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A Multiwavelength Dynamical State Analysis of ACT-CL J0019.6+0336

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Abstract: In our study, we show a multiwavelength view of ACT-CL J0019.6+0336 (which hosts a radio halo), to investigate the cluster dynamics, morphology, and ICM. We use a combination of XMM-Newton images, Dark Energy Survey (DES) imaging and photometry, SDSS spectroscopic information, and 1.16 GHz MeerKAT data to study the cluster properties. Various X-ray and optical morphology parameters are calculated to investigate the level of disturbance. We find disturbances in two X-ray parameters and the optical density map shows elongated and axisymmetric structures with the main cluster component southeast of the cluster centre and another component northwest of the cluster centre. We also find a BCG offset of \( \sim 950 \text{ km/s} \) from the mean velocity of the cluster, and a discrepancy between the SZ mass, X-ray mass, and dynamical mass (\( M_{X,500} \) and \( M_{SZ,500} \) lies \( > 3\sigma \) away from \( M_{\text{dyn,500}} \)), showing that J0019 is a merging cluster and probably in a post-merging phase.

Keywords: diffuse radio emission; galaxy clusters; multiwavelength; radio halo; cluster dynamical state; morphology; ICM; turbulence; bent tailed galaxies

1. Introduction

Multiwavelength observations of galaxy clusters carry an abundance of information, giving insight into the intracluster medium (ICM), its thermal and non-thermal components, and the cluster dynamics. Galaxy clusters merge with neighbouring clusters through the most energetic events in the Universe since the big bang [1]. The merger disturbs the cluster’s natural state through merger-driven turbulence and often leave imprints on the ICM.

Turbulence in the ICM also affects individual sources in the cluster, particularly radio galaxies [2]. There are differences in radio emission from cluster galaxies and isolated galaxies that are generally attributed to the interaction of the emission with the ICM. Most
radio emissions from galaxies have a compact radio source associated with the active
galactic nuclei (AGN) and extended regions of radio emission (radio lobes) quite distant
from the compact radio source. The morphology of extended radio emission associated with
individual galaxies (such as bent tail galaxies) is strongly influenced by the environment in
which the galaxy exists [3]. The motion of galaxy clusters through the dense cluster gas
is widely recognized as the mechanism which produces bent radio source tails and the
dynamic pressure responsible for a variety of other observed source shapes [3].

A fraction of the gravitational energy released during violent merger events is con-
verted into magnetic-field amplification and results in the acceleration of high energy
particles in the ICM, giving rise to diffuse radio sources [4]. Cluster-scale diffuse radio
sources have steep spectra, exhibit low surface brightness, and are broadly categorized into
three groups: radio halos, radio mini halos, and radio relics. Radio halos and mini halos
are centrally located sources, whereas relics are found as filamentary, elongated shapes
located at the cluster periphery (see [3] for a recent review).

Diffuse radio emission and multiwavelength studies provide more knowledge about
the link between the underlying thermal and non-thermal processes present in the ICM [6].
Several open questions regarding cluster diffuse emission exist. It is not understood why
some merging clusters host radio halos and relics while others do not show any evidence
of extended radio emission [7]. The correlation between non-thermal diffuse emission
and the dynamical properties of the host clusters is unclear and the underlying formation
mechanisms are not well understood [7]. These open questions have larger implications
on topics and processes which include dark matter, AGN feedback mechanisms, and the
large-scale structure of the universe. The sensitivity of new generation radio telescopes will
enable studies of large and diverse samples of clusters at higher redshift and lower mass.

Previous studies of radio diffuse emission were limited to high mass and low redshift
samples, and clusters with high redshifts were generally solitary detections [8,9]. Large
sample studies, such as Cuciti et al. [10], are crucial to a complete understanding of the
evolutionary life of radio diffuse emission, formation theories of radio sources, and
analysis of radio source statistics. The MeerKAT Exploration of Relics, Giant Halos and
Extragalactic Radio Sources (MERGHERS; [11]) project aims to be a mass-selected sample
of ∼200 Sunyaev–Zel’dovich (SZ) selected clusters. These clusters will cover a suitable
range of criteria to ensure an unbiased study. In preparation for such a large project, a pilot
sample of 13 clusters has been obtained using the MeerKAT telescope [12].

In this paper, we perform a multiwavelength follow-up of ACT-CL J0019.6+0336
(hereafter J0019), presented in Knowles et al. [12], to further understand the dynamics of the
cluster. The paper is structured as follows. In Section 2, we show existing multiwavelength
data on J0019. Section 3 shows the morphology analysis methods and results, followed by
the discussion and conclusion in Sections 4 and 5 respectively. In this paper, we adopt a
ΛCDM flat cosmology with $H_0 = 70$ km s$^{-1}$ Mpc$^{-1}$, $\Omega_\Lambda = 0.7$, $\Omega_m = 0.3$.

2. Observations and Data

J0019 lies at $z = 0.266$ and was first detected in a catalogue of galaxy clusters derived
by Zwicky et al. [13] (ZwCl 0017.0+0320) and was later observed by the Atacama Cos-
mology Telescope (ACT; [14]). In this section, we describe the multi-wavelength data and
MeerKAT observations used for this study. The relevant cluster properties are given in
Table 1.
Table 1. Published properties of J0019. The position, redshift, SZ mass, and radio properties are from Knowles et al. [12], with the mass being the weak lensing-calibrated value. Integrated 0.5–2.0 keV X-ray luminosity is calculated within $R_{500} = 1389.5$ kpc (spherical radius that encloses an average density equal to 500 times the critical density at the cluster redshift, $z = 0.266$).

<table>
<thead>
<tr>
<th>Property</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>R.A. (J2000)</td>
<td>4.91085 deg</td>
</tr>
<tr>
<td>Dec. (J2000)</td>
<td>3.60879 deg</td>
</tr>
<tr>
<td>Redshift</td>
<td>0.266</td>
</tr>
<tr>
<td>$M_{500c,\text{SZ}} (10^{14} M_\odot)$</td>
<td>$10.2 \pm 2.28$</td>
</tr>
<tr>
<td>$L_{500,\text{X}} (10^{44} \text{ ergs/s})$</td>
<td>$4.94 \pm 0.032$</td>
</tr>
<tr>
<td>$S_{1.16\text{GHz}} \text{ (mJy)}$</td>
<td>$9.16 \pm 0.57$</td>
</tr>
<tr>
<td>$\log P_{1.4\text{GHz}} \text{ (W/Hz)}$</td>
<td>$24.13 \pm 0.06$</td>
</tr>
</tbody>
</table>

2.1. Millimetre

Studies such as Basu [15] and Cuciti et al. [16], have shown that samples selected via their SZ signal show a higher detection rate than X-ray selected samples and SZ selected samples have the benefit of the flux limit translating directly into a mass limit [17]. Various SZ telescopes now provide large SZ-detected cluster samples in which galaxy clusters are detected via the distortion of the cosmic microwave background (through inverse Compton scattering of high-energy electrons). J0019 was observed in both Planck and ACT surveys [18,19]. We make use of the ACT observations because they have higher sensitivity and resolution than the Planck observations.

J0019 is part of the fifth release of the ACT cluster catalogue ([20], ACT DR5) which contains more than 4000 SZ selected optically confirmed clusters [18]. J0019 was observed as part of the larger SZ-selected sample from ACT that forms the targets for MERGHERS with a signal-to-noise ratio, $S/N = 25.23$ (at the reference 2.4′ filter scale) [18]. This shows that J0019 is a very massive and energetic cluster.

2.2. Optical

We make use of available Dark Energy Survey (DES) DR1 photometry for J0019. ZCLUSTER$^3$ estimates the photometric redshift of a particular galaxy cluster using multi-band optical and infrared photometry. The optical photometry is obtained using the ZCLUSTER code and the photometric redshift of the cluster is estimated by the weighted sum of individual galaxy probability distributions in the direction of each galaxy [14]. These probability distributions are found using a template fitting method as seen in Benitez [21], where a set of default galaxy spectral energy distribution (SED) templates [18] are fitted to the observed broadband SED of each galaxy. Moreover, spectroscopic observations are available for 17 possible cluster members (within $R_{500}$) in the Sloan Digital Sky Survey (SDSS) DR16 database [22].

2.3. Radio

J0019 has been observed by the MeerKAT telescope for $\sim 24$ min at L-band ($\sim 0.9–1.7$ GHz) ([12], PI:Knowles). First, second, and third-generation MeerKAT calibration algorithms have been applied to the radio data to correct for direction independent and dependent effects. First and second-generation calibration was applied to the field using the oxkat v1.0$^2$ pipeline [23], a Python-based reduction pipeline for MeerKAT. Third-generation calibration has been applied to the field using DDFACET$^3$ and Killms$^4$ to correct for corrupting artifacts in the radio image. The KATBEAM$^3$ package is used to create a primary beam corrected final image for analysis. A large scale filtered image (contours overlaid in Figure 1) is also created through an image-plane filtering technique introduced by Rudnick [24], in which extended emission without compact sources is highlighted by filtering out emission on a scale 1–3 times the synthesised beam. The left panel of Figure 1 shows the final full-resolution radio map of the J0019 cluster region with large-scale filtered image contours overlaid. A radio halo ($\sim 810$ kpc in size) has been detected.
with a 1.16 GHz flux density of $9.16 \pm 0.57$ mJy. The power value has been estimated as $P_{1.4 \text{GHz}} = 24.13 \pm 0.06$ W/Hz by using a fiducial spectral index, $\alpha = -1.3 \pm 0.4$ and extrapolating the MeerKAT flux densities to 1.4 GHz (adopting the spectral power law convention of $S_\nu \propto \nu^\alpha$).

**Figure 1.** Left: Full-resolution MeerKAT L-band radio image for J0019 with large scale filtered image contours overlaid. Contours are at $[-3, 3, 5, 10] \times \sigma$ (52 $\mu$Jy/beam) of the filtered image. The synthesized beam for both the full-resolution (red ellipse, 7.9$''$ $\times$ 7.2$''$, p.a. 162$^\circ$) and filtered map (black circle, 21$''$) are indicated in the lower left of the panel. The dashed red circle denotes $R_{500,pc}$ centred on the ACT SZ peak, which is marked by a black cross. Full resolution central rms noise is 18.1 $\mu$Jy/beam. Right: XMM-Newton combined PN, MOS1, and MOS2 images of J0019 with the large-scale radio contours overlaid. The black cross represents the ACT SZ peak and the green Xs are centred on the peak radio emission of sources found in the MeerKAT radio image. Regions of the image masked after removing X-ray point sources are indicated by yellow, dashed, excluded circles. The image is unbinned and has been smoothed by a Gaussian with a kernel radius of 6 pixels (1 px = 4.35$''$).

2.4. X-ray

We use archival data from the XMM-Newton space telescope to provide an X-ray view of J0019 and probe its morphology. The only XMM data available for J0019 was obs ID: 0693010301, a pointed, 47.9 ks observation of the cluster taken in 2012. Events from all three EPIC (PN, MOS1, and MOS2) cameras have been used. The XMM Cluster Survey ([25], XCS) cleaned the data and performed source detection with a custom version of WAVDETECT [26] (called the XCS Automated Pipeline Algorithm (XAPA)). XAPA is run on merged EPIC (PN+MOS1+MOS2) images, with a pixel size of 4.35$''$, that have been generated using events within an energy range of 0.5–2.0 keV. Once sources in an image have been located, XAPA classifies them as either point or extended.

We use XMM: Generate and Analyse (XGA; [27]), a new open-source Python module created for XCS, to generate soft-band (0.5–2.0 keV) count-rate maps. XGA then uses the XCS source catalogue to produce a mask to remove irrelevant sources. It identifies the XCS source that matches J0019, then masks all other sources. The final processed X-ray image is shown in the right panel of Figure 1.

The unabsorbed X-ray luminosity in the 0.5–2.0 keV energy band was measured by performing simultaneous fits to spectra generated within $R_{500,pc}$ of J0019, centred on the Knowles et al. [12] coordinates. We fit absorbed (with $tbabs$, [28]) plasma emission models (APEC, [29]) to the three spectra, with temperatures linked, metallicity fixed at 0.3 $Z_\odot$, redshift fixed at 0.266, and nH fixed at 0.029 $\times$ 10$^{22}$ cm$^{-2}$ (taken from the full-sky HI survey by the [30]). The fit is performed using XGA’s XSPEC [31] interface, SAS v17.0.0, and XSPEC v12.10.1. We measure a temperature of $T = 6.7 \pm 0.1$ keV within $R_{500}$ and a bolometric X-ray luminosity of $L_X = (1.86 \pm 0.02) \times 10^{45}$ ergs/s. Comparing to the $M_X-T_X$ relation in [32], we estimate a mass within $R_{500}$, $M_X \sim (6.5 \pm 1.3) \times 10^{14} M_\odot$. 
3. Cluster Morphology

Galaxy clusters form via a hierarchical sequence of accretion and mergers of smaller interacting substructures [6]. Current observations of radio diffuse emission favour theories of merger-driven radio halo and radio relic formations [5]. Therefore, it is important to understand the dynamical state of J0019. Numerous studies, e.g., Böhringer et al. [33], Wen and Han [34] have analyzed the substructure of clusters to determine their dynamical states using multiwavelength data. By studying the form or shape of the cluster we can infer the dynamics. We use available multiwavelength information to estimate a variety of X-ray and optical-derived parameters.

3.1. Optical Morphology

ZCLUSTER retrieves and uses redshift probability distributions to measure a galaxy cluster’s photometric redshift. Thereafter, we create a projected 2-D density map of the galaxy cluster using the right ascension (in degrees), declination (in degrees), photometric redshift of the cluster, and a catalogue of redshift probability distributions of individual galaxies. We use the right ascension and declination of each galaxy to position them onto a 2-D map as done in Wen and Han [34]. The galaxy coordinates are converted into Cartesian coordinates and centred at the input R.A. and Dec. (SZ peak), on a projected 4 Mpc × 4 Mpc map. The map is made up of 20 × 20 bins, each with 0.2 Mpc length. The positions of the galaxies in Cartesian coordinates are obtained by finding the angular distance between the coordinates of individual galaxies and the centre of the cluster.

The peak of the probability distribution gives the maximum likelihood of a galaxy having a specific redshift [14]. We obtain a weighting for our plot by integrating the probability distribution of each galaxy around the photometric redshift (bounds of integration from \( z_{\text{phot}} - \Delta z \) to \( z_{\text{phot}} + \Delta z \)). The ZCLUSTER code takes the above-calculated quantities and produces a 2-D projected density map of the galaxy cluster. Thereafter, the map is convolved with a Gaussian kernel (with smoothing length = 2 pixels) to smooth the map.

A density uncertainty map (right panel of Figure 2) is created using Monte Carlo simulations. We add Gaussian-distributed noise to the galaxy photometry at each Monte Carlo step and recompute the photometric redshifts, probability distributions, and density maps. After 1000 iterations, the 68th percentile of the maps is taken as an estimate of the error map.

![Figure 2. Left: Density map of J0019 obtained from ZFIELD using photometry from the DES DR1 database and \( z = 0.266 \). The projected density map is 4 Mpc × 4 Mpc with each pixel 0.2 Mpc × 0.2 Mpc in size. Black large-scale radio contours are overlaid and the red cross indicates the SZ peak of the cluster. The units of the colour scale are counts per pixel. Right: Corresponding error map for the J0019 galaxy density map.](image-url)
The density map for J0019 (left panel of Figure 2) is made using zField (a subclass of zCluster) in which we have used photometric redshift, \( z = 0.266 \), and a redshift range, \( \Delta z = 0.2 \). The weighting for this density map is obtained by integrating the probability distribution for each galaxy around \( z = 0.266 \) (bounds of integration from 0.266 − 0.2 to 0.266 + 0.2).

3.1.1. Centre Shift

The centre shift method used in this work is a simple method to determine the dynamical state of a cluster by calculating the offset between the SZ peak and the peak in the optical density map. The centre of the cluster generally hosts the brightest cluster galaxy (BCG). In relaxed clusters, the centre of galaxy clusters traced by the gas using X-ray, or the SZ effect, coincides (within a few kpc) with the centre of the galaxy cluster traced optically. This can be contrasted with the centres of mergers (determined using X-ray and SZ data), which can be hundreds of kpc away from the centroid determined optically as seen in Pandge et al. [35].

We measure the distance between the coordinates of the maximum of the density plot and the coordinates of the cluster SZ peak [36]. After finding the coordinates of the peak value of the density plot (left panel of Figure 2), we use the distance Formula (1) to find the centre shift.

\[
CS = \sqrt{(x_m - x_0)^2 + (y_m - y_0)^2 \times 0.2}, \quad (1)
\]

where \( x_m \) and \( y_m \) are the x and y coordinates of the peak of the density plot and \( x_0 \) and \( y_0 \) are the coordinates of the centre of the plot (SZ peak). We multiply the distance formula by the size of 1 pixel in Mpc (0.2), to obtain CS in the projected Mpc.

3.1.2. Asymmetry Parameter

Another way to identify the morphology and dynamical state of a galaxy cluster is to study and classify the asymmetry of the cluster. The morphology of astronomical systems is closely related to their symmetry. If a cluster is asymmetric, then it has been disturbed in some way and formed substructures [34]. An efficient way to estimate the asymmetry of a cluster is to flip the density map horizontally and vertically and then compare it to the original image [37]. The asymmetry parameter is calculated using the following formula [37]:

\[
A^2 = \frac{\sum (I_0 - I_\phi)^2}{\sum (2I_0^2)}, \quad (2)
\]

where \( I_0 \) is the original image and \( I_\phi \) is the image that has been flipped horizontally and vertically.

3.2. X-ray Morphology

In this section, we show the calculations for the three main X-ray morphological parameters to further study the cluster dynamics and compare our results with those available in the literature. Following the work of Santos et al. [38], Poole et al. [39], and Buote and Tsai [40], we calculate the power ratio, centroid shift value, and concentration parameter of J0019.

3.2.1. Concentration Parameter

Relaxed clusters generally host luminous cool cores in their centres whereas merging clusters show core disturbance [41]. The concentration parameter is defined as the ratio of the X-ray integrated fluxes in the cluster core and the larger-scale region of the galaxy cluster [38]. The concentration parameter \( c_{SB} \), following Santos et al. [38], is calculated in two circular regions; a core radius of 100 kpc and an outer radius of 500 kpc using:

\[
c_{SB} = \frac{S < (100 \text{ kpc})}{S < (500 \text{ kpc})}, \quad (3)
\]
where $S$ is the X-ray surface brightness centred on the X-ray peak, within a particular radius (i.e., 100 kpc and 500 kpc). Cassano et al. [7] classifies a cluster as relaxed if $c_{SB} > 0.2$ and disturbed if $c_{SB} < 0.2$.

### 3.2.2. Centroid Shift

The centroid shift method is the most robust X-ray morphology estimator, as it is most sensitive to the dynamical state while being least sensitive to image quality [39]. The centroid shift is defined as the distance between the X-ray peak and the centre of mass, in units of aperture radius, $R_{ap}$, and is calculated using the following equation:

$$ w = \frac{1}{N-1} \left[ \sum \left( \Delta_i - \langle \Delta \rangle \right)^2 \right]^{1/2} \times \frac{1}{R_{ap}} $$

where $\Delta_i$ is the distance between the X-ray peak and the centroid of the $i$-th aperture and $\Delta$ is the mean of the sample of $\Delta_i$s.

This equation calculates the offset ($\Delta_i$) between the X-ray peak and the centroid using a series of concentric circles. The centroid of the $i$-th aperture is the centre of mass within the $i$-th aperture. A sample of $\Delta_i$s are computed within a circular aperture of radius $R_{ap}$, ranging from $R_{ap} = 500$ kpc to $0.05 \times R_{ap}$, decreasing in steps of 5%, as seen in Poole et al. [39]. Weißmann et al. [42] classifies a cluster as relaxed if $w < 0.01$ and disturbed if $w > 0.01$.

### 3.2.3. Power Ratios

The power ratio probes the underlying mass distribution to analyze the substructure of a cluster. It is calculated using a multipole decomposition of the potential of the two-dimensional projected mass distribution within a certain aperture [40]. Instead of mass, it is applied to the X-ray surface brightness. The third-order power ratio is one of the best indicators of the dynamical state of a cluster [43]. The general $m$-order power ratio (where $m > 0$) is defined as:

$$ P_m = \frac{1}{2m^2R_{ap}^2} (a_m^2 + b_m^2) $$

and

$$ P_0 = a_0 \ln(R_{ap})^2 $$

where $a_0$ is the total intensity within aperture $R_{ap}$. $a_m$ and $b_m$ are generic moments in polar coordinates, $R$ and $\phi$, given by:

$$ a_m(r) = \int_{R' \leq R_{ap}} S(x')(R')^m \cos(m\phi') d^2x' $$

$$ b_m(r) = \int_{R' \leq R_{ap}} S(x')(R')^m \sin(m\phi') d^2x' $$

The $P_2$ quadrupole power represents the cluster ellipticity, $P_3$ represents the bimodal distribution, and $P_4$ is similar to $P_2$, but has power on much smaller scales. $P_3$ is most suitable to identify the presence of substructures or asymmetries [43]. We make use of the normalized hexapole moment, $P_3 / P_0$, to provide a clear measure of substructure calculated in an aperture of radius, $R_{ap} = 500$ kpc, centred on the X-ray cluster centroid as seen in Buote and Tsai [40]. Based on previous studies [42], a cluster is disturbed if $P_3 / P_0 > 10^{-7}$.

### 3.3. Results and Comparison with the Literature

To compare our optical results to the literature, the above methods were applied to a sample of 98 galaxy clusters found in Wen and Han [34], with known cluster dynamical states. With the set of results obtained from this sample, we can categorize clusters as relaxed or unrelaxed based on the criteria shown in Figure 3.
We find that galaxy clusters with an asymmetry value, $A^2 > 0.1$ or a centre shift value, $CS > 0.4$ are all dynamically disturbed. Galaxy clusters with centre shift values, $CS < 0.4$, and asymmetry values $< 0.1$ could be classified as either relaxed or disturbed.

Applying the methods from Sections 3.1.1 and 3.1.2 on the density map of J0019 (3), we obtain a centre shift value, $CS = 0.20 \pm 0.16$ Mpc, and an asymmetry value, $A^2 = 0.068 \pm 4.9 \times 10^{-4}$.

The statistical uncertainties on the above parameters are estimated by Monte Carlo simulations in which we add Gaussian distributed noise to the galaxy photometry. The density map is recreated at each Monte Carlo step and after 1000 iterations, the standard deviation of the sample is taken as an estimate of the uncertainty.

The X-ray image has been smoothed with a Gaussian kernel of smoothing length 4 pixels for calculations. The following results were obtained: $P_3/P_0 = (8.93 \pm 0.4) \times 10^{-8}$, $w = 0.03 \pm 0.01$, $c_{SB} = 0.13 \pm 0.002$. The analysis methods in Sections 3.1.1 and 3.1.2 have previously [42] been applied to a sample of 80 galaxy clusters observed with the XMM-Newton telescope and have classified a cluster as disturbed when $P_3/P_0 > 10^{-7}$ and $w > 0.01$. Another study [7] also characterized cluster substructures based on the power ratios, centroid shift, and X-ray brightness concentration parameter estimated from Chandra X-ray images. They define a cluster to be dynamically disturbed if its morphological parameters satisfy the following conditions: $P_3/P_0 > 1.2 \times 10^{-7}$, $c_{SB} < 0.2$ and $w > 0.012$. Two of the estimated morphological parameters ($c_{SB}$ and $w$) satisfy the above conditions, classifying J0019 as a dynamically disturbed cluster with $P_3/P_0$ falling $\sim 3\sigma$ away from the threshold value of $10^{-7}$.

The statistical uncertainties on the morphological parameters are estimated from a set of 1000 Monte Carlo simulations in which Poisson/shot noise is added to the observational
XMM X-ray data. Uncertainties are estimated from the standard deviation of a sample of 1000 simulated images with added noise.

To verify our calculations of the three morphological parameters, we use XMM-Newton data on a known merging cluster, ACT-CL J0528.8−3927, and compare the derived morphological parameters to the known literature values found in Weißmann et al. [42] and Lovisari et al. [44]. All estimated values obtained by our code lie within 1σ of the literature values.

3.4. Optical Redshift Distribution

The redshift distribution of 17 spectroscopically confirmed SDSS DR16 [22] cluster member galaxies are used to gauge any disturbance in the cluster. This distribution is shown in Figure 4 where there is an indication of bimodal structure in the histogram.

![Figure 4](image)

Figure 4. Histogram showing the redshift distribution for 17 spectroscopically confirmed cluster members within $R_{500}$ from SDSS, DR16. Here, $v = 0$ is defined as the cluster systemic redshift of $z = 0.266$. The bin width is 350 km/s. The vertical black dashed line shows the velocity of the BCG of the main component.

The standard deviation of the redshift distribution corresponds to the velocity dispersion, $\sigma_v$. Following Munari et al. [45], we use the calculated velocity dispersion and the galaxies-based scaling relation to determine $M_{200}$ and $R_{200}$, using a value of $h = 0.7$ in the Munari et al. [45] equation. For comparison with other results, we convert our $M_{200}$ estimates into $M_{500}$, following the appropriate relation given in Duffy et al. [46]. The quantities $M_{500}$ and $R_{500}$ have been rescaled from $M_{200}$ and $R_{200}$ assuming a concentration parameter, $c = 3$. We integrate a NFW profile Navarro et al. [47] and interpolate to determine $M_{500}$ and $R_{500}$. We estimate uncertainties on all cluster properties by bootstrap resampling 1000 times.

From the mean redshifts of the components, we find a mean velocity at $z = 0.265$. The velocity dispersion, $\sigma_v$, is $810 \pm 160$ km/s. We obtain a dynamical mass of $M_{200} = (4.1 \pm 1.8) \times 10^{14} M_\odot$ and $M_{500} = 2.5 \pm 1.2 \times 10^{14} M_\odot$.

4. Discussion

It is believed that diffuse radio sources are associated with dynamically disturbed clusters. We can conclude that J0019 is undergoing a merger event through the detection of the radio halo shown in [12]. A radio halo in J0019 would require a significantly substantial merger to power it. Previous studies [42] have classified a cluster as disturbed when $P_3/P_0 > 10^{-7}$ and $w > 0.01$. The authors of [7] also characterized cluster substructures
based on the power ratios, centroid shift, and X-ray brightness concentration parameter estimated from Chandra X-ray images. They define a cluster to be dynamically disturbed if its morphological parameters satisfy the following conditions: $P_3/P_0 > 1.2 \times 10^{-7}$, $c_{SB} < 0.2$, and $w > 0.012$. The morphological estimators for J0019 (X-ray concentration and X-ray centre shift values), classify J0019 as a merging cluster and the X-ray $P_3/P_0$, optical centre shift, and optical asymmetry values are inconclusive. X-ray morphology parameters (particularly the power ratios) are largely insensitive to substructures along the line of sight. Spectroscopy of more cluster members are needed to gauge any disturbed morphology in this direction.

The optical density map shows a displacement (0.2 Mpc) between the ACT cluster SZ peak and the peak in the density map. Moreover, the density map is fairly elongated and axisymmetric. The density map shows two components, one along a northwestern direction (top) and another along a southeastern direction (bottom). This is inferred through the ‘tails’ or ‘trails’ seen in these regions of the density map.

The radio halo in the left panel of Figure 1 roughly follows the thermal gas in the X-ray image (right panel of Figure 1) as seen in the level edge present at the bottom of the diffuse structure and its overall shape in both images. The thermal gas and diffuse emission follow an elongated shape with an extended bottom left feature which hosts the BCG. The X-ray image orientation of the extended bottom left feature indicates that gas in the subcluster is undergoing ram pressure stripping as it interacts with the main cluster component.

Furthermore, there is a velocity offset of $\sim 950$ km/s between the mean velocity of the cluster and that of the BCG, further indicating that J0019 is a disturbed cluster. There is also a discrepancy between the SZ mass ($M_{SZ500}$), the X-ray estimated mass ($M_{X500}$), and dynamical mass estimated from the optical redshift distribution ($M_{dyn500}$). $M_{X500}$ and $M_{SZ500}$ lie $> 3\sigma$ away from $M_{dyn500}$, suggesting that the two components, northwest and southeast of the cluster centre, have already merged and J0019 is now observed in a post-merger phase. The gas traced by X-ray and SZ emission is completely disturbed during a previous collision, showing an enhanced X-ray emission and hence, a $T_X - M_X$ overestimation.

The angular resolution and short baselines of the MeerKAT telescope allow us to investigate both low surface brightness diffuse emission as well as compact radio sources. The radio contours overlaid on the optical image are shown in Figure 5 where 10 compact radio sources are identified within a radius, $R = 500$ kpc. The compact source properties are listed in Table 2.

Only one of the radio sources is found in the Faint Images of the Radio Sky at Twenty-Centimetres survey (FIRST [48]). This source is labelled S5 and has an integrated flux density, $S = 0.56$ mJy. Source S3 in Figure 5 is associated with the BCG of the main cluster component and S4 is associated with the BCG of the infalling subcluster. This is inferred from the optically visible BCGs. There are two BCGs that are spatially separated, as seen in the DES image in Figure 5. These BCGs are also separated in velocity space as shown in Figure 4. This provides support for the existence of two distinct galaxy populations.

Moreover, an arched structure is found within the radio halo in projection. The optical image (Figure 5) shows no optical counterpart at the centre of this arched structure, S1, suggesting that it is embedded in the diffuse structure. There are, however, two optical galaxies at either end of the bent radio structure. Spectroscopy of all cluster members are needed to gauge if these galaxies are related to the radio emission or not.
Figure 5. Optical gri-band DES image with MeerKAT contours overlaid. Contours are without short baselines and start at [3, 4.5, 6, 10, 20, 40, 80] × σ rms where the rms is the full resolution central noise (\(\sim 18.1\) µJy/beam). The white dashed circle represents \(R_{500}\). The radio image beam is shown as the white ellipse in the lower left corner (7.9′ × 7.2′, p.a. 12°) and the red cross marks the position of the ACT SZ peak. Individual radio galaxies are labelled from S1 to S10.

Table 2. Properties of cluster region radio sources. Source labels are shown in Figure 5. The R.A. and Dec. values are for the peak source emission in the radio map. Flux densities for the sources are measured in the 1.16 GHz map. Spectroscopic redshifts, \(z_{\text{Spec}}\), are from SDSS DR16. * C: compact; T: resolved with tailed emission.

<table>
<thead>
<tr>
<th>Source</th>
<th>R.A. (deg)</th>
<th>Dec. (deg)</th>
<th>Type (^a)</th>
<th>(S_{1.16\text{GHz}}) (mJy)</th>
<th>(z_{\text{Spec}})</th>
<th>Notes</th>
</tr>
</thead>
<tbody>
<tr>
<td>S1</td>
<td>4.90205</td>
<td>3.59912</td>
<td>T</td>
<td>0.30 ± 0.03</td>
<td></td>
<td>No optical counterpart found</td>
</tr>
<tr>
<td>S2</td>
<td>4.91539</td>
<td>3.59455</td>
<td>C</td>
<td>0.67 ± 0.03</td>
<td></td>
<td></td>
</tr>
<tr>
<td>S3</td>
<td>4.91164</td>
<td>3.59829</td>
<td>C</td>
<td>0.04 ± 0.01</td>
<td>0.26904</td>
<td>Associated with BCG of the main component</td>
</tr>
<tr>
<td>S4</td>
<td>4.90705</td>
<td>3.62907</td>
<td>C</td>
<td>1.70 ± 0.03</td>
<td>0.26415</td>
<td>Associated with BCG of infalling subcluster</td>
</tr>
<tr>
<td>S5</td>
<td>4.89830</td>
<td>3.63448</td>
<td>C</td>
<td>0.69 ± 0.02</td>
<td></td>
<td>Detected in FIRST(0.56 mJy)</td>
</tr>
<tr>
<td>S6</td>
<td>4.91872</td>
<td>3.60286</td>
<td>C</td>
<td>0.10 ± 0.01</td>
<td></td>
<td></td>
</tr>
<tr>
<td>S7</td>
<td>4.91153</td>
<td>3.64177</td>
<td>T</td>
<td>0.66 ± 0.03</td>
<td></td>
<td></td>
</tr>
<tr>
<td>S8</td>
<td>4.92836</td>
<td>3.60541</td>
<td>C</td>
<td>0.15 ± 0.02</td>
<td></td>
<td></td>
</tr>
<tr>
<td>S9</td>
<td>4.92806</td>
<td>3.63719</td>
<td>C</td>
<td>0.11 ± 0.01</td>
<td></td>
<td></td>
</tr>
<tr>
<td>S10</td>
<td>4.90081</td>
<td>3.61702</td>
<td>C</td>
<td>0.05 ± 0.01</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

5. Conclusions

A low surface brightness radio halo was detected in J0019 [12] with the MeerKAT telescope at L-band. The halo flux density estimated by Knowles et al. [12] is \(S_{1.16} = 9.16 ± 0.57\) mJy. The calculated k-corrected 1.4 GHz radio power, \(P_{1.4\text{GHz}}\), is \((1.0 ± 0.3) \times 10^{24}\) W/Hz and has been estimated by extrapolating the 610 MHz flux density to 1.4 GHz using a theoretically motivated spectral index of \(\alpha = -1.3 ± 0.4\). In this study, we performed a dynamical state analysis of J0019 using both XMM-Newton X-ray observa-
tions and optical DES photometry. The estimated morphology parameters are as follows: 

\[ \frac{P_3}{P_0} = (8.93 \pm 0.4) \times 10^{-8}, \ w = 0.03 \pm 0.01, \ c_{SB} = 0.13 \pm 0.002, \ A_2 = 0.068 \pm 4.9 \times 10^{-4}, \]

and \( CS = 0.20 \pm 0.16 \text{ Mpc} \).

Although the cluster is known to be involved in a merger event through the detection of the radio halo, the X-ray \( \frac{P_3}{P_0} \) value, optical centre shift, and optical asymmetry values in the morphology analysis are inconclusive. The X-ray centroid shift and concentration values indicate a dynamically disturbed system.

The optical density map shows two cluster components: one along the northwestern and another in the southeastern direction. The BCG of the main component lies in the southern region and the BCG of the infalling cluster lies in the northern region as seen in the DES gri-band image and optical galaxy density map. Moreover, the density map shows that the cluster is elongated and axisymmetric.

The bimodality detected in the velocity distribution and the offset of the BCG with respect to the mean velocity of the cluster suggests that the cluster merger has a component in the line-of-sight direction. However, the DES density maps, the MeerKAT radio halo, and XMM X-ray surface brightness shape shows an elongation in the north–south direction, which suggests that the collision also presents a velocity component contained in the plane of the sky and follows this north–south axis.

With an elongated shape in the X-ray map, the clear bimodality and axisymmetric structures in the optical maps, two X-ray morphology parameters, the BCG offset by \( \sim 950 \) km/s from the mean velocity, and the discrepancy between the SZ mass \( (M_{SZ,500} = (10.2 \pm 2.28) \times 10^{14} M_\odot) \), X-ray mass \( (M_{X,500} \sim (6.5 \pm 1.3) \times 10^{14} M_\odot) \), and dynamical mass \( (M_{dyn,500} = (2.5 \pm 1.2) \times 10^{14} M_\odot) \), we clearly show that J0019 is a merging cluster and probably in a post-merging phase.

We can aim to obtain spectroscopy for the full field to identify more cluster members, infer cluster dynamics, and model how the merger event takes place in future work. An additional aim could be to study both the spectral index map of the diffuse structure and the polarized radio map. The energetics of the system can also be studied in future work with simulations.

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**Conflicts of Interest:** The authors declare no conflict of interest.

**Notes**

2. [https://github.com/IanHeywood/oxkat](https://github.com/IanHeywood/oxkat) (accessed on 1 June 2020.)
4. [https://github.com/saopicc/killMS](https://github.com/saopicc/killMS) (accessed on 1 July 2020).
6. The XMM-Newton observations include European Photon Imaging Camera (EPIC) data from the two MOS (Metal Oxide Semi-conductor) CCD arrays and the pn CCD array.
7. XMM: Generate and Analyse GitHub (accessed on 1 May 2021).
8. Spherical radius that encloses an average density equal to 500 times the critical density at the cluster redshift, z = 0.266.
9. Corrected for the cosmology in this paper.

**References**
