Galaxy and Mass Assembly (GAMA): merging galaxies and their properties

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ABSTRACT

We derive the close pair fractions and volume merger rates for galaxies in the Galaxy and Mass Assembly (GAMA) survey with \(-23 < M_r < -17\) \((\Omega_M = 0.27, \Omega_\Lambda = 0.73, H_0 = 100 \text{ km s}^{-1} \text{ Mpc}^{-1})\) at \(0.01 < z < 0.22\) (look-back time of \(<2\) Gyr). The merger fraction is approximately 1.5 per cent \(\text{Gyr}^{-1}\) at all luminosities (assuming 50 per cent of pairs merge) and the volume merger rate is \(\approx 3.5 \times 10^{-4} \text{ Mpc}^{-3} \text{ Gyr}^{-1}\). We examine how the merger rate varies by luminosity and morphology. Dry mergers (between red/spheroidal galaxies) are found to be uncommon and to decrease with decreasing luminosity. Fainter mergers are wet, between blue/discy galaxies. Damp mergers (one of each type) follow the average of dry and wet mergers. In the brighter luminosity bin \((-23 < M_r < -20)\), the merger rate evolution is flat, irrespective of colour or morphology, out to \(z \sim 0.2\). The makeup of the merging population does not appear to change over this redshift range. Galaxy growth by major mergers appears comparatively unimportant and dry mergers are unlikely to be significant in the build-up of the red sequence over the past 2 Gyr. We compare the colour, morphology, environmental density and degree of activity (BPT class, Baldwin, Phillips & Terlevich) of galaxies in pairs to those of more isolated objects in the same volume. Galaxies in close pairs tend to be both redder and slightly more spheroid dominated than the comparison sample. We suggest that this may be due to ‘harassment’ in multiple previous passes prior to the current close interaction. Galaxy pairs do not appear to prefer significantly denser environments. There is no evidence of an enhancement in the AGN fraction in pairs, compared to other galaxies in the same volume.

Key words: galaxies: formation – galaxies: interactions.

1 INTRODUCTION

Mergers and interactions are believed to be among the primary pathways for galaxy formation and evolution: they are expected to drive star formation histories, morphologies, internal kinematics, chemical evolution and nuclear activity, among others, and reflect a long history of processing through different environments.\(^1\) The merger rate of galaxies, its evolution and its dependence on mass, luminosity, colour, morphology and environment (to name some) provide important clues to theories of galaxy formation and are an

\(^1\) The literature on this subject is very large and growing by the day; we cannot hope to be just or comprehensive in a brief introduction. However, Baugh (2006) and the recent textbook by Mo, van den Bosch & White (2010) provide good entry points to this subject.
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2 DATA SET AND METHODOLOGY

The data set we use is the GAMA redshift survey. This is fully described in Driver et al. (2011), see Baldry et al. (2010) for the input catalogue and Hopkins et al. (2013) for details of the spectroscopic analysis); in the following, we give a brief summary of its main properties as used in this paper. GAMA as used in this paper is equivalent to the GAMA-I sample obtained during the first three years of the survey and frozen for internal team use and consists of three 48 deg$^2$ regions on the celestial equator, at RA of 9 h, 12 h and 15 h, containing photometry over a wide range of wavelengths, including vacuum UV from GALEX, u, g, r, i, z from the SDSS (York et al. 2000; Ahn et al. 2012), YJHK from 2MASS and UKIDSS, mid-IR from the WISE survey at 3.6, 4.5, 12 and 22 µm, as well as other sources of archival photometry (e.g. Herschel, Planck). Spectroscopy to a limiting magnitude of $r = 19.4$ ($r = 19.8$ for the 12$^h$ region) has been obtained for over 120 000 galaxies with the AAOmega multifibre spectrograph on the Anglo-Australian Telescope, supplementing existing data sets from the 2dF Galaxy Redshift Survey (Colless et al. 2001), the SDSS and others, to reach a high degree (>97 per cent) of spectroscopic completeness (Hopkins et al. 2013). Each region in the sky is surveyed multiple times, which aids in reaching high completeness at small angular separations. Because of this, GAMA is ideal to carry out a close pair analysis of galaxy populations to measure the merger rate and explore the influence of interactions and mergers on galaxy properties.

We search for pairs in GAMA using the methods developed by Patton et al. (2000, 2002) and used in our previous work on the subject (De Propris et al. 2005, 2007, 2010). Following previous studies, we define galaxies as being in a close pair if their projected separation on the sky is $5 < r_p < 20$ kpc (the lower limit is imposed to avoid selecting high surface brightness regions within galaxies, such as H II regions) and if their velocity difference is $|\Delta V| < 500$ km s$^{-1}$. If these objects are truly bound, we expect that they will merge, by dynamical friction in $<1$ Gyr (the actual time-scales are likely to be somewhat longer, see below for details).

Fig. 1 shows the absolute magnitude of galaxies in GAMA (only a subset is shown for clarity) as a function of redshift. Here, all magnitudes and $k + e$ corrections are as given in the GAMA-I
Figure 1. Galaxies in the GAMA survey (we plot only a small subset for reasons of clarity) plotted as absolute magnitude versus redshift (top panel). The coloured boxes on the plot show the volume-limited samples that we search for pairs in. The limits in magnitude and redshifts are shown in Table 1. The two orange boxes show the adopted samples for galaxies with \(-23 < M_r < -20\) and \(0.0100 < z < 0.1724\) and \(0.1724 < z < 0.2193\), respectively. The two redshift intervals are chosen to yield approximately \(z_{\text{av}}\) at magnitude for a slowly evolving galaxy (discussed in text) having \(r\) at \(0.0100 < z < 0.0100\)

\[N_{\text{pairs}} = \sum_{i} w_{i} N_{c,i},\]

where the sums are over the \(i = 1, \ldots, N_{1}\) galaxies in the primary sample, and the total companion luminosity is

\[L_{c} = \sum_{i} w_{i} L_{c,i} = \sum_{i} \sum_{j} w_{i} w_{j} L_{c,i,j},\]

where \(N_{c,i}\) and \(L_{c,i}\) are the number and luminosity (respectively) of galaxies in the secondary sample which are close companions (by the above definition) to the \(i\)th galaxy in the primary sample:

\[N_{c,i} = \sum_{j} w_{j} N_{c,i,j} = \sum_{j} \frac{w_{j} w_{c,i,j}}{S_{c}(z_{j})},\]

and

\[L_{c,i} = \sum_{j} w_{j} L_{c,i,j} = \sum_{j} \frac{w_{j} w_{c,i,j}}{S_{c}(z_{j})} L_{j,i},\]

where the sums run over the \(j\) secondary galaxies that fulfil the criteria of being dynamically close to the \(i\)th primary galaxy. The raw pair fractions are then weighted to correct for sources of spatial and spectroscopic incompleteness (e.g. areas where no fibres were placed such as within the haloes of bright stars, or limits to slit placement for close companions).

The components \(S_{c}(z)\) and \(S_{l}(z)\) are weights to correct for the change in density of the secondary galaxies as a function of redshift due to the apparent flux limit of the sample and convert a luminosity-limited sample into what would be expected from a volume-limited sample. They are calculated as the ratio of the integrated luminosity function over the luminosity sampled to the integrated luminosity

catalogue (Driver et al. 2011). The first step in our analysis is to select galaxies within a series of volume-limited boxes in absolute magnitude and redshift space. In order to measure a pair fraction, we need to restrict ourselves to a relatively small range of luminosities, as we need objects to have similar clustering properties (Patton et al. 2000; Patton & Atfield 2008) and the clustering length is known to vary as a function of galaxy luminosity (Norberg et al. 2001). However, we also need our samples to contain enough galaxies to determine the pair fraction to sufficient accuracy, given that only \(<3\) per cent of galaxies (based on previous work) are likely to fulfil our definition of a close pair. As in Patton et al. (2000, 2002) and De Propris et al. (2005, 2007), we draw ‘volume-limited’ samples spanning 3 mag in luminosity (Fig. 1). Potentially, this includes mergers with luminosity ratios of up to 1:10, but the majority of close pairs will have luminosity ratios of \(<1: 4\) and are therefore ‘major’ mergers (see below). Table 1 shows the luminosity and redshift ranges we study in this work, as well as the numbers of objects in each volume-limited box and the number of galaxies in close pairs (this need not be an even number, because of possible triplets). Note that for the brighter luminosity bin we consider, we are able to draw two samples, at \((z) = 0.09\) and \((z) = 0.19\), and study the evolution of the merger rate between these two redshifts for galaxies with \(-23 < M_r < -20\).

Because our sample (like all samples) is luminosity limited, we need to impose a lower absolute magnitude limit to our search, rather than using all galaxies in the appropriate luminosity range. This corresponds to the luminosity of a slowly evolving galaxy having the apparent magnitude limit of the survey \((r = 19.4)\) at all redshifts considered (thick red line in Fig. 1). If this is not done, it is possible for galaxies to enter or leave the sample depending on their star formation history (Patton et al. 2000). We use a Bruzual & Charlot 2003 model with \(z_f = 3\), \(Z = Z_{\odot}\) and \(\tau = 1\) Gyr, evolving passively, as a template for the slowest evolving object we are likely to encounter. When drawing the volume-limited boxes shown in Fig. 1, this model sets the high redshift limit for each luminosity interval we examine.

Each pair consists of a primary and secondary member; in this case we choose to make the two samples coincide so each galaxy is counted twice, once as a central and next as a satellite of the other. Following Patton et al. (2000, 2002) and our previous work, the number of close companions per galaxy is

\[N_{c,i} = \sum_{j} w_{j} N_{c,i,j} = \sum_{j} \frac{w_{j} w_{c,i,j}}{S_{c}(z_{j})},\]

and

\[L_{c,i} = \sum_{j} w_{j} L_{c,i,j} = \sum_{j} \frac{w_{j} w_{c,i,j}}{S_{c}(z_{j})} L_{j,i},\]

where the sums run over the \(j\) secondary galaxies that fulfil the criteria of being dynamically close to the \(i\)th primary galaxy. The raw pair fractions are then weighted to correct for sources of spatial and spectroscopic incompleteness (e.g. areas where no fibres were placed such as within the haloes of bright stars, or limits to slit placement for close companions).

The components \(S_{c}(z)\) and \(S_{l}(z)\) are weights to correct for the change in density of the secondary galaxies as a function of redshift due to the apparent flux limit of the sample and convert a luminosity-limited sample into what would be expected from a volume-limited sample. They are calculated as the ratio of the integrated luminosity function over the luminosity sampled to the integrated luminosity

<table>
<thead>
<tr>
<th>Luminosity range (mag)</th>
<th>Redshift range</th>
<th>(N_{\text{obj}})</th>
<th>(N_{\text{pairs}})</th>
<th>Pair fraction (per cent)</th>
<th>Volume merger rate ((10^{-3} \text{Mpc}^{-2} \text{Gyr}^{-1}))</th>
</tr>
</thead>
<tbody>
<tr>
<td>(-23 &lt; M_r &lt; -20)</td>
<td>(0.0100 &lt; z &lt; 0.1724)</td>
<td>19 816</td>
<td>644</td>
<td>3.30 ± 0.10</td>
<td>3.33 ± 0.11</td>
</tr>
<tr>
<td>(-23 &lt; M_r &lt; -20)</td>
<td>(0.1724 &lt; z &lt; 0.2193)</td>
<td>17 093</td>
<td>494</td>
<td>3.03 ± 0.14</td>
<td>2.60 ± 0.12</td>
</tr>
<tr>
<td>(-22 &lt; M_r &lt; -19)</td>
<td>(0.0100 &lt; z &lt; 0.1442)</td>
<td>20 936</td>
<td>699</td>
<td>3.39 ± 0.13</td>
<td>4.39 ± 0.17</td>
</tr>
<tr>
<td>(-21 &lt; M_r &lt; -18)</td>
<td>(0.0100 &lt; z &lt; 0.0939)</td>
<td>8084</td>
<td>256</td>
<td>3.24 ± 0.21</td>
<td>4.26 ± 0.28</td>
</tr>
<tr>
<td>(-20 &lt; M_r &lt; -17)</td>
<td>(0.0100 &lt; z &lt; 0.0607)</td>
<td>3151</td>
<td>56</td>
<td>1.84 ± 0.24</td>
<td>2.51 ± 0.33</td>
</tr>
</tbody>
</table>
function in a hypothetical volume-limited sample over the same redshift. In our case, as we are using volume-limited subsets of the data (see Fig. 1), these components are equal to 1 (no weight).

The other two weights, \( w_b \) and \( w_s \), correct for boundary effects. The first one takes care of areas (within the search radius) around each galaxy where no companion can be found because they lie beyond the edges of the survey on the sky or within areas occulted by bright stars, ‘drill holes’ (for guide probes), image defects and satellite trails or other cosmetic issues. This weight is the reciprocal of the fraction of the \( \pi r^2 \) area around each galaxy which is occulted in some fashion. The second weight is for galaxies close (within 500 km s\(^{-1}\)) to the redshift limits of each volume-limited box, and accounts for possible companions missed because they lie just beyond these boundaries. As in Patton et al. (2000, 2002), the value of this weight is set at 2. Similar weights need to be applied to the primary sample as well, to account for missed primaries in the same manner (recall that both primary and secondary samples are the same):

\[
w_{b1}^i = w_{b1} w_{v1} S_{b1}(z),
\]

\[
w_{s1}^i = w_{s1} w_{v1} S_{s1}(z),
\]

these weights being the reciprocals of those applied to the secondary sample \( w_{b2}^i, w_{v2}^i \) (Patton et al. 2000).

Although GAMA is highly complete and is intended to have no pair bias at close angular separations, we need to correct for potential pairs missed because of redshift incompleteness, especially at small separations. We carry this out in the following manner: we first select potential pairs with galaxies that have no redshifts by the projected separation criterion only, and if such a pair exists, we assign to the galaxy the same redshift and \( k + e \) corrections as its primary and require that the secondary galaxy falls within the selection criteria in luminosity. We then estimate the fraction of ‘true’ pairs in this sample by carrying out the same analysis on the photometric sample (only) and on the redshift sample (only) separately, by using only the projected separation criterion. Since we know the true pair fraction in the redshift sample, we can use the ratio between the pair fractions in the ‘photometric’ and ‘redshift’ samples to derive an incompleteness correction. This is similar to the approach used by De Propris et al. (2005) for the MGC and by Masjedi et al. (2006) and Masjedi, Hogg & Blanton (2008). Naturally, this assumes that pairs are missed in a random fashion and potential companions have the same redshift distributions as the overall GAMA survey. This completeness correction amounts to \( \sim 1 \) to 4 per cent between \( M_r = -23 \) and \( M_r = -17 \), respectively, with some increase towards lower luminosities, in agreement with GAMA’s high overall completeness and broad lack of bias at small separations.

An alternative approach is to select galaxies by stellar mass rather than luminosity. For the GAMA data set the galaxy stellar mass was calculated by Taylor et al. (2011) who gives a formula involving velocity dispersion of the disc-galaxy, assuming a single time-scale of 0.5 Gyr for all objects. If they adopted the fitting formulae by Kitzbichler & White (2008), these would result in merger time-scales of 0.8–3.2 Gyr depending on galaxy luminosities, which are somewhat longer than ours. Compared to the more conventional estimates from dynamical friction, the merger time-scales are about 30 per cent longer, but as pointed out by Jiang et al. (2008) these are due to the poor performance of the simple estimate of the Coulomb logarithm. As we give explicit values for the raw pair fraction (with completeness corrections), as well as details on the luminosities and volumes sampled, future studies or theoretical comparisons can adopt any appropriate time-scale and apply it to the data presented here.

Jian et al. (2012) show that the merger time-scale depends sensitively on the mass ratio of mergers, being shorter for equal masses and considerably longer for more minor mergers. Most of our pairs, on the other hand, have similar luminosity ratios (about 4/5 of all objects are <1:4 mergers). The time-scales we calculate above are therefore appropriate for nearly equal mass mergers as in Jian et al. (2012), although our values may be a slight underestimate of the actual merger time-scale. We tabulate the derived volume merger rates in units of \( 10^{-2} \) Mpc\(^{-3}\) Gyr\(^{-1}\) in Table 1. These amount to mean luminosity accretion rates of \( 0.5 \times 10^9 < L/\odot < 5 \times 10^9 \) per Gyr.
for galaxies with $-23.0 < M_r < -17$. If we use the stellar masses computed following Taylor et al. (2011), these luminosity accretion rates correspond to stellar mass accretion rates of $\mathcal{O}(10^9) M_\odot \, \text{Gyr}^{-1}$, equivalent to growth factors of $\approx 5$ per cent of the current mass of each galaxy per Gyr.

We present these quantities in Fig. 2. The pair fraction (top panel) is a nearly constant 3 per cent at all luminosities. Note that the actual merger fraction is 1/2 of these values, if we assume that only 50 per cent of close pairs actually merge. The bottom panel of this figure shows the equivalent for volume merger rates. In both cases, we observe a decline in the merger fraction (and rate) in the lower luminosity bin, and a possible decrease at high luminosities as well; the former is in agreement with observations by Patton & Atfield (2008). These trends are more significant for volume merger rates, because of the different volume densities and merger time-scales for each luminosity bin. We are also able to measure the pair fraction and volume merger rate for galaxies in the brightest luminosity bin within two redshift intervals (centred on $z = 0.09$ and 0.19) to measure the evolution of the merger rate since $z < 0.2$ (approximately 2 Gyr in the past). We do not show the higher redshift data points in Fig. 2 and subsequent figures to avoid confusion. Our measurements are consistent with a flat or even declining merger rate evolution, in agreement with previous observations for similarly massive galaxies (Lin et al. 2004; Kartaltepe et al. 2007; Lin et al. 2008; de Ravel et al. 2009) at higher redshifts, which are also consistent with $R_{\text{inf}}(z) \sim (1 + z)^{-\alpha}$ (Lotz et al. 2011).

### 3.1 Dependence on colour and morphology

We now consider how the merger rate depends on the colour of galaxies. For example, mergers between gas-less galaxies (‘dry’ mergers, with galaxies already on the red sequence) have been proposed as a possible mechanism to grow the red sequence in mass without increasing the intrinsic scatter in the colour–magnitude relation and other scaling factors (Bell et al. 2004; van Dokkum 2005; Faber et al. 2007; Skelton, Bell & Somerville 2012). Obviously, we have no knowledge of the gas content of these galaxies, but we can use the $u - r$ colour to separate quiescent and actively star-forming galaxies, with the dividing line at $u - r = 2.2$ (e.g. Strateva et al. 2001; Baldry et al. 2004). This is acceptable for our purposes, as there is a moderate degree of correlation between colour and gas content, in the sense of red galaxies being more gas poor (Toribio et al. 2011). Here, ‘dry’ mergers are between two ‘red’ galaxies, ‘wet’ mergers between two ‘blue’ galaxies and ‘damp’ mergers contain one member of each.

Another possible definition uses morphology. Galaxy profiles (single Sérsic models) have been derived for GAMA galaxies in Kelvin et al. (2012). We define galaxies with $n > 2.5$ as spheroid dominated and those with $n < 2.5$ as disc dominated. More spheroid-dominated galaxies are known to be more gas poor and we can therefore use this as a proxy of relative ‘dryness’. Here, dry mergers are those between two spheroid-dominated galaxies, while wet mergers are those between two disc-dominated objects, with damp mergers having one member of each class. The derived pair fractions and volume merger rates for dry, wet and damp mergers (selected by colour and morphology, as above) are tabulated in Tables 2 and 3, respectively.

Fig. 3 shows the pair fractions and volume merger rates for galaxies, split into dry, wet and damp mergers. For the brighter luminosity range, we also tabulate separate results for the samples at $(z) = 0.08$ and 0.19 (but note that as in the previous figure, we do not show this data point in Fig. 3 to avoid confusion).

We observe dry mergers to be rare in our most luminous sample at $z < 0.2$, in agreement with previous results by Masjedi et al. (2006, 2008), Wen, Liu & Han (2009), Robaina et al. (2010) and Jiang, Hogg & Blanton (2012). The dry merger rate also decreases monotonically as a function of luminosity. There are very few faint dry mergers. It appears therefore unlikely that the red sequence can be built by this process at least since $z \sim 0.2$ as probed in this study. Most mergers at the faint end are wet and their relative contribution increases with decreasing $M_r$. Damp mergers are intermediate between the behaviour of red and blue pairs. This is true for samples selected by colour or Sérsic index. The intermediate behaviour of damp mergers suggests that galaxies in pairs are selected almost at random from the parent population and therefore that presence in a close pair does not strongly affect the properties of the galaxies. We will examine this in greater detail below. In the brightest luminosity bin, we also find evidence of only flat evolution in the merger rate, irrespective of colour or morphology, since $z \sim 0.2$, suggesting that there has been no change in the makeup of merging pairs (at least for massive galaxies) over the last $\sim 2$ Gyr of the history of the Universe.

### 3.2 Cosmic variance

As GAMA consists of three separate regions, we calculate the pair fractions and volume merger rates for galaxies in each separate region to estimate the cosmic variance. This is shown in Fig. 4 for the full sample as well as for dry, wet and damp mergers (by colour...
The cosmic variance over volumes of $\sim 10^3 h^{-3} \text{Mpc}^3$ is approximately a factor of 2, in reasonable agreement with the estimate by Lopez-Sanjuan et al. (2014) from the ALHAMBRA survey.

### 4 THE PROPERTIES OF GALAXIES IN PAIRS

We now compare the properties of galaxies in pairs (colour, morphology, environmental density and degree of activity, such as star formation or AGNs, measured from the BPT index of Baldwin, Phillips & Terlevich 1981) to those of more isolated (i.e. not in close pairs) galaxies in the volume-limited region from which the pairs are drawn. We divide all galaxies into bins of colour, morphology, environmental density and BPT class and compute the fractions of galaxies in each bin, for objects in close pairs and for their parent sample (more isolated systems) in the same volume. We then calculate the fractional difference (pairs minus parent sample) in each colour (morphology, environment, BPT class) bin. A negative value means that there is a relative deficit of galaxies in close pairs in the appropriate colour (morphology, etc.) bin (i.e. they are less frequent) and vice versa for a positive value. In this way, we can assess how membership in a close pair and the ongoing interaction affects the properties of galaxies (relative to similar objects and morphology). We have applied small shifts in luminosity in this figure for reasons of clarity. The cosmic variance over volumes of $\sim 10^3 h^{-3} \text{Mpc}^3$ is approximately a factor of 2, in reasonable agreement with the estimate by Lopez-Sanjuan et al. (2014) from the ALHAMBRA survey.
Figure 4. Pair fractions (top panels) and volume merger rates (bottom panels) for galaxies in each individual GAMA region. The left-hand panels show all objects, the middle panels plot pairs selected by colour and the right-hand panels pairs selected by their Sérsic index. The size of the bars shows the range of values for the pair fraction (without the error bars) and volume merger rate across all regions and hence the cosmic variance in these quantities. The black bars are for all objects, while red, blue and green bars are for dry, wet and damp mergers, respectively, the same colour scheme as adopted in previous and subsequent figures. We have applied small shifts in luminosity (+0.25 mag for the blue bars and –0.25 mag for the green bars) about the mid-point of each bin (as in Tables 1 and 2) so that the bars do not overlap.

Figure 5. Fractional difference in the distribution of $u - r$ colours for galaxies in pairs and the parent sample. A negative number means that there are fewer pairs, in proportion, in the colour bin considered, with respect to the distribution of the volume-limited box in Fig. 1 and vice versa. The length of each bar shows the 1σ range of the fractional difference at each colour, while the width of the bar shows the colour range we study. The luminosity ranges considered are found in the figure legend.

The colour distribution of galaxies is bimodal, as is well known (Strateva et al. 2001; Baldry et al. 2004), with red galaxies becoming more prominent in the higher luminosity bins. In all cases, except the faintest luminosity range where most galaxies are blue, we note that pairs are overrepresented in the red peak and less frequent in the blue peak (Fig. 5). This argues that presence in a close pair does
not trigger but rather tends to quench star formation. Similar claims of a greater frequency of red galaxies in pairs were earlier reported by Darg et al. (2010a) and Ellison et al. (2010). Chou, Bridge & Abraham (2012) also find that close pairs tend to quench, and not enhance, star formation in their sample. Bergvall, Laurikainen & Aalto (2003) argued that there is only a modest enhancement in star formation rates for a sample of highly interacting objects. In our previous study of pairs in the MGC (De Propris et al. 2005), we
Figure 8. Fractional difference in the distribution of BPT classes for galaxies in pairs and the parent sample. A negative number means that there are fewer pairs, in proportion, in the bin considered, compared to the parent sample of galaxies in the volume-limited boxes of Fig. 1 and vice versa. The length of the bar shows the 1σ range of the fractional difference. The BPT class is indicated in the figure.

noticed a similar excess of red pairs, as well as an excess of blue pairs (and a deficit of green valley systems in pairs), which is not seen here. Darg et al. (2010a) suggest that the apparent increase in the fraction of blue galaxies in pairs in some previous studies may be due to greater visibility of gas-rich, star-forming objects, especially if pairs and mergers are selected by morphology.

The distribution of Sérsic indices shows a deficit of ‘discy’ galaxies among pairs, and a slight excess of more spheroid-dominated objects. This may imply a degree of morphological evolution as part of the merger process. We found a similar trend in MGC data (De Propris et al. 2005) with an excess of E/S0 and Sc/Sd galaxies in pairs compared to the MGC sample, and a deficit of intermediate spirals (see also Darg et al. 2010a). We also concluded (as do Darg et al. 2010a) that this is due to a combination of morphological evolution and induced star formation. Most of these changes (in colour and morphology) appear to take place close to the ‘transition’ colour of $u - r = 2.2$ or Sérsic index of $n = 2.5$. This may indicate that mergers and interactions do not cause abrupt changes in star formation rates or morphologies, but might simply accelerate or enhance an ongoing process of secular evolution (as in the ‘harassment’ scenario by Moore et al. 1996).

However, we cannot exclude the possibility that environment plays a role. Patton et al. (2011) note a similar trend of increasing spheroid fractions among galaxies in pairs and attribute it to the denser environments in which pairs may reside, yielding larger spheroids because of the morphology–density relation. Deng & Zhang (2013) use a volume-limited sample in the SDSS and reach similar conclusions. We therefore look at the environmental density of galaxies in pairs versus other (more isolated) galaxies in the same volume. Here, environmental density is defined using an index developed by Haas, Schaye & Jeeson-Daniel (2012) who use the distance to the seventh nearest neighbour within 1000 km s$^{-1}$ and which is believed to provide a halo-independent measure of the environment of galaxies; this is more appropriate to our purposes (and less model dependent) as we are chiefly interested in exploring whether differences in colour or morphology can be attributed to environmental effects. We derive this index from our GAMA data by applying the recipe of Haas et al. (2012): a shorter distance (i.e. closer seventh neighbour) implies a denser environment and vice versa (higher value of this index implies a less dense environment). Generally, GAMA samples a variety of environments, but most galaxies lie in the field or groups, with only one large cluster known in the GAMA fields.

We show the fractional difference of galaxies (in pairs minus all other objects not in close pairs, as with previous figures) in Fig. 7. This does not support the hypothesis that galaxy pairs reside in more dense environments, in agreement with previous work by Darg et al. (2010b) and Ellison et al. (2010) that finds only a weak environmental dependence in the SDSS pairs (but see Kampczyk et al. 2013, who find a much stronger environmental trend, although this is at $z \sim 0.8$ in the COSMOS field). The weak environmental dependence is somewhat surprising as simulations indicate a much stronger effect of environmental density than here observed (Jian et al. 2012).

Finally, we consider the BPT diagram (Fig. 8) to measure the degree of activity in galaxies in pairs and their parent sample. This classifies all emission-line galaxies as Star Forming, Seyfert, LINER, composite spectra between LINER and Seyfert or between Star Forming and Seyfert (here termed Composite) on the basis of emission-line intensity ratios. We use the classification of Kauffmann et al. (2003) and Tremonti et al. (2004) as applied to SDSS DR8 (and later releases) data by Thomas et al. (2013). This is only a subset of all data in the GAMA survey but it includes about 1/3 of all galaxies in our sample and there is no apparent bias towards or against galaxies in pairs. In agreement with the relative weakness of the blue peak in the colour distribution of Fig. 5, we find a smaller fraction of active star-forming galaxies among pairs, but it is surprising to see that there is little or no evidence of increased AGN activity among galaxies in pairs compared to the volume-limited parent sample of objects. In all magnitude ranges we consider, there are fewer
Seyfert or Composite (LINER+Seyfert or Star Forming+Seyfert) objects in the pairs sample, while LINERs are only marginally more common.

5 DISCUSSION

We have measured pair fractions and merger rates for galaxies in the GAMA survey, with $-23 < M_r < 5\log h < -17$ in a series of volume-limited samples and explored how the merger rate varies with luminosity, colour or morphology. We then considered how the properties of galaxies in close pairs (colour, morphology, environmental density and BPT class) compare to those of more isolated objects in the same volume to understand how galaxies are affected by mergers and interactions.

The merger fractions for galaxies in the GAMA survey are (assuming that 50 per cent of mergers are dynamically bound) $\sim 1.5$ per cent almost irrespective of luminosity. These are generally in good agreement with previous estimates of the pair fraction and merger rates in the local universe, as well as with the well-known flat evolution out to high redshift, as in the compilation by Lotz et al. (2011). Our previous local ($z < 0.1$) estimate of the merger rate from the MGC (De Propris et al. 2005, 2007) is in good agreement. Comparison of the actual volume merger rates requires understanding assumptions concerning the space density of galaxies and BPT class) compare to those of more isolated objects in the same volume to understand how galaxies are affected by mergers and interactions.

Finally, the lower redshift points of Kartaltepe et al. (2007) and Lin et al. (2008) are reasonably close to our values, despite the somewhat different selection criteria in each case.

The volume merger rates are seen to decrease significantly at both high and low luminosities in our sample. Given the measured accretion rates and evolution of the merger rate, we estimate that the total stellar mass growth (given the luminosity accretion rates and reasonable mass-to-light ratios) of galaxies in the past 1/2 of the Hubble time due to mergers lies in the range of 10–30 per cent. This assumes merger time-scales as discussed above, a 50 per cent ‘true’ merger fraction and the flat evolution of the merger rate found in this work to $z \sim 0.2$ and measured to $z = 1.2$ by other studies (e.g. Lin et al. 2008; de Ravel et al. 2009). Lopez-Sanjuan et al. (2012) estimate a similar $\approx 30$ per cent for the growth rate of galaxies since $z \sim 1$. However, these trends are more massive than the GAMA sample which is generally between $9 < \log M_*/M_\odot < 11$.

The trend of merger rates with luminosity and redshift is instead not in good agreement with predictions from dark matter simulations such as Murali et al. (2002) and Maller et al. (2006). This is significantly lower than estimates of the mass growth of galactic haloes in simulations (e.g. van den Bosch 2002; Cattaneo et al. 2011), which on average double the mass of galaxies since $z \sim 1$. The discrepancy may be reduced if the merger rate evolution with redshift is faster but our data for the brighter galaxies are consistent with a flat or declining evolution since $z = 0.2$, irrespective of colour or morphological type, as previously observed for similarly luminous galaxies at higher redshifts by Bundy et al. (2004), Lin et al. (2004, 2008) and de Ravel et al. (2009). In the compilation of Lotz et al. (2011), the merger rate evolves as $(1 + z)^{-0.1}$ for luminosity (or mass) selected samples, in agreement with very slow growth rates by major merging for massive galaxies. Robotham et al. (2014) also find that the evolution of the merger rate is close to flat.

Our data allow us to consider the local dry, wet and damp merger rate. Dry mergers (between either red or spheroid-dominated galaxies) are observed to be rare in the local universe, in agreement with several previous studies (Masjedi et al. 2006, 2008; Wen et al. 2009; Robaina et al. 2010; Jiang et al. 2012), while we also observe that dry mergers are less frequent as a function of decreasing luminosity. We find evidence of slow evolution for at least the more massive (luminous) systems. In our close pair study of luminous red galaxies at $(z) \sim 0.55$, we measured a stringent upper limit of $< 0.8$ per cent per Gyr, which is consistent with flat evolution of the dry merger rate out to at least $z = 0.6$. Chou et al. (2012) also confirm the slow growth of the dry merger rate as a function of redshift out to $z = 0.1$. It is therefore unlikely that the red sequence is formed by dry mergers at least in the past 1/3 of the Hubble time. Most mergers between lower luminosity galaxies are instead wet (and the rate increases towards lower luminosities) while damp mergers are intermediate.

When we compare the colour distribution of galaxies in pairs with that of galaxies in their volume-limited parent sample, we observe a deficit of blue objects and an excess of red galaxies. This argues for suppression of star formation in galaxy pairs, as noted earlier by Darg et al. (2010a) and Ellison et al. (2010). This is somewhat surprising in the light of theoretical expectations that mergers and interactions boost the star formation and activity rates (e.g. Mihos & Hernquist 1996). We confirm this from our analysis of BPT classes where there is a lower fraction of star-forming galaxies among close pairs compared to galaxies in the same volume. Bergvall et al. (2003) found only a modest degree of enhancement of star formation for interacting pairs, contributing little to the luminosity; Lin et al. (2007) also find only a modest enhancement of star formation in galaxy pairs in a wide range of merger stages; Li et al. (2008) instead find that the enhancement in star formation, in a sample of star-forming galaxies, increases with decreasing separation but does not depend on the luminosity ratio, although it is stronger for lower mass galaxies; for massive galaxies, Robaina et al. (2009) find that only 10 per cent of star formation is due to major interactions (albeit at higher redshifts than studied here); Patton et al. (2011) also detect an excess of red galaxies in pairs and attribute this to denser environments and hence more significant spheroids (because of the morphology–density relation); Chou et al. (2012) confirm that most close pairs are red and not blue; Deng & Zhang (2013) also find evidence that interactions are not generally responsible for star formation bursts in two volume-limited samples drawn from the SDSS; Cluver et al. (2014) study star formation for galaxies in groups and claim that galaxies with a close neighbour have lower star formation rates. Similarly, we find that, by comparing the distributions of S´ersic indices for galaxies in pairs and their parent sample, there is a relative lack of ‘discs’ and excess of ‘spheroids’ among paired galaxies as in Patton et al. (2011). Spheroids are generally associated with suppression of star formation, but it is not clear whether this is a cause or an effect (McGee et al. 2011). If pairs lie in more dense environments, they will contain more prominent spheroids (by the standard morphology–density relation) and therefore have lower star formation rates (as suggested by Patton et al. 2011), explaining the excess of red galaxies among pairs found above. However, in our sample, we do not find that galaxies in pairs lie in significantly different environments (as do Darg et al. 2010b and Ellison et al. 2010 who find only a mild environmental
dependence). This suggests a degree of morphological evolution is also associated with star formation suppression in these objects. Li et al. (2008) show that tidal interactions may indeed cause an increase in galaxy concentration, which is an indicator of more significant spheroids. However, this is not due to the fact that the galaxies lie in more dense environments and are more spheroid dominated (and therefore more quiescent). Rather, it appears that the extra spheroid light is related to the suppression of star formation and presence in a close pair, although it is not possible to determine from the present sample whether quenching or morphological evolution occur first (Li et al. 2008).

One mechanism by which we may explain these findings is ‘harassment’ (Moore et al. 1996). Galaxies in these close pairs are likely to have undergone a series of previous encounters and close passes; Patton et al. (2013) argue that about 2/3 of the star formation in close pairs takes place at separations above 30 kpc. This leads to a gradual suppression of star formation as gas is exhausted or driven out, while dynamical relaxation leads to secular evolution of spheroids without mergers. Chou et al. (2012) also suggest that in their sample star formation is triggered at large radii and then suppressed to form red close pairs. Robotham et al. (2013) point out the role of mergers and interactions in suppressing star formation in most galaxy pairs, while Trihun et al. (2013) present an analysis of the properties of single and widely spaced paired galaxies in the SDSS and find evidence of delayed quenching, over time-scales of a few Gyr, consistent with a ‘harassment’ scenario as proposed here. On the other hand, an Integral Field Unit study of the Mice pair, shows that little star formation has been induced so far (Wild et al. 2014).

The BPT classifications, from emission-line diagnostics, support the observation that galaxies in pairs are less active than their counterparts in the same volume. As remarked above, there is a deficit of star-forming galaxies in pairs. Surprisingly, there is also no evidence of excess AGN activity in galaxies in pairs. This is surprising as one expects mergers to trigger nuclear activity (Mihos & Hernquist 1996; Springel, di Matteo & Hernquist 2005; Hopkins et al. 2008). In these models, AGNs occur at or shortly after final coalescence so we should expect double QSOs in close pairs (Comerford et al. 2009). Bessiere et al. (2012), however, find QSOs at several stages of the merging process even before the final merger. On the other hand, Cisternas et al. (2011) show that AGNs do not lie in more disturbed hosts and conclude that mergers are not necessarily responsible for nuclear activity. Teng et al. (2012) search for double QSOs in pairs of galaxies where one member is already known to host an AGN and find only one such case out of 12 objects. This suggests that nuclear activity is not necessarily related to the ongoing mergers or interactions in these objects. Ellison et al. (2008) also report on a relative lack of AGN activity in close pairs from the SDSS although Patton et al. (2011) claim to find some enhancement from the same sample. Deng, Yu & Jiang (2013) also argue that there is less than 1σ significance for AGN triggering in mergers from a volume-limited sample of SDSS galaxies. We find no evidence of AGN activity in strongly interacting mergers and merger remnant post-starburst galaxies in De Propris & Melnick (2014), although there are indications of centrally concentrated star formation and inside-out quenching in these objects. Other studies find a stronger correlation between presence in a close pair and AGN activity (Patton et al. 2011). In some cases, these discrepancies may be due to issues of selection, including the use of emission-line diagnostics rather than X-ray or mid-IR fluxes.

At face value our findings minimize the importance of major mergers and interactions for galaxy evolution and argue that most galaxy evolution takes place via internal and secular processes, as well as minor mergers, at least at low redshift. This is also the conclusion of several studies in this field (e.g. Lopez-Sanjuan et al. 2012; Huang et al. 2013; Kaviraj et al. 2013; McLure et al. 2013; Fritz et al. 2014) and is broadly in agreement with more recent models where minor mergers and secular evolution may play more important roles in galaxy evolution in the past 1/2 of the Hubble time (e.g. Guo & White 2008; Parry, Eke & Frenk 2009).

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