^*)} \int_{2t-T}^t \frac{1}{\sqrt{T - \tau}} d\tau = 2c\lambda_1(\sqrt{2} - 1)\sqrt{T - t} e^{pN(t^*)}$$

where $t^* = 2t - T$

Combining the last inequality with 9 yields

$$e^{N(t^*)} \leq \frac{C}{2c(\sqrt{2} - 1)(T - t)^{\frac{p+1}{2p(pq-1)} + \frac{1}{2p}}} = \frac{\frac{q+1}{2^{2(pq-1)}}C}{2c(\sqrt{2} - 1)(T - t^*)^{\frac{q+1}{2(pq-1)}}}.$$

It follows that, there exists a constant $c_1 > 0$ such that

$$e^{N(t^*)} (T - t^*)^{\frac{q+1}{2(pq-1)}} \leq c_1.$$

Similarly, we can find $c_2 > 0$ such that

$$e^{M(t^*)} (T - t^*)^{\frac{p+1}{2(pq-1)}} \leq c_2.$$

This leads to, there exist $C_1, C_2 > 0$ such that

$$\max_{\overline{B}_R} u(x, t) \leq \log C_1 - \frac{\alpha}{2} \log(T - t), \quad , \tag{10}$$

$$\max_{\overline{B}_R} v(x, t) \leq \log C_2 - \frac{\beta}{2} \log(T - t) \quad . \tag{11}$$

for $0 < t < T$

4. Blow-up Set

In this section, we study the blow-up set for problem 1, showing that the blow-up can only occur on the boundary. To prove this result, we recall the following lemma proved in a previous article [6].

Lemma 4.1. Let w be a continuous function on the domain $\bar{B}_R \times [0, T)$ and satisfies

$$\left. \begin{aligned} w_t &= \Delta w, & (x, t) \in B_R \times (0, T), \\ w(x, t) &\leq \frac{C}{(T-t)^m}, & (x, t) \in S_R \times (0, T), \quad m > 0 \end{aligned} \right\}$$

Then for any $0 < a < R$,

$$\sup\{w(x, t): 0 \leq |x| \leq a, 0 \leq t < T\} < \infty.$$

Proof: Set

$$h(x) = (R^2 - r^2)^2, r = |x|,$$

$$z(x, t) = \frac{C_1}{[h(x) + C_2(T-t)]^m}.$$

We can show that:

$$\Delta h - \frac{(m+1)|\nabla h|^2}{h} = 8r^2 - 4n(R^2 - r^2) - (m+1)16r^2 \geq -4nR^2 - 16R^2(m+1),$$

and

$$z_t - \Delta z = \frac{C_1 m}{[h(x) + C_2(T-t)]^{m+1}} (C_2 + \Delta h - \frac{(m+1)|\nabla h|^2}{h+C_2(T-t)}) \geq \frac{C_1 m}{[h(x) + C_2(T-t)]^{m+1}} (C_2 - 4nR^2 - 16R^2(m+1)).$$

Let $C_2 = 4nR^2 + 16R^2(m+1) + 1$, and take C_1 to be large such that

$$z(x, 0) \geq w(x, 0), \quad x \in B_R.$$

Let $C_1 \geq C(C_2)^m$, which implies that

$$z(x, t) \geq w(x, t) \quad \text{on } S_R \times [0, T).$$

Then from the maximum principle [10], it follows that

$$z(x, t) \geq w(x, t), \quad (x, t) \in \bar{B}_R \times (0, T)$$

and hence

$$\sup\{w(x, t): 0 \leq |x| \leq a, 0 \leq t < T\} \leq C_1(R^2 - a^2)^{-2m} < \infty,$$

for $0 \leq a < R$.

Theorem 4.2 Let (u, v) be a blow-up solution to problem1, and $T > 0$ is the blow-up time. Then (u, v) can only blow-up on the boundary.

Proof: By using equations 10 and 11, we obtain

$$u(R, t) \leq \frac{C_1}{(T-t)^{\frac{\alpha}{2}}}, \quad v(R, t) \leq \frac{C_2}{(T-t)^{\frac{\beta}{2}}}$$

for $t \in (0, T)$.

From Lemma 4.1, it follows that

$$\sup\{u(x, t): (x, t) \in B_a \times [0, T)\} \leq C_1(R^2 - a^2)^{-\alpha} < \infty,$$

$$\sup\{v(x, t): (x, t) \in B_a \times [0, T)\} \leq C_1(R^2 - a^2)^{-\beta} < \infty,$$

for $a < R$. Therefore, if $x \in B_R$, it cannot be a blow-up point.

5. Conclusions

This paper is concerned with the blow-up properties of positive solutions to a system of two heat equations defined on a ball in R^n associated with coupled Neumann boundary conditions of exponential type. Firstly, the upper bounds of blow-up rate estimates are derived. Secondly, the blow-up set is considered. The main conclusion is that the blow-up in this problem only occurs on the boundary.

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