Dark Energy Surveyed Year 1 results: calibration of cluster mis-centring in the redMaPPer catalogues

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Dark Energy Surveyed Year 1 results: calibration of cluster mis-centring in the redMaPPer catalogues


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ABSTRACT

The centre determination of a galaxy cluster from an optical cluster finding algorithm can be offset from theoretical prescriptions or N-body definitions of its host halo centre. These offsets impact the recovered cluster statistics, affecting both richness measurements and the weak lensing shear profile around the clusters. This paper models the centring performance of the redMaPPer cluster finding algorithm using archival X-ray observations of redMaPPer-selected clusters. Assuming the X-ray emission peaks as the fiducial halo centres, and through analysing their offsets to the redMaPPer centres, we find that ∼75 ± 8 per cent of the redMaPPer clusters are well centred and the mis-centred offset follows a Gamma distribution in normalized, projected distance. These mis-centring offsets cause a systematic underestimation of cluster richness relative to the well-centred clusters, for which we propose a descriptive model. Our results enable the DES Y1 cluster cosmology analysis by characterizing the necessary corrections to both the weak lensing and richness abundance functions of the DES Y1 redMaPPer cluster catalogue.

Key words: galaxies: clusters: general – dark energy – X-rays: galaxies: clusters.
1 INTRODUCTION

The abundance of galaxy clusters is a sensitive probe of cosmological models (see reviews and the referenced literature in Allen, Evrard & Mantz 2011; Weinberg et al. 2013). Cluster cosmology studies from the latest optical imaging surveys such as the Dark Energy Survey (the DES Collaboration, in preparation) will deliver significant improvement in precision over previous studies based on the Sloan Digital Sky Survey (SDSS, Dietrich et al. 2014), correlated structures and their projection effect (Erickson, Cunha & Evrard 2011; Costanzi et al. 2019), mass profile modelling uncertainties (McClintock et al. 2019), and the contamination of cluster member galaxies in the lensing measurements (Varga et al. 2018).

One such important systematic effect is the mis-identification of cluster centres (Johnston et al. 2007a,b; Melchior et al. 2017; Simet et al. 2017; McClintock et al. 2019; the DES Collaboration, in preparation). Cluster observables, e.g. gravitational shear profiles, must be compared to models in order to derive constraints on parameters, and the models are based on some definition of cluster centre, described theoretically or on the basis of the matter density field in N-body simulations. Cluster lensing studies based on data from DES (McClintock et al. 2019) require accurate knowledge of cluster mis-centring fraction and offset distribution in order to forward model masses using analytic halo profiles.

Optical cluster finders often attempt to identify a central galaxy as the centre (Koester et al. 2007; Hao et al. 2010; Rykoff et al. 2014; Oguri et al. 2018). These central galaxies are typically quenched of star formation activities and may appear to be the brightest galaxy in a cluster (see studies in Skibba et al. 2011; Lauer et al. 2014; Hoshino et al. 2015). The identification of cluster central galaxies may seem straightforward given their dominant appearances, but mis-identifications or offsets relative to any other theoretical definition of centre is inevitable. Because massive haloes experience growth through mergers, cluster central galaxies can be displaced from the local gravitational potential minimum (e.g. Martel, Robichaud & Burlat 2014). A related effect is when a second galaxy within the halo is chosen as the centre by the cluster finding algorithm. For a colour-based scheme focused on the reddest galaxies, this may happen if the central galaxy of the host halo has experienced recent star formation (e.g. McDonald et al. 2012; Donahue et al. 2015) or if a merging event brings in two nearly identical central galaxies of the progenitor haloes, as in the case of the Coma cluster (e.g. Vikhlinin et al. 2001). Another cause of mis-centring is when galaxies lying outside the primary host halo, but aligned in projection, are chosen as the central galaxy by the cluster finding algorithm.

The centring performance of optical cluster finding algorithms has been characterized with various methods. Cluster hot gas is an excellent tracer of the cluster potential as the dominant baryonic mass component. Cluster X-ray or thermal Sunyaev–Zel’dovich (tSZ) observation centres, identified as the centroids or the peaks of the surface brightness, are often used to calibrate the optically selected centres (see examples of X-ray studies in Lin & Mohr 2004; Stott et al. 2012; Mahdavi et al. 2013; Lauer et al. 2014; Rozo & Rykoff 2014; Rykoff et al. 2016; Zhang et al. 2016 and examples of tSZ studies in Song et al. 2012; Saro et al. 2015). Other than calibration to centres identified in multiwavelength observations, cluster centring has been characterized through comparing cluster radial profiles (lensing or galaxy number count) to those of a cluster sample with well-known centres from X-ray or optical data (Hikage et al. 2018; Luo et al. 2017) and through examining the velocity and separation distribution of cluster satellite galaxies (Skibba et al. 2011).

In this paper, we characterize the centring performance of the redMaPPer cluster finding algorithm – a method for identifying galaxy clusters from optical imaging data. The redMaPPer algorithm excels in producing a complete and efficient cluster sample with an accurate richness mass proxy and precise redshift estimations, as characterized with multiwavelength and spectroscopic data (Rykoff et al. 2012; Rozo & Rykoff 2014; Rozo et al. 2015a,b; Saro et al. 2015; Murata et al. 2018). Cluster catalogues constructed from SDSS (Rykoff et al. 2016) and DES data (Rykoff et al. 2016) are used to derive cosmological constraints in Costanzi et al. (2018) and the DES Collaboration (in preparation). In terms of the cluster centre identification, redMaPPer is not exempt from occasional mis-identification of the central galaxy.

The centring distribution of the redMaPPer algorithm has been studied using almost all of the aforementioned methods (Rozo & Rykoff 2014; Sato et al. 2015; Rykoff et al. 2016; Hikage et al. 2017). In this paper, we model the cluster centring distribution with 211 high-signal-to-noise X-ray cluster detections associated with the redMaPPer SDSS DR8 and DES Year 1 samples from the Chandra public archives, and the model constraints are then validated with X-ray cluster detections from the XMM public archives. We focus on the modelling aspects of redMaPPer centring performance in this paper, while the X-ray data processing procedures are presented in Hollowood et al. (2018) and Giles et al. (in preparation). The data set and methods we employ allow us to analyse the redMaPPer centring performance with the highest precision to date. We also develop a model to characterize how mis-centring affects redMaPPer richness estimation and discuss the impact of cluster mis-centring on cluster cosmology analyses. It is the first time that this effect has been quantified.

This paper is a companion paper to the DES and SDSS cluster weak lensing and cosmology studies presented in McClintock et al. (2019), Costanzi et al. (2018), and the DES Collaboration (in preparation). It uses similar data products to Farahi et al. (2019).

Throughout this paper, we assume a Flat Λ cold dark matter cosmology with $h = 0.7$ and $\Omega_m = 0.3$.

2 DATA

2.1 The redMaPPer catalogues

The redMapper algorithm examines galaxy colour, spatial overdensity, and galaxy luminosity distribution to identify possible galaxy clusters. Cluster centres are placed on a central galaxy candidate according to the colour, luminosity, and galaxy overdensity computed around the galaxy. Up to five central galaxy candidates are recorded for each cluster with probabilities assigned to them. The cluster centre is chosen to be the most probable one. The redMaPPer algorithm also estimates a richness as a mass proxy, $\lambda$, which is a probabilistic count of red sequence galaxies within an aperture centred on the central galaxy candidate. Detailed presentation of this algorithm can be found in Rykoff et al. (2014, 2016).

The redMaPPer samples studied in this paper have been derived from both SDSS (Rykoff et al. 2014, 2016) and DES (Rykoff et al. 2016; McClintock et al. 2019) data. We use the redMaPPer SDSS 6.3.1 sample (Costanzi et al. 2018) derived from SDSS DR8 (Aihaara...
et al. 2011) photometric data and the redMaPPer DES-Y1-6.4.17 volume-limited catalogue which is based on the DES Y1 gold catalogue (Drlica-Wagner et al. 2018), with cluster richnesses ≥ 20. For the SDSS redMaPPer sample, we consider clusters in the redshift range of 0.1–0.35 which are nearly volume-limited. For the DES redMaPPer sample, we select clusters between redshifts 0.2 and 0.7.

2.2 Comparison of centres in redMaPPer catalogues

In this paper, we treat the DES and SDSS redMaPPer catalogues as two independent samples and characterize their centring performances separately. We examine the offset distribution between redMaPPer centres of the overlapping clusters in the DES and SDSS samples, to estimate an upper limit of the well-centred cluster fraction. We match between the DES and SDSS redMaPPer samples to identify the overlapping clusters in the redshift range of 0.2–0.35. For a pair of DES and SDSS clusters to be considered a match, their redshift difference, |Δz|, must be less than 0.05 to account for the scatter in photometric redshifts and possible blending effects. The redMaPPer centres must be within 1 Rλ (Rλ = (λ/100)^0.2 h^−1 Mpc). The radius aperture is chosen because redMaPPer does not consider the clusters to be the same if their centring offset is larger than Rλ. The richness estimations derived from SDSS and DES data have an average relation of λDES = (0.88 ± 0.03) × λSDSS + (3.28 ± 1.20). We look for matches to λ > 20 DES clusters in the SDSS redMaPPer sample by lowering the SDSS redMaPPer λ threshold to 5 to account for the λ difference and scatter between DES and SDSS.

There are 150 DES clusters with λ > 20 and 38 786 SDSS clusters with λ > 5 in the overlap region of both catalogues after applying redshift, position, and mask cuts. Of these 150 DES clusters, 148 have SDSS matches given the criteria listed above. 15 of these 148 clusters have at least two matches and the most likely SDSS match was chosen by inspection of redshift, position, and richness. As the purpose of this matching process is to verify the centring consistency in SDSS/DES, we further remove three clusters in the total sample (148) because of our poor confidence in the match: the SDSS matching candidates have large richness differences with their respective DES clusters. This further reduces our matching sample size to 145.

Fig. 1 shows the scaled offset distribution between SDSS and DES centres for the matched clusters. For 77 per cent of the matched clusters, their SDSS and DES centres are within 0.05 Rλ (corresponding to ∼50 kpc at λ = 20, close to the size of a typical central galaxy, Zibetti et al. 2005; Stott et al. 2011). We consider these clusters as consistently centred between the two catalogues. The remaining 23 per cent of the clusters comprise a long tail in the SDSS and DES offset distribution up to 1 Rλ. The inconsistency indicates that for at least one of the samples, the mis-centred fraction is ≥0.23/2 = 0.115. As we do not have sufficient information (i.e. enough X-ray observations) in the DES/SDSS samples to analyse which has a greater rate of well-centring, we decide to independently analyse the SDSS and DES redMaPPer samples in this paper.

1After further investigation, the other two DES clusters appear to have SDSS matches with the same central galaxy selections, but the redshift differences between the DES and SDSS match are ∼0.06, and thus did not pass our strict redshift difference cuts.

![Figure 1](https://academic.oup.com/mnras/article-abstract/487/2/2578/5490401/figure1)
centred on the redMaPPer centre, the pipeline first determines the X-ray centroid, re-centres, and then iteratively finds X-ray centroids until convergence is reached within 15 kpc. A cluster is considered to be detected if the signal-to-noise ratio within a final 500 kpc aperture centred on the converged centroid is greater than 5. For detected clusters, MATCHa analysis proceeds with attempts to measure \( L_X \), \( T_x \), and centroid within a set of apertures including 500 kpc, \( r_{2000}, f_{500} \), and core-cropped \( r_{300} \). Visual checks of non-detected redMaPPer clusters are employed to examine if any redMaPPer clusters with \textit{Chandra} observations were omitted in the process. We find one SDSS redMaPPer cluster with a large offset of 1.82 Mpc (1.40 \( R_x \)) for this cluster) between the X-ray centroid and the redMaPPer centre, possibly overlooked because of the initial 500 kpc X-ray centroid searching criteria. Since this omission makes up less than 1 per cent of the total SDSS \textit{Chandra} sample, we do not consider it in further analyses. No similar cases were found in the DES Y1 sample.

For the clusters with X-ray detections, MATCHa additionally determines the position of the X-ray peak starting from the reduced, exposure-corrected, and point source subtracted images. Images are smoothed with a Gaussian with \( \sigma = 50 \) kpc width, and the peak is defined to be the brightest pixel in this smoothed image. All peaks are then visually examined. In a small number of cases relic point source emission or the removal of a point source near the cluster peak are found to bias the peak determination. The peak position is adjusted after accounting for the point source emission. In addition, two failure modes are flagged and removed from the sample. First, for the centring analysis we remove clusters falling on or near a chip edge in the X-ray observation such that the position of the X-ray peak could not be reliably determined. Second, in a few cases the identified X-ray cluster is clearly not the redMaPPer cluster (e.g. a bright foreground or background cluster in the same observation), and these clusters are likewise removed (see Hollowood et al. 2018 for further detail). Moreover, there are some special redMaPPer mis-centring cases, denoted as mis-percolations in Hollowood et al. (2018), because these cases are related to a ‘percolation’ procedure of redMaPPer. In these cases there is a spatially close pair of clusters with similar redshifts, and the one with a less luminous X-ray detection is assigned with a greater richness. Hollowood et al. (2018) manually associates the richer redMaPPer candidate with the more luminous X-ray detection and removes the less rich system from the X-ray samples.

Finally, to improve the accuracy of X-ray peak location, among all the clusters identified in Hollowood et al. (2018), we further impose a signal-to-noise cut removing clusters with a signal-to-noise ratio less than 6.5 within a 500 kpc aperture. In the end, 144 redMaPPer SDSS clusters are identified with X-ray peak centres in the \textit{Chandra} archival data.

The compilation of the DES redMaPPer \textit{Chandra} sample follows a similar process to the SDSS redMaPPer sample, with the exception that the X-ray peaks of the DES sample are initially identified around the redMaPPer centres within 500 kpc due to pipeline re-factoring. The peak identifications are visually examined and adjusted if needed. Overall, 67 DES redMaPPer clusters are identified with \textit{Chandra} observations of signal-to-noise ratio higher than 6.5.

In Fig. 2, we show the redshift and richness distributions of the SDSS and DES redMaPPer \textit{Chandra} samples. Fig. 3 shows the scaled offset distributions between the X-ray peaks and the redMaPPer centres. The tail of the offset distribution indicates a population of mis-centred clusters. Examples of the mis-centred clusters can be found in Hollowood et al. (2018).

3 THE X-RAY REDMAPPER OFFSET

3.1 Model

In this paper, we use the X-ray peaks as the cluster fiducial centre, and model the offsets between X-ray peaks and redMaPPer centres to characterize the redMaPPer centring distribution. When redMaPPer misidentifies the galaxy at the cluster centre, an offset between X-ray peaks and the redMaPPer centres is expected. On the other hand, when the redMaPPer centres are correct, the X-ray peaks may still be offset from them because of the different dynamics and relaxation time-scales of gas and galaxies (e.g. see studies about cluster state and X-ray galaxy offsets in Mantz et al. 2015; Kim, Peter & Wittman 2017; Roberts, Parker & Hlavacek-Larrondo 2018), as well as observational uncertainties in identifying the X-ray peaks, but these offsets tend to be small, i.e. less than tens of kpc. Therefore, we expect the well and mis-centred redMaPPer clusters to have different offset distributions to the X-ray peaks, and we model the redMaPPer and X-ray offset as a mixture of well-centred and mis-centred components, written as:

\[
P(x|\rho, \sigma, \tau) = P_{\text{cent}}(x|\sigma) + (1 - P_{\text{cent}}(x|\sigma)) \times P_{\text{miscent}}(x|\tau),
\]

\[
P_{\text{cent}}(x|\sigma) = \frac{1}{\sigma} \exp\left(-\frac{x}{\sigma}\right),
\]

\[
P_{\text{miscent}}(x|\tau) = \frac{1}{\tau^2} \exp\left(-\frac{x}{\tau}\right).
\]

In the above model, the peaked exponential distribution characterized by the parameter \( \sigma \), \( P_{\text{cent}}(x|\sigma) \), describes the X-ray offset distribution of the correctly centred clusters. The second component, \( P_{\text{miscent}}(x|\tau) \), a Gamma distribution of shape parameter 2 and characterized by a scale parameter \( \tau \), describes the offset of the mis-centred clusters. This distribution has a heavy tail and the choice of this distribution is inspired by the \textit{Chandra} SDSS sample having an extended offset distribution. The fraction of well-centred clusters are modelled by the \( \rho \) parameter. In total, the models consists of three parameters \( \rho, \sigma \) and \( \tau \). The redMaPPer and X-ray centre offset, \( x \), is computed as

\[
x = r_{\text{offset}}/R_x.
\]

\[
R_x = (\lambda/100)^{0.2} \ h^{-1} \ \text{Mpc},
\]

The photometric and astrometric uncertainties of galaxy positions are negligible in DES compared to X-ray (Drlica-Wagner et al. 2018).
The \( R_\lambda = (\lambda/100)^{0.2} \, h^{-1} \, \text{Mpc} \) scaled offset distribution between the redMaPPer centres and the X-ray emission peaks for the redMaPPer SDSS and DES samples from the Chandra archival observations, with the inset zooming on the mis-centred component, starting at \( r_{\text{offset}}/R_\lambda = 0.05 \). The distribution can be fitted with two components – a concentrated component that represents the well-centred redMaPPer clusters, and an extended component that represents the mis-centred redMaPPer clusters. The best-fitting SDSS offset model is shown as the solid lines (black: well-centred model, red: mis-centred model), with the shaded regions representing the uncertainties.

where \( P(\rho, \sigma, \tau) \) is the prior distribution of \( \rho, \sigma, \) and \( \tau \) listed in Table 1, which is chosen to be flat and independent parameters. We sample the posterior distribution of \( \rho, \sigma, \) and \( \tau \) using a Markov Chain Monte Carlo (MCMC) method.

### 3.2 Model constraints

The aforementioned X-ray redMaPPer offset model is constrained separately for the Chandra SDSS and DES redMaPPer samples. Table 1 lists and Fig. 4 shows the posterior constraints of the model parameters including the correctly centred fraction \( \rho \), and the mis-centred characteristic offset \( \tau \), as well as the characteristic redMaPPer X-ray offset, \( \sigma \). The SDSS sample yields higher
precision because of the larger sample sizes. The fraction of well- centred clusters $\rho$ and the mis-centring offset $\tau$ for the mis-centred clusters are mildly different from the DES redMaPPer sample which displays a hint of having a higher fraction of well-centred clusters at a 1.5σ significance level. For the well-centred clusters, the characteristic redMaPPer X-ray offset, $\sigma$, of the DES sample is larger than the respective parameter of the SDSS sample, reflecting the limited angular resolution of X-ray peak identification and the higher redshift range of the DES sample (and therefore lower physical separation resolution of the DES X-ray peak identification).

The redMaPPer algorithm computes a centring probability, $P_{\text{cen}}$, as an indicator of whether or not the selected central galaxy is the right choice. We do not find the values of $P_{\text{cen}}$ to accurately reflect the centring statistics of the redMaPPer sample. Specifically, we study the dependence of the centring performance by separately constraining the centring model of the redMaPPer SDSS sample in two $P_{\text{cen}}$ ranges, $\geq 0.9$ and $<0.9$, respectively. Clusters of $P_{\text{cen}} \geq 0.9$ have centring fraction of 0.76 ± 0.05 while clusters of $P_{\text{cen}} < 0.9$ have a notably lower centring fraction of 0.49 ± 0.09. Although a larger $P_{\text{cen}}$ value does indicate a better centring performance, it does not reflect the real centring fractions at the face values.

We have tested the dependence of the centring parameters through constraining the model with the SDSS sample in different ranges of richnesses, X-ray temperatures or luminosities, and for serendipitous versus targeted observations. We do not find significant differences in the centring parameter constraints. A larger set of X-ray observations would be needed to reveal any trends.

Interpretation of the results from this centring offset characterization depends on the adopted models. Different functional forms of the mis-centring offset – a Rayleigh distribution for the mis-centred component, and a Gaussian distribution for the well-centred clusters – have been attempted in previous studies (Saro et al. 2015; Rykoff et al. 2016; Hikage et al. 2017). We have tested alternative models (listed in Table 2) for the mis-centring and centring distributions with the Chandra SDSS sample. As the common parameter shared across different models, the well-centred cluster fraction, $\rho$, is consistently constrained to be in the range of 64 per cent to 70 per cent for SDSS.

We use the Bayesian deviance information criterion (DIC, Gelman et al. 2003) to compare the fitness of the models listed in Table 2. We compute the Bayesian DIC values (DIC, Gelman et al. 2003) by sampling the posterior constraints of the alternative and nominal models presented in the paper. DIC is computed as

$$
\text{DIC} = -2\log(p(\bar{x}_i|\theta)) + 2\log(p(x_i|\hat{\theta})).
$$

$log(p(\bar{x}_i|\theta))$ is the probability of the observed offsets averaged over the posterior mis-centring model, and $log(p(x_i|\hat{\theta})$ is the probability of the observed offsets given the best-fitting mis-centring model. Lower DIC values indicate better fitting of the model, and a DIC difference larger than 2 is considered significant. This DIC comparison strongly favours the nominal model described in this section.

Notably, when adopting a Rayleigh distribution for the mis-centred component, our SDSS posterior values on the well-centred cluster fraction, $\rho$, and the mis-centred characteristic offset, $\tau$ are highly consistent with the previous study in Hikage et al. (2017) which adopts a similar model and is based on a similar SDSS redMaPPer catalogue. The posterior precision of the parameters has improved significantly in our analysis.

### 3.3 Model validation

We check the model goodness-of-fit with a posterior prediction test. Specifically, we compare the fractions of clusters in an offset range from the data to the predictions from the constrained models. The procedures are as the follows.

(i) From the measurements of $x = r_{\text{offset}}/R_{\text{c}},$ record the number of clusters with offsets larger than a comparison value $x_0$.

(ii) Take one set of model parameters, $\rho$, $\sigma$, and $\tau$ from the MCMC posterior constraints, denoted as $\rho_i, \sigma_i$, and $\tau_i$. Randomly draw a set of centring offsets, $\{x_i\},$ from the offset model (equation 1) with the above set of posterior model parameters. The number of random draws should match the size of the X-ray redMaPPer sample being tested.

(iii) With the above set of centring offsets sampling, $\{x_i\},$ record the number of offsets larger than a comparison value $x_0$, $N(x_i > x_0)$.
4 MIS-CENTRING IMPACT ON RICHNESS SCALING RELATION

When a cluster is mis-centred, the redMaPPer richness estimation may become biased, and the bias depends on the mis-centring offset. In this section, we use the redMaPPer catalogue itself to constrain a model that describes the bias of \( \lambda \) upon a mis-centring offset.

As mentioned in Section 3, the redMaPPer algorithm selects the five most probable central galaxies and stores the \( \lambda \) estimations computed at each of the centres. The default redMaPPer centre is chosen as the one with the highest centring probability. We make use of this information to construct a richness shift versus offset model.

Assuming that there existed a redMaPPer catalogue with the most probable centres always being the correct ones, the \( \lambda \) estimations computed at the other four centres will be affected by the mis-centring offset. The \( \lambda \) bias between the real centre and each of the four remaining centre candidates, and the positional offsets between them, can be used to constrain a \( \lambda \) versus centring offset model.

We use the lambda offsets between the second and the first (redMaPPer default centre) most likely centres, and the distance offset between them, to model the richness versus centring offset (Section 4.1). The model is further validated with X-ray data (Section 4.2).

![Figure 5](https://academic.oup.com/mnras/article-abstract/487/2/2578/5490401)

### Table 3. Parameter constraints for the lambda offset versus centring offset model as in \( \gamma(x) = \exp(-x^2/a^2) \) and \( \sigma_{\gamma}(x) = a \times \arctan(bx) \).

<table>
<thead>
<tr>
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<th>SDSS ( a )</th>
<th>SDSS ( b )</th>
<th>SDSS ( a )</th>
</tr>
</thead>
<tbody>
<tr>
<td>Prior</td>
<td>(0, 10)</td>
<td>(0, 10)</td>
<td>(0, 10)</td>
</tr>
<tr>
<td>Posterior</td>
<td>1.64 ± 0.02</td>
<td>1.76 ± 0.07</td>
<td>0.241 ± 0.008</td>
</tr>
</tbody>
</table>

<table>
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<th></th>
<th>DES ( a )</th>
<th>DES ( b )</th>
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<tr>
<td>Prior</td>
<td>(0, 10)</td>
<td>(0, 10)</td>
<td>(0, 10)</td>
</tr>
<tr>
<td>Posterior</td>
<td>1.66 ± 0.06</td>
<td>1.43 ± 0.22</td>
<td>0.26 ± 0.04</td>
</tr>
</tbody>
</table>

### 4.1 Model and model constraints

We quantify the fractional shift of \( \lambda, y = \lambda_{\text{miscentred}}/\lambda_{\text{true}} \) due to mis-centring. This shift is quantified as being dependent on the scaled mis-centring offset in terms of \( r_x \). Upon trial of different analytical forms, we model the probabilistic distribution of \( y \) as a Gaussian distribution:

\[
y \sim N(\bar{\gamma}(x), \sigma_{\gamma}(x)),
\]

with the mean and the dispersion, \( \bar{\gamma}(x) \) and \( \sigma_{\gamma}(x) \), both depending on \( x \). The mean is positive, decreases with larger offsets starting from 1, and asymptotically reaches 0 for large offsets (Fig. 6). The dispersion is positive, increases with larger offsets starting from 0, and asymptotically reaches a constant for large offsets (Fig. 6).

Specifically, we choose \( \bar{\gamma}(x) \) to be a Gaussian function between \( y \) and \( x \), \( \bar{\gamma}(x) = \exp(-x^2/a^2) \) with \( a \) being a model parameter. \( \sigma_{\gamma}(x) \) is chosen as \( \sigma_{\gamma}(x) = a \times \arctan(bx) \) and \( a \) and \( b \) are model parameters.

The process of evaluating the above model goes as follows:

(i) We compute the separation between the first and the second most probable centres, \( r_{\text{sep}} \), assuming the cluster photometric redshift from the redMaPPer algorithm for both of the galaxies. We scale \( r_{\text{sep}} \) as \( r_{\text{sep}}/R_x \) to become quantity \( x \).

(ii) We compute the relative \( \lambda \) offset between the first and the second most probable centres as \( y = \lambda_{2a}/\lambda_{1a} \).

(iii) We repeat the above process for each cluster \( i \) in the redMaPPer catalogue, acquiring a measurement data set \( \{x_i\} \) and \( \{y_i\} \). This data set is shown in Fig. 6.

(iv) With the above measurement data set, we constrain the model parameters of \( y \) with the following likelihood:

\[
L = \sum_{x_i<0.1} \left( -\frac{\ln(\gamma(x_i))}{2\sigma_{\gamma}(x_i)^2} - \ln(\sigma_{\gamma}(x_i)) \right).
\]

The likelihood is sampled with an MCMC algorithm.

(v) Note that only the data points with \( x_i > 0.1 \) are used in the fitting process. This helps eliminate overfitting at small \( x \), and improves the fitting results at large \( x \).

The same measurements and modelling processes are performed for the SDSS and DES redMaPPer catalogues separately. Fig. 6 shows the data points as well as the best-fitting models in the first two columns. For comparison, the red lines/points are the data running averages and running dispersions in the \( x \) bins (bin widths indicated by the \( x \) error bars in the second column). The MCMC posterior constraints are shown in Fig. 7 and listed in Table 3. \( a \) does not appear to be covariant with \( a \) or \( b \), but \( a \) and \( b \) appear to be highly covariant.

Overall, the impact of mis-centring on cluster richness is mild with a low scatter at small mis-centring offset, but grows with a larger offset, reaching a bias ratio \( \lambda_{\text{miscentred}}/\lambda_{\text{true}} \) of 0.5 at \( r_{\text{sep}}/R_x \sim 1.40 \).
Figure 6. The richness versus offset distributions and models. The cluster richness tends to be biased lower by mis-centring (a, e), and the mean of the bias can be characterized by a Gaussian function (b, f). The biases have large dispersions (c, g), which can be further characterized by an arctan function (c, g). Posterior predictive checks show that the models are tightly constrained from the data and fit the data well (d, h).
4.2 Model validation with X-ray centres

As a test of the richness offset model, we make use of the Chandra-redMaPPer samples (Section 2.3) and compare the X-ray peak redMaPPer centred \( \lambda \) to the prediction from the model. We rerun the redMaPPer \( \lambda \) algorithm with the X-ray peaks as the cluster centres. The procedures are equivalent to the original \( \lambda \) estimation with the exception of a ‘percolation’ process (Rykoff et al. 2014), which re-evaluates \( \lambda \) upon masking neighbouring redMaPPer clusters. The \( \lambda \) estimations on X-ray peaks do not go through the ‘percolation’ process as the run does not consider redMaPPer clusters not present in the X-ray sample. In this test, to ensure that the ‘percolation’ process is negligible, we remove clusters whose \( \lambda \) changed by 10 per cent in the initial redMaPPer percolation process.\(^3\)

We calculate the \( \lambda \) offsets versus the distance offsets between the X-ray peaks and redMaPPer centres. The \( \lambda \) offsets are calculated as \( \lambda_i = \text{r}_{\text{RM}} \text{redMaPPer} / \text{r}_{\text{x-ray}} \), and the corresponding X-ray and redMaPPer miscentring offsets as \( x_i = \text{r}_{\text{RM}} / \text{r}_{\text{x-ray}} \), where \( r_j \) is evaluated with the X-ray centred \( \lambda \). Clusters of centring offsets \( x_i \) less than 0.1 are considered well-centred (a similar cut is applied when deriving the model in Section 4.1) and do not enter the test. In Fig. 8, we show the derived \( \lambda \) and centring offsets from the X-ray observations. The constrained models in Section 4.1 appear to be qualitatively consistent with these offsets.

To quantify the fitness of the model from the previous section, we compute the following \( \chi^2 \) discrepancy given model parameters and observations:

\[
\chi^2_n(\alpha, a, b) = \sum_{i=1}^{n} \left[ \frac{y_{x,i} - \bar{y}(x_{x,i})}{\sigma_{y_x}(x_{x,i})} \right]^2.
\]

As described in the previous section, \( \alpha \) is the parameter of the \( \bar{y}(x) \) function, and \( a \) and \( b \) are the parameters of the \( \sigma_{y_x}(x) \) function. If the model quantitatively describes cluster mis-centring correctly, and the values of \( \alpha, a, \) and \( b \) are accurate, \( \chi^2_n(\alpha, a, b) \) will appear to be drawn from a chi-squared distribution, \( \chi^2(n) \), with the degree of freedom, \( n \), matching the number of \( \{y_i, x_i\} \) observations. This chi-squared distribution is known as the posterior predictive density for the \( \chi^2 \) discrepancy.

We perform Bayesian posterior predictive assessment on the fitness of the model following the process in Gelman, Meng & Stern (1996) and Meng (1994). The process includes calculating a posterior predictive \( p \)-value (PPP value) which is the classical \( p \)-value averaged over the posterior model parameter distribution. With the posterior distribution of \( \alpha, a \) and \( b \) sampled with the MCMC method, the procedure goes as follows:

(i) Take one set of \( \alpha, a \), and \( b \) from the MCMC posterior constraints, donated as \( \alpha_j, a_j, \) and \( b_j \).

(ii) Calculate the \( \chi^2 \) discrepancy for \( \alpha_j, a_j, \) and \( b_j \), denoted as \( \chi^2_{j} \).

(iii) Repeat the process for each set of \( \alpha, a \), and \( b \) values from the MCMC posterior chain.

(iv) For each \( \chi^2_{j} \), randomly draw a value, \( q_j \), from a standard \( \chi^2(n) \) distribution. For the posterior set of \( \{\chi^2_{j}, q_j\} \), record the fraction of \( \chi^2_{j} \geq q_j \) as \( Pb_2 \), and the frequency of \( \chi^2_{j} \leq q_j \) as \( Pb_1 \). We use a two sided \( p \)-value definition, that \( Pb = \min(Pb_1, Pb_2) \) as the PPP value.

We compute the posterior \( \chi^2 \) discrepancy given X-ray \( \{y_i, x_i\} \) observations shown in Fig. 8. The posterior \( \chi^2 \) discrepancy values are expected to occupy a highly probable interval of a chi-squared distribution, \( \chi^2(n) \) (posterior predictive density), with the degree of freedom \( n \) matching the number of \( \{y_i, x_i\} \) observations from X-ray peak centres. This appears to be the case for both the SDSS and DES X-ray redMaPPer sample. The Bayesian posterior predictive \( P \)-values (PPP values) are 0.18 and 0.50 respectively for the SDSS and DES X-ray redMaPPer samples. Both of the PPP values are above a 0.025 model rejection threshold, indicating consistency between the constrained models and the offsets derived with X-ray peak centres.

We also compute the posterior \( \chi^2 \) discrepancy for the original \( \{y_i\}, \{x_i\} \) observations that are used to constrain the model. The distributions of \( \{\chi^2_{j}\} \) are shown as the grey shaded histogram in Fig. 6, along with the probability density of a chi-squared distribution, \( \chi^2(n) \), for comparison. The distribution of \( \{\chi^2_{j}\} \) is in good accordance with the expected \( \chi^2(n) \) distribution, for both the SDSS and DES redMaPPer samples, indicating the goodness of the fit and tightness of the model parameter constraints. The PPP values for the SDSS and DES samples are 0.50 and 0.49, respectively, matching the expectation of a well-posited model.

Note that when deriving the model, the available redMaPPer catalogues already have imperfect centre selections. According to the previous section, the majority (~70 per cent) of the clusters in the redMaPPer catalogue are correctly selected. We have attempted to select clusters with a higher centring probability (using the redMaPPer \( P_{\text{cen}} \) quantity), but the samples selected on \( P_{\text{cen}} \) display a hint of performing slightly worse in the validation test with X-ray data. Because of a concern that this \( P_{\text{cen}} \) selection may have biased the cluster sample, and that the selection still cannot ensure a 100 per cent well centred subsample, we do not apply the selection in this paper and emphasize on our derived model passing the validation test with X-ray data. In the future, it would be desirable to quantify the richness offset model with X-ray centred richness when larger X-ray redMaPPer clusters become available.

\(^3\)Two from each of the SDSS and DES Y1 samples.
5 IMPLICATIONS FOR DES CLUSTER COSMOLOGY

5.1 Mis-centring model in DES cluster weak lensing analysis

In DES stacked cluster lensing studies, the cluster lensing signals are fitted to an analytic model to determine cluster mass and concentration. To summarize, McClintock et al. (2019) adopts the following method to correct for mis-centring:

\[
\Sigma(r|M, c) = \rho \Sigma_{\text{cent}}(r|M, c) + (1 - \rho) \Sigma_{\text{miscent}}(r|M, c).
\]  

(7)

In the above equations, \( \Sigma(r|M, c) \) is the cluster mass profile model with mass \( M \) and concentration \( c \). \( \Sigma_{\text{cent}}(r|M, c) \) denotes the mass profile model for well-centred clusters, while \( \Sigma_{\text{miscent}}(r|M, c) \) is the mass profile model for mis-centred clusters. The mis-centred profile is averaged over the angle, \( \theta \), and magnitude, \( R \), of the radial vector to the correct centre,

\[
\Sigma_{\text{miscent}}(\vec{r}|M, c) = \Sigma_{\text{cent}}(\vec{r} + \vec{R}|M, c),
\]

\[
\Sigma_{\text{miscent}}(r|M, c) = \frac{1}{2\pi} \int d\phi dR P_{\text{miscent}}(R)
\]

\[
\times \Sigma_{\text{cent}}(\sqrt{r^2 + R^2 + 2rR\cos(\theta)}|M, c).
\]

The distribution for the magnitude of the radial offset \( R \) is \( P_{\text{miscent}}(R) \), which is described by a parameter \( \tau \). The model does not account for the offset between central galaxies and X-ray peaks because haloes are assumed to be centred on a massive dark matter substructure hosting a central galaxy. Based on the analyses carried out in this study and a companion SDSS redMaPPer centring study of a complete sample with Swift observations in von der Linden et al. (in preparation) and McClintock et al. (2019) adopted the prior \( \rho = 0.75 \pm 0.08 \) and \( \tau = 0.17 \pm 0.04 \). These values are consistent with the Chandra DES constraints presented in Section 3, but also encompass the results of the SDSS samples presented in this paper and in von der Linden et al. (in preparation).

5.2 Sensitivity of cluster mass estimation to the mis-centring model

We determine the sensitivity of the mass calibration to variations in the values of the mis-centring parameters. To do so, we create a fiducial mass profile and analyse how much the measured masses deviate from the truth if the mis-centring model is inaccurate. Following the recipe in McClintock et al. (2019), the fiducial mass profile model combines an NFW profile and a two-halo model of \( M_{200} = 10^{14} M_{\odot} \), concentration of 5 and \( R_5 = 1 \) Mpc. We compute the mis-centred lensing signal by adopting fiducial values \( \rho = 0.68 \) and \( \tau = 0.15 \). We fit this synthetic weak lensing data with a minimum \( \chi^2 \) method assuming a range of values for both \( \rho \) and \( \tau \), and measure the bias of the best-fitting mass and concentration as a function of these two parameters. The fitting process is restricted to the 0.2–30 Mpc radius range as in the DES Y1 weak lensing study (McClintock et al. 2019) and the profile measurement uncertainty is assumed to be due to shape noise only, and therefore scales with radius as \( r^{-1} \).

Fig. 9 shows the best-fitting mass and concentration parameters as a function of \( \rho \) and \( \tau \). We find that the best-fitting mass is insensitive to the assumed \( \rho \) value, whereas the recovered concentration is biased. Allowing the concentration parameter to vary effectively decouples the recovered mass from \( \rho \). By contrast, variations in \( \tau \) have a non-negligible impact on the best-fitting mass. Uncertainties in \( \tau \) at a level of \( \pm 0.04 \) are comparable to the...
constraint in this paper, results in a mass uncertainty of $\pm 0.015$ dex.\footnote{0.015 dex means $\delta \log M_{200c} = 0.015$.}

These results are in qualitative agreement with the results of McClintock et al. (2019), though computed with two important methodological differences. Specifically: (1) McClintock et al. (2019) has considered more systematic effects other than mis-centring, including using a semi-analytic covariance matrix (Gruen et al. 2015), and accounting for boost-factor corrections (Varga et al. 2018). These changes will affect the relative weighting of radial inner to outer scales, thereby impacting the sensitivity of the mass posteriors to the mis-centring parameters. (2) When constraining the richness-mass scaling relation parameters, McClintock et al. (2019) treats the mis-centring parameters for each richness and redshift bin as independent, which reduces the relative importance of mis-centring in their analysis. Together, these differences reduce the sensitivity of the scaling relation amplitude from the 0.015 dex we estimated here to 0.78 per cent, as quoted in McClintock et al. (2019). Nevertheless, it is clear from fig. 10 in McClintock et al. (2019) that the mass posterior in a single bin is largely insensitive to $\rho$, but is degenerate with $r$, as illustrated in our toy model analysis above.

These conclusions, however, rely on the assumption that the cluster mass profile is not correlated with the cluster mis-centring effect in optical data. Future cluster lensing analysis may wish to further investigate this assumption, e.g. through examining the cluster mass distribution in X-ray selected clusters (Das et al., private communication).

5.3 Cluster abundance

The $\lambda$ offset caused by cluster mis-centring introduces bias and scatter into the lambda–mass scaling relation. We study the scatter increase with a test based on a N-body dark matter simulation (Habib et al. 2016). Richnesses are prescribed to each of the simulation dark matter haloes following the richness–mass relation in Saro et al. (2015), with a richness scatter, $\sigma_{\ln(M)}$, of 25 per cent. We perturb the assigned richnesses with the Chandra SDSS offset model presented in Section 3 and the richness bias model presented in this section, and find the richness scatter, $\sigma_{\ln(M)}$, to increase by 2 per cent.

The bias and scatter manifest themselves in the number count of clusters selected by $\lambda$, which is a fundamental input to cluster abundance cosmology. Mis-centring tends to lower the richness estimation and the numbers of clusters above a richness threshold selected by the mis-centred richnesses would be lower than those selected by richnesses without mis-centring. The average mass of the clusters selected by the mis-centred richnesses tend to be higher. Testing with simulations shows that the numbers of clusters selected by the mis-centred richnesses is lower by $\sim 2$ per cent, and the average cluster masses increase by $\sim 0.5$ per cent. Costanzi et al. (2018) estimates these shift to have negligible effects for SDSS and DES Y1 cluster cosmology constraints since they are significantly smaller than the remaining systematic uncertainties of cluster abundance and mass estimations. Nevertheless, Costanzi et al. (2018) and the DES Collaboration (in preparation) corrected the data vectors with the factors listed above to account for the mis-centring effect. Future cluster cosmology analysis in the coming years of DES and LSST may wish to explicitly incorporate the mis-centring richness relations as the mis-centring effect on cluster abundance becomes more substantial compared to the statistical uncertainty.

6 CENTRING STUDY WITH XMM DATA

During the preparation of this paper, an additional sample of redMaPPer clusters (SDSS and DES) with archival X-ray observations from the XMM–Newton space telescope became available. As the centring analyses in this paper are optimized for Chandra observations, especially the centring offset model for well-centred clusters is optimized for Chandra PSFs, we do not attempt to combine the XMM and Chandra observations in the modelling processes. Rather, we use the XMM defined X-ray centres to explore the robustness of the fits presented in Section 3. The sample selection and the XMM data analysis methods are described in Section 6.1, and Section 6.2 describes the centring comparison results from XMM.

6.1 X-ray data processing

The redMaPPer–XMM joint samples (SDSS and DES) were constructed as follows. First, the redMaPPer centroids were compared to the aim points of observations in the XMM public archive.\footnote{The archive match used in this analysis was carried out on 2018 August.} redMaPPer clusters with centroids falling outside 13 arcmin of...
an aim point were excluded from the samples. Secondly, the mean and median XMM exposure time was determined within a 10 arcsec radius of the redMaPPer centroid. For this we used exposure maps produced by the XMM Cluster Survey (XCS, Lloyd-Davies et al. 2011). Any redMaPPer clusters with mean exposure times of <3 ks, and/or median exposure times of <1.5 ks, were excluded from the samples. We note that the median filter was necessary because some clusters straddle both active and inactive regions of the field of view (FOV), e.g., those lying close to the FOV edge. If a given cluster was observed multiple times by XMM, only the observation with the longest exposure time, at the redMaPPer centroid, was used in subsequent analyses. Third, the remaining redMaPPer cluster centroids were compared to the list of extended sources detected using the XCS Automated Pipeline Algorithm (XAPA). Any redMaPPer clusters lying further than 2 h⁻¹ Mpc (assuming the redMaPPer redshift) of such a source were excluded from the samples. At this stage, the SDSS and DES redMaPPer–XCS samples comprised of 356 and 282 clusters, respectively.

The X-ray peaks for the clusters in the redMaPPer–XMM joint samples were then determined using a method that closely follows that used for the Chandra analyses, as described in Section 2.3. An initial peak location was found in the respective merged (PN, MOS1, MOS2) XMM image, after smoothing with a σ = 50h⁻¹ kpc Gaussian (assuming the redMaPPer redshift). As with the Chandra analysis, other sources (i.e. those assumed not to be associated with the redMaPPer cluster) are masked out before the smoothing takes place. For this, we use the XAPA source catalogue of point-like sources. If there are multiple XAPA extended source within the image, then only the one closest to the redMaPPer centroid is left unmasked. The peak location is then the brightest pixel in the masked, smoothed image, within a radius 1.5 × Rₗ of the RM position.

The initial peak selection was occasionally erroneous. For example when there was a very bright point source in the XMM FOV. Such sources ‘bleed’ into the surrounding region of the detector, meaning that the default point source mask size was not large enough to remove them completely. Such cases were easily identified by eye using SDSS (or DES) and XMM ‘postage stamp’ images. Most of these cases could be corrected by adjusting the size of the point source mask, and then re-running the peak finding script. However, for some, the point source ‘bleeding’ was so pronounced that the respective cluster had to be removed from the redMaPPer–XMM joint sample. Another reason for the initial peak being erroneous was the mis-percolation issue described in Section 2.3. The mis-percolation cases were also identified using eye-ball checks. For these it was necessary to adjust the extended source mask, so that the closest source was now masked, but the second closest was not.

Once the second run of peak finding has been completed (and the new peak positions have been confirmed by eye), the remaining clusters were eye-balled again. At this stage, more redMaPPer clusters were removed from the sample: (i) those where the XAPA extended source is clearly not associated with the redMaPPer cluster (i.e. it is a foreground background cluster in projection), and (ii) those where the redMaPPer central galaxy falls in an XMM chip gap. After the various cuts described above, the SDSS–XCS and DES–XCS samples contained 248 and 109 sources, respectively.

We further apply off axis angle and SNR cuts to the XMM samples. The SNR was determined in the same way as for the Chandra sample i.e. within 500h⁻¹ kpc, using 0.5–2.0 keV XMM images. For the XMM–SDSS sample, we require the detections S/N to be >6.5, and the redMaPPer centres to be within 8.5 arcmin of the aim point (or 6.5 arcmin away from the FOV edge assuming a 15 arcmin FOV radius). For the XMM–DES sample, we again require the detections S/N to be >6.5, but allow the redMaPPer centres to be up to 10.5 arcmin of the aim point. The FOV cuts ensure that the corresponding redMaPPer centres are more than 500 kpc away from the XMM FOV at z = 0.1 for the SDSS sample, or at z = 0.2 for the DES sample. The S/N cuts were imposed to match the 6.5 S/N cuts in the Chandra analysis. With these FOV and S/N cuts, the final SDSS and DES XMM samples contain 163 and 66 clusters, respectively. Fig. 10 shows their richness and mass distributions.

Further details of the redMaPPer–XMM joint sample development, the peak measurements, the signal-to-noise estimation and the individual mis-percolation cases can be found in Giles et al. (in preparation).

6.2 The XMM–Chandra and XMM–redMaPPer offsets

A subsample of the redMaPPer clusters, 54 in the SDSS sample, and 25 in the DES sample, are analysed by both the XMM and Chandra analyses. With these overlapping cases, we compare the XMM peak measurements to those from Chandra. Fig. 11 shows the offset distribution between XMM and Chandra peak identifications for the same redMaPPer clusters, scaled by their Rₗ. The XMM and Chandra peak identifications are highly consistent: their separations are within 0.05 Rₗ for 53/54 of the overlapping SDSS clusters, and 20/25 of the overlapping DES clusters. The separations have a wider distribution for the DES redMaPPer sample reflecting its higher redshift range, and hence higher X-ray peak identification uncertainties in terms of physical distances.

Given the consistency between XMM and Chandra X-ray peak identifications, we constrain the redMaPPer centring offset model proposed in Section 3 with the XMM peaks. Figs 12 and 13 respectively show the offset distributions between the XMM peaks and the redMaPPer centres for the SDSS and DES samples, with comparisons to their corresponding Chandra offset distributions. Figs 14, 15, and Table 4 show the model parameter constraints from these XMM samples. The σ parameter in our redMaPPer centring offset model (equation 1) represents the X-ray peak offset to cluster central galaxy for well-centred clusters, which is further smeared by X-ray peak identification uncertainty and X-ray telescope PSFs. Since σ is in the unit of physical distance, the σ difference can be driven by the different angular resolutions of XMM and Chandra at low redshift, and other X-ray peak identification uncertainties, and X-ray peak-galaxy centre separations at higher redshift. Hence, we do not expect this parameter to agree between XMM and Chandra.
Figure 11. The $R_\lambda$ [$R_\lambda = (\lambda/100)^{0.2} h^{-1} \text{Mpc}$] scaled offset distribution between the Chandra and XMM peak identifications for the same redMaPPer clusters.

Figure 12. The $R_\lambda$ [$R_\lambda = (\lambda/100)^{0.2} h^{-1} \text{Mpc}$] scaled offset distribution between the redMaPPer centres and the X-ray emission peaks for the redMaPPer SDSS samples from the XMM archival observations. The distribution can be fitted with two components – a concentrated component that represents the well centred redMaPPer clusters, and an extended component that represents the mis-centred redMaPPer clusters. The best-fitting SDSS offset model is shown as the solid lines (black: well-centred model, red: mis-centred model), with the shaded regions representing the uncertainties. As a comparison, we also show the corresponding offset distribution from the analysis with Chandra archival data (blue dashed histogram).

Figure 13. The $R_\lambda$ [$R_\lambda = (\lambda/100)^{0.2} h^{-1} \text{Mpc}$] scaled offset distribution between the redMaPPer centres and the X-ray emission peaks for the redMaPPer DES samples from the XMM archival observations. The distribution can be fitted with two components – a concentrated component that represents the well centred redMaPPer clusters, and an extended component that represents the mis-centred redMaPPer clusters. The best-fitting DES offset model is shown as the solid lines (black: well-centred model, red: mis-centred model), with the shaded regions representing the uncertainties. As a comparison, we also show the corresponding offset distribution from the analysis with Chandra archival data (blue dashed histogram).

Figure 14. Centring offset parameter constraints (equation 1) for the XMM SDSS (grey) redMaPPer samples. 76 ± 6 per cent of the SDSS redMaPPer clusters appear to be well centred (indicated by the $\rho$ parameter). For the mis-centred clusters, their mis-centring offsets is characterized by a Gamma distribution with a characteristic offset (the $\tau$ parameter) of 0.16 ± 0.03 $R_\lambda$. As a comparison, we also show the corresponding model constraints from the SDSS analysis with Chandra archival data (blue). The prior central values adopted in DES cosmological analyses as described in Section 5 are shown as the black dashed lines.

Figure 15. Centring offset parameter constraints (equation 1) for the XMM DES (grey) redMaPPer samples. 74 ± 10 per cent of the DES redMaPPer clusters appear to be well centred (indicated by the $\rho$ parameter). For the mis-centred clusters, their mis-centring offsets is characterized by a Gamma distribution with a characteristic offset (the $\tau$ parameter) of 0.20 ± 0.07 $R_\lambda$. As a comparison, we also show the corresponding model constraints from the DES analysis with Chandra archival data (blue). The prior central values adopted in DES cosmological analyses as described in Section 5 are shown as the black dashed lines.
and the $\sigma$ difference is especially larger for the lower redshift SDSS redMaPPer samples. For the other two parameters of the model, $\rho$ and $\tau$, which respectively represent the well-centred fractions of redMaPPer and the centring offset of mis-centred redMaPPer clusters, the constraints from XMM are consistent with those from Chandra for both the SDSS and DES redMaPPer samples within two standard deviations.

Notably, no selection cuts in the XMM analysis were made to make its centring model constraints better match the Chandra results, yet their centring offset results are consistent with each other. We conclude that the redMaPPer mis-centring offset modelling presented in this paper are robust upon investigation with archival XMM data.

### 7 Summary

This analysis makes use of the archival X-ray observations to constrain the centring performance of the redMaPPer cluster finding algorithm. We calibrate the well-centred fraction of redMaPPer clusters for both the SDSS and DES samples with the X-ray emission peaks. The offsets between the redMaPPer centres and X-ray peaks are well modelled by a two-component distribution, which indicates that $69^{+1.5}_{-1.3}$ per cent and $83.5_{-7.5}^{+15}$ per cent of the clusters are well centred in the SDSS and DES samples. The offset distribution of the mis-centred redMaPPer clusters are modelled with a Gamma distribution, and cluster mass modelling appears to be more sensitive to the accuracy of these mis-centring offsets than the mis-centring fraction.

With the upcoming DES Year 3 and Year 5 data, we expect the redMaPPer centring constraints to continue improving with $\sim 2$ times larger overlapping samples between DES and archival X-ray observations, which may permit us to quantify the dependence of the centring parameters on cluster properties, such as cluster richness, redshift, X-ray temperature, and luminosity. The current improvement has already lowered the cluster weak lensing mass modelling uncertainties due to mis-centring, to the extent of being in-substantial comparing to the other modelling systematic effects. Since mis-centring is often assumed to be uncorrelated with cluster mass distributions in weak lensing analyses, with the anticipated level of improvement, one may wish to investigate the correlations of mis-centring with other cluster lensing systematic effects, such as cluster mass modelling uncertainties, cluster orientation, triaxiality, and projection (McClintock et al. 2019).

The cluster richness estimates tend to be biased lower by mis-centering. In this paper, we propose a richness bias model to describe the effect, which is validated by X-ray centred richness measurements. The richness bias is offset dependent, low for clusters with small mis-centering offset, but larger than 50 per cent for severely mis-centred clusters. Cluster cosmology studies based on full depth DES data or LSST data should explicitly account for this effect to avoid biased cosmological parameter inferences.

Code used in this analysis is available from [https://github.com/yzyhang/center_modeling_yz](https://github.com/yzyhang/center_modeling_yz).

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### Table 4. Centring offset Parameter constraints (equation 1) for the XMM DES and SDSS redMaPPer samples.

<table>
<thead>
<tr>
<th>Prior</th>
<th>$\rho$</th>
<th>$\sigma$</th>
<th>$\tau$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\rho$</td>
<td>[0.3, 1]</td>
<td>$[0.0001, 0.1]$</td>
<td>$[0.08, 0.5]$</td>
</tr>
<tr>
<td>XMM SDSS posterior</td>
<td>$0.781^{+0.055}_{-0.038}$</td>
<td>$0.0432^{+0.0063}_{-0.0059}$</td>
<td>$0.201^{+0.026}_{-0.039}$</td>
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<tr>
<td>XMM DES posterior</td>
<td>$0.815^{+0.059}_{-0.085}$</td>
<td>$0.053^{+0.012}_{-0.011}$</td>
<td>$0.185^{+0.066}_{-0.041}$</td>
</tr>
</tbody>
</table>
REFERENCES

Habib S. et al., 2016, New Astron., 42, 49
McDonald M. et al., 2012, Nature, 483, 349
Oguri M. et al., 2018, PASJ, 70, 520
Rykoff E. S. et al., 2016, ApJS, 224, 1

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