Double Chooz $\theta_{13}$ measurement via total neutron capture detection

Article (Accepted Version)


This version is available from Sussex Research Online: http://sro.sussex.ac.uk/id/eprint/102883/

This document is made available in accordance with publisher policies and may differ from the published version or from the version of record. If you wish to cite this item you are advised to consult the publisher's version. Please see the URL above for details on accessing the published version.

Copyright and reuse:
Sussex Research Online is a digital repository of the research output of the University.

Copyright and all moral rights to the version of the paper presented here belong to the individual author(s) and/or other copyright owners. To the extent reasonable and practicable, the material made available in SRO has been checked for eligibility before being made available.

Copies of full text items generally can be reproduced, displayed or performed and given to third parties in any format or medium for personal research or study, educational, or not-for-profit purposes without prior permission or charge, provided that the authors, title and full bibliographic details are credited, a hyperlink and/or URL is given for the original metadata page and the content is not changed in any way.

http://sro.sussex.ac.uk
First Double Chooz $\theta_{13}$ Measurement via Total Neutron Capture Detection

The Double Chooz Collaboration

H. de Kerret$^{a,d}$, T. Abrahao$^a$, H. Almazan$^o$, J. C. dos Anjos$^o$, S. Appel$^j$, J. C. Barriere$^k$, I. Bekmam$^a$, T. J. C. Bezerra$^a$, L. Bezrukov$^l$, E. Blucher$^d$, T. Brugi`ere$^q$, C. Buck$^o$, J. Busenitz$^b$, A. Cabrera$^{11,d,aa}$, M. Cerrada$^h$, E. Chauveau$^d$, F. Chimenti$^{12}$, O. Corpaceh$^d$, J. V. Dawson$^d$, Z. Djuricic$^c$, A. Etenko$^h$, H. Furuta$^a$, I. Gil-Botella$^a$, A. Givaudan$^d$, H. Gomez$^d$, L. F. G. Gonzalez$^c$, M. C. Goodman$^c$, T. Harai$^m$, J. Haser$^c$, D. Hellwig$^k$, A. Hourlier$^{d,k}$, M. Ishitsuka$^{k,l}$, J. Jochum$^w$, C. Jollet$^y$, K. Kald$^{a,n}$, M. Kaneda$^b$, M. Karakad$^t$, T. Kawasaki$^l$, E. Kemp$^d$, D. Krynd$^d$, M. Kuze$^c$, T. Lachenmaier$^w$, C. E. Lane$^l$, T. Lasserre$^{b,k}$, C. Lastoria$^h$, D. Lhuillier$^k$, H. P. Lima Jr$^a$, M. Lindner$^b$, J. M. Lpez-Castaño$^b$, J. M. LoSecco$^b$, B. Lubandsorzhie$^v$, J. Maeda$^{u,m}$, C. Mariani$^i$, J. Marcic$^{i,5}$, J. Martinov$^o$, T. Matsubara$^{u,f}$, G. Mention$^a$, A. Meregaglia$^i$, T. Miletici$^{17}$, R. Milincic$^{i,5}$, A. Minotti$^{k,18}$, D. Navas-Nicol$^a$, P. Novella$^{b,19}$, L. Oberauer$^v$, M. Obolensky$^d$, A. Onillon$^{d,k}$, A. Oralbaev$^n$, C. Palomares$^b$, I. M. Pepe$^e$, G. Pronost$^{11}$, J. Reichenbacher$^{b,11}$, B. Reinhold$^{d,15}$, S. Schöner$^v$, S. Schoppmann$^a$, L. Scola$^k$, R. Sharkanov$^i$, V. Sibille$^{k,13}$, V. Sinev$^l$, M. Skorokhvatov$^n$, P. Soldin$^a$, A. Stahl$^a$, I. Stancio$^b$, L. F. F. Stokes$^w$, F. Sukhert$^d$, S. Sukhotin$^u$, T. Sumiyoshi$^u$, Y. Sun$^{b,4}$, C. Veyssiere$^k$, B. ViAUD$^i$, M. Vivier$^k$, S. Wagner$^{d,k}$, C. Wiebusch$^a$, G. Yang$^{11,12}$, and F. Yerm$^t$

$^a$III. Physikalisches Institut, RWTH Aachen University, 52056 Aachen, Germany

$^b$Department of Physics and Astronomy, University of Alabama, Tuscaloosa, Alabama 35487, USA

$^c$Argonne National Laboratory, Argonne, Illinois 60439, USA

$^d$APC, CNRS/IN2P3, CEA/IRFU, Observatoire de Paris, Sorbonne Paris Cité University, 75205 Paris Cedex 13, France

$^e$Centro Brasileiro de Pesquisas Físicas, Rio de Janeiro, RJ, 22290-180, Brazil

$^f$CENBG, CNRS/IN2P3, Université de Bordeaux, F-33175 Gradignan, France

$^g$The Enrico Fermi Institute, The University of Chicago, Chicago, Illinois 60637, USA

$^h$Centro de Investigaciones Energéticas, Medioambientales y Tecnológicas, CIEMAT, 28040, Madrid, Spain

$^i$Department of Physics, Drexel University, Philadelphia, Pennsylvania 19104, USA

$^j$Institute of Nuclear Research of the Russian Academy of Sciences, Moscow 117312, Russia

$^k$IRFU, CEA, Université Paris-Saclay, 91191 Gif-sur-Yvette, France

$^l$Department of Physics, Kitasato University, Sagamihara, 252-0373, Japan

$^m$Department of Physics, Kobe University, Kobe, 657-8501, Japan

$^n$NRC Kurchatov Institute, 123182 Moscow, Russia

$^o$Max-Planck-Institut für Kernphysik, 69117 Heidelberg, Germany

$^p$University of Notre Dame, Notre Dame, Indiana 46556, USA

$^q$IPHC, CNRS/IN2P3, Université de Strasbourg, 67037 Strasbourg, France

$^r$SUBATECH, CNRS/IN2P3, Université de Nantes, IMT-Atlantique, 44307 Nantes, France

$^s$Research Center for Neutrino Physics, Tohoku University, Sendai 980-8578, Japan

$^t$Department of Physics, Tokyo Institute of Technology, Tokyo, 152-8551, Japan

$^u$Department of Physics, Tokyo Metropolitan University, Tokyo, 192-0397, Japan

$^v$Physik Department, Technische Universität München, 85748 Garching, Germany

$^w$Kepler Center for Astro and Particle Physics, Universität Tübingen, 72076 Tübingen, Germany

$^x$Universidade Estadual de Campinas-UNICAMP, Campinas, SP, 13083-970, Brazil

$^y$Center for Neutrino Physics, Virginia Tech, Blacksburg, Virginia 24061, USA

$^z$LNCA Underground Laboratory, IN2P3 - CEA, Chooz, France

$^\dagger$Current Institutions

July 27, 2021

*Deceased.
Neutrinos were assumed to be massless particles until the discovery of neutrino oscillation process. This phenomenon indicates that the neutrinos have non-zero masses and the mass eigenstates ($\nu_1$, $\nu_2$, $\nu_3$) are mixing of their flavour eigenstates ($\nu_e$, $\nu_\mu$, $\nu_\tau$). The oscillations between different flavour eigenstates are described by three mixing angles ($\theta_{12}$, $\theta_{23}$, $\theta_{13}$), two differences of the square of neutrino masses ($\Delta m^2_{21}$, $\Delta m^2_{31}$) and a charge conjugation parity symmetry (CP) violating phase $\delta_{\text{CP}}$. The Double Chooz (DC) experiment, located near the Chooz Electricité de France reactors, France, measures the oscillation parameter $\theta_{13}$ using reactor neutrinos. In this paper, DC reports its latest $\theta_{13}$ result, $\sin^2(2\theta_{13}) = 0.105 \pm 0.014$, exploiting its multi-detector configuration, iso-flux baseline, reactor-off data and a novel total neutron capture detection technique. Since $\theta_{13}$ was the last unknown mixing angle, which is necessary to measure $\delta_{\text{CP}}$, the result has contributed to the completion of a quest of the neutrino oscillation studies lasting half a century and to pave way toward the CP violation measurement. In addition, DC provides the most precise measurement of the reactor neutrino flux to date, given by the mean cross section per fission $\langle \sigma_f \rangle = (5.71 \pm 0.06) \times 10^{-43}\text{cm}^2/\text{fission}$. Due to transitions between neutrino flavours ($\nu_e$, $\nu_\mu$, $\nu_\tau$), the neutrino masses are generated and the mass eigenstate of the neutrino system ($\nu_1$, $\nu_2$, $\nu_3$) becomes a superposition of the flavour eigenstate. When considering the simpler two flavour ($\nu_\alpha$, $\nu_\beta$) case, the mass eigenstate ($\nu_1$, $\nu_2$) expresses as

$$\begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_\alpha \\ \nu_\beta \end{pmatrix}. \quad (1)$$

If we started with $\nu_\alpha$, we might observe a $\nu_\beta$ at a certain distance $L$, due to neutrino oscillations. The probability of the $\nu_\beta$ appearance as a function of the distance is

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta \sin^2 \frac{\Delta m^2 L}{4E}, \quad (2)$$

where $E$ is neutrino energy and $\Delta m^2 = m_3^2 - m_1^2$ is the difference of square of $\nu_2$ and $\nu_1$ masses. The probability oscillates due to the interference between the amplitudes of propagation: $\nu_\alpha \rightarrow \nu_1 \rightarrow \nu_\beta$) and ($\nu_\alpha \rightarrow \nu_2 \rightarrow \nu_\beta$). The disappearance probability of $\nu_\beta$ is therefore expressed as

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - P(\nu_\alpha \rightarrow \nu_\beta). \quad (3)$$

The first experimental evidence for neutrino oscillations and, still, most of the information today relies on high precision disappearance measurements with about 50 years of history.

The establishment of the neutrino oscillation phenomenon [1][2][3] came as a solution of atmospheric and solar neutrino anomalies around the year 2000. Those results had indicated two consequences: a new oscillation mode, labelled $\theta_{13}$, and the possibility to observe CP-violation, if $\theta_{13}$ was sizeable. The reactor neutrino experiments CHOOZ [4] and Palo Verde [5] set upper limit of $\sin^2 2\theta_{13} < 0.15$ already before 2001. The Double Chooz (DC) group was formed in 2006 to measure the $\theta_{13}$ more precisely making use of near and far detector configuration [6]. The DC experiment has played a pioneering role in this oscillation channel by providing the first positive evidence, in 2011 [7], in combination with the $\nu_\mu \rightarrow \nu_e$ appearance results of T2K [8] and MINOS [9] experiments. The establishment of $\theta_{13}$ awaited the Daya Bay experiment’s observation in 2012 [10]: confirmed soon after by the RENO experiment [11]. Today’s world best value [12] is driven by the statistical combination of the latest published $\theta_{13}$ results [13][14][15][16][17]. A reassuring feature for the field is that all reactor $\theta_{13}$ experiments are redundant which is critical to ensure a robust and unambiguous result. Therefore, multi-experimental validation framework is highly beneficial. A working group formed by all three experiments is on-going with the goal to assess both internal consistency and coherent systematics treatment of each experiment.

Besides reactor experiments, neutrino beam experiments, such as T2K [18], NOvA [19] and MINOS [20], are also sensitive to $\theta_{13}$ via the sub-dominant appearance $\nu_\mu \rightarrow \nu_e$ and $\nu_\mu \rightarrow \nu_\tau$ oscillation modes. However, their ability for a high precision measurement of $\theta_{13}$ is limited by uncertainties and unknowns such as the $\delta_{\text{CP}}$ and $\theta_{23}$ octant degeneracy. Conversely, this channel allows them to explore neutrino oscillation CP-violation directly. Currently, the beam experiments provide the first CP-violation explorations [18][19] by using the value of $\theta_{13}$ from reactor experiments as input. The latest data allow beam experiments to obtain the first hints of non-zero $\delta_{\text{CP}}$. This result embodies a remarkable demonstration of synergy and compatibility across both reactor and beam measurements.

In this article, DC reports its multi-detector results for the first time with its 4th data release comprising 865 days exposure. DC measures reactor $\nu_e$ coming from EDF company Chooz twin reactors by identical near and far neutrino detectors. The $\theta_{13}$ value was extracted from the disappearance and spectrum distortion of the reactor $\nu_e$ caused by the baseline difference of the two detectors. Most of the systematic uncertainties are cancelled out by using functionally identical detectors observing the same reactor $\nu_e$ sources. DC can provide clean analysis results by using several unprecedented techniques such as the exploitation of the effective iso-flux site geometry which cancels possible difference between the two reactors, a model-independent background estimation via reactor power modulation including reactor-off data and the total neutron capture detection technique which significantly increases the neutrino event statistics. Those analysis details are also explained. In addition, the DC near detector is used to characterise the rate and shape difference between data and predicted spectra, and measure the most precise neutrino flux to date, given by the mean cross-section per fission $\langle \sigma_f \rangle$. The latter can be used as reference in other reactor neutrino experiments for an accurate neutrino event rate prediction.

### The Double Chooz Experiment

The DC experiment relies on two identical detectors [21] and two of the most powerful pressurised water reactors of the N4 plant series, whose full power is 8.5 GW thermal power
Figure 1: Laboratory and DC Detectors. The underground LNCA (Laboratoire Neutrino Champagne-Ardenne) site allows for an almost iso-flux geometry (left) to the two identical DC detectors (right) yielding major inter-detector cancellation of reactor flux and detection systematics. Active BG rejection is achieved by the exploitation of the multi-layer (blue shaded) liquid-scintillator design whose light is read out by many photomultipliers via a Flash-ADC deadtime-less electronics. The “Inner Detector” (ID) is subdivided into 3 optically coupled volumes: i) ν-Target (GdT, 10 m³ liquid scintillator Gd 1 g/l loaded), ii) γ-Catcher (GC, 23 m³ liquid scintillator), and iii) buffer (100 m³ non-scintillating oil). The “Inner Veto” (IV, 0.5 m thick liquid scintillator) fully surrounds the ID while the “Outer Veto” (OV, tracking plastic scintillator strip) is placed on the top. The IV tags external rock γ’s (anti-Compton veto), fast-neutrons and cosmic μ’s while the OV only sees cosmic μ’s, covering the ID chimney region. An external inert shield surround the IV: 15 cm steel (FD) and 1 m water (ND). The glove-box allows clean and safe deployment of the same calibration sources (252Cf, 60Co, 68Ge, 137Cs) in both the ND and FD.

(i.e. \( \sim 10^{21} \bar{\nu}_e/s \) flux). The near (ND) and far (FD) detectors are located at the respective average distance of \(~400\) m and \(~1050\) m to the Chooz reactors (B1 and B2). The \( \theta_{13} \) signature manifests itself as a rate deficit with an up to \(~10\%) spectral distortion in the FD relative to the almost undistorted ND spectrum. Thus, DC performs a “rate+shape” \( \theta_{13} \) measurement whose statistical uncertainty is dominated by the FD. There are three types of sources of systematic uncertainties: detection (including the estimation of the neutrino energy) and background are internally constrained with DC data, while the reactor flux relies on an external reactor model. The model is commonly used by most reactor experiments while here it is customised to the specific DC reactor conditions. The simple Chooz multi-detector site geometry enables to place the ND at the effective iso-flux position relative to the FD. This means implying the condition \( L_{B1,ND}/L_{B1,FD} \approx L_{B2,ND}/L_{B2,FD} \) for each reactor-detector pair distance (L). This way, both the FD and the ND are exposed to both reactors with the same fraction. From the single-detector (SD) to the multi-detector (MD) configurations, major systematics cancellation occurs by virtue of correlations due to identical detectors (detection systematics) and the iso-flux reactor geometry (flux systematics). The site and detectors are briefly described in Fig. 1 – see Appendix for details. In this release, 481 days of data from single detector operation (FD-I, April 2011 until January 2013) previous to commissioning of the ND and 384 days of data with both detectors FD and ND (FD-II, January 2015 until April 2016) are combined. The result presented here supersedes our previous [13, 14].

Each N4 reactor typically runs at maximum power allowing the lowest power uncertainty (0.5%), or else they stop a few weeks once per year to refuel. The Chooz total reactor power modulation allows for “2-reactors” (both on), “1-reactor” (either on) and the unique BG only “0-reactor” (both off) data sets. An exposure of \( \sim 25 \) days of 0-reactor data is available. Past FD-I SD results [13, 14] employed the Bugey4 experiment data [23] to compensate the ND absence, improving the overall systematics.

The ND monitors the rate+shape of the flux, thus reducing the \( \theta_{13} \) uncertainty from most of the reactor physics and running configuration. The ND is a direct reactor monitor of the FD-II (iso-flux) and indirect to the FD-I. The iso-flux implies that the neutrino fluxes are expected to be largely correlated across detectors, with negligible impact from reactor power or composition variations. This correlation translates into an almost total rate+shape flux-error cancellation [6, 24]: a

<table>
<thead>
<tr>
<th>Spectrum</th>
<th>Single reactor uncertainty (%)</th>
<th>Inter-Reactor SD</th>
<th>MD</th>
</tr>
</thead>
<tbody>
<tr>
<td>Bugey4 ( \sigma_f )</td>
<td>2.19/0.66(^{14})</td>
<td>Corr.</td>
<td>Corr.</td>
</tr>
<tr>
<td>Fission Fractions</td>
<td>0.69/0.78(^{14})</td>
<td>Corr.</td>
<td>Uncorr.</td>
</tr>
<tr>
<td>Thermal Power</td>
<td>0.47</td>
<td>Corr.</td>
<td>Uncorr.</td>
</tr>
<tr>
<td>Total</td>
<td>2.27/1.68(^{14})</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

\(^{1}\)The detector locations are slightly off from iso-flux while this has negligible impact.
Figure 2: The TnC Detection Principle. The IBD acceptance criteria are widely opened to integrate over all capture-$\gamma$'s: $\sim2.2$ MeV (H-n), $\sim5.0$ MeV (C-n) and $\sim8$ MeV (Gd-n). The overwhelming accidental BG (ND) is rejected over $>4$ orders of magnitude below 3.5 MeV (top) using vetoes and the ANN selection. Excellent data (blue points) to MC (red area) agreement is found below 3.5 MeV (top) using vetoes and the ANN selection. The mean capture time $\tau_{\text{capture}}$ is $\sim200$ $\mu$s in a metal-free organic liquid scintillator. The energy scale uncertainty has negligible impact ($<0.05\%$) due to 1.3 MeV cut. The selection efficiency (bottom) of the Gd-only $[13]$ (left) confines IBD’s to the Gd presence. The TnC (right) yields detection in the full volume; i.e. Gd-T ($>95\%$) and GC ($>80\%$). The relative yields are $\sim61.3\%$ H-n, $\sim38.2\%$ Gd-n and $\sim0.5\%$ C-n.

Neutrino Detection by Total Neutron Capture

Since the discovery of the neutrino, reactor $\bar{\nu}_e$ are typically detected via the inverse-beta-decay (IBD: $\bar{\nu}_e + p \rightarrow n + e^+$) [31] interactions on free protons (i.e. H nuclei) via a coincidence technique, where the prompt trigger ($e^+$) is followed by the delayed trigger (neutron capture) several tens of $\mu$s later. The mean capture time $\tau_{\text{capture}}$ is $\sim200$ $\mu$s in a metal-free organic liquid scintillator. The energy scale uncertainty has negligible impact ($<0.05\%$) due to 1.3 MeV cut. The selection efficiency (bottom) of the Gd-only [13] (left) confines IBD’s to the Gd presence. The TnC (right) yields detection in the full volume; i.e. Gd-T ($>95\%$) and GC ($>80\%$). The relative yields are $\sim61.3\%$ H-n, $\sim38.2\%$ Gd-n and $\sim0.5\%$ C-n.

The IBD accep-
tance criteria are widely opened to integrate over all capture-$\gamma$'s: $\sim2.2$ MeV (H-n), $\sim5.0$ MeV (C-n) and $\sim8$ MeV (Gd-n). The over-
whelming accidental BG (ND) is rejected over $>4$ orders of mag-
nitude below 3.5 MeV (top) using vetoes and the ANN selection. Excellent data (blue points) to MC (red area) agreement is found in the delayed energy distribution after the rejection. The energy scale uncertainty has negligible impact ($<0.05\%$) due to 1.3 MeV cut. The selection efficiency (bottom) of the Gd-only [13] (left) confines IBD’s to the Gd presence. The TnC (right) yields detection in the full volume; i.e. Gd-T ($>95\%$) and GC ($>80\%$). The relative yields are $\sim61.3\%$ H-n, $\sim38.2\%$ Gd-n and $\sim0.5\%$ C-n.

The novel TnC has several remarkable features. Mainly, the TnC integrates all n-capture elements, thus the $\bar{\nu}_e$ detection is independent from specific capture details. This implies that the detection volume expands to both the GdT (Gd-target) and GC (gamma-catcher), as shown in Fig. 1, so the TnC volume increases $\sim3\times$ as compared to Gd-only. This boost in statistic is critical for the DC sensitivity. In addition, the statistical limitation of the selection systematics is reduced as the selection efficiency per volume increases close to 100% in GdT and $\sim80\%$ in GC due to the ANN accidental rejection. The wider TnC acceptance integrates over most MC inaccuracies, including the complex spill-in/out effect across the GdT-to-GC boundary. The only relevant boundary is the simpler GC-to-buffer with only H-C on both sides. So, TnC matches MC better while all element-dependent terms, such as Gd or H fractions, are irrelevant. A small concentration of Gd was found inside the GC of the ND due to leakage from the GdT. The described element independence also makes the TnC leak insensitive demonstrating selection stability within 0.1% in both the ND and FD.

unique DC feature as compared to other reactor-$\theta_{13}$ experiments [25]. Instead, the FD-I benefits only from partial error cancellation. The ND provides the reference oscillation spectrum for both FD-I and FD-II for the $\theta_{13}$ measurement. The shape differences between the FD-I and ND due to the different fuel composition are less than 0.5%. The flux systematics reduce from 2.27% (1.68% with Bugey4 data – see Table 1) in the SD case to $\sim0.66\%$. It strongly reduces to $\sim83\%$ in MD configurations. The uncorrelated uncertainty between the ND and FD-I predictions is estimated to be $0.66\%$. It strongly reduces to $\sim0.1\%$ between the ND and FD-II predictions because of the iso-flux configuration and simultaneous data taking. In brief, the unique DC geometry grants a framework for a high-precision measurement by cancellation of the flux systematics.

Despite the ND, the reactor $\bar{\nu}_e$ prediction model remains an important element to the analysis, key for the FD-I. The same strategy employed in past publications [13] is adopted. An external reactor $\bar{\nu}_e$ prediction model is used [24, 27] where $^{235}$U, $^{239}$Pu, $^{241}$Pu fissile isotope contributions rely on ILL data [25]. DC uses a measurement [29] for the $^{238}$U prediction, while Daya Bay and RENO use summation methods. Dedicated Chooz reactor simulations provide the fission fraction ($\alpha_f$) evolution considering the thermal power data ($P_{\text{th}}$) and re-fuel inventories. Bugey4 data aid both to constrain the predicted rate [30] and to improve the precision.

$^{2}$~900 IBD/day (ND) and ~140 IBD/day (FD) with 2 reactors on.
The IBD Signal & Backgrounds

The BG consists of all physical events mimicking the IBD time-space coincidence implied by the TnC selection. This includes accidental and correlated BGs. BG rejection is much harder whenever there is a neutron in the final state. Due to the small overburden, cosmogenic BGs are dominant: fast-neutrons and unstable isotopes produced from 12C spallation, such as 9Li. The signature of fast-neutrons is a recoil on H, as prompt, followed by the n-capture, as delayed. One or more neutrons can participate in a fast-neutrons coincidence. 9Li undergoes a $\beta$-n decay (including $\alpha$’s), but no indications of $^8$He production are found \textsuperscript{33}. Since the detector chimney is an effective tagging hole to vertical $\mu$’s, there is a potential background when they stop inside the active detector volume. Both $\mu$ decay at rest (Michel e\textsuperscript{±}) and $\mu$ capture have been carefully studied \textsuperscript{34} and rejected to a negligible level. Lastly, the accidental BG is caused by two independent events, mainly from radioactive isotopes inside or around the detectors.

An offline 1.25 ms veto is imposed after each tagged $\mu$, thus rejecting the resulting fast neutrons’ captures ($>5 \times 10^{17}$) and stopping-$\mu$’s. The live time loss is 5.4\% (FD) and 25.5\% (ND). DC has developed a multi-veto method yielding large BG rejection: a rejection factor of 200 is reached for the FD and 35 for the ND as compared to simple time coincidence. The vetoes are defined identically across ND and FD. Cosmogenic BG rejection relies on direct $\mu$ and/or neutron tagging

\textsuperscript{3}ND and FD are, respectively, at $\sim$30 m and $\sim$100 m rock overburden depth, so their cosmic $\mu$ rates are, respectively, $\sim$240 s\textsuperscript{−1} and $\sim$45 s\textsuperscript{−1}. 

### Table 2: Detection Uncertainties

<table>
<thead>
<tr>
<th>Uncertainty (%)</th>
<th>SD</th>
<th>MD</th>
</tr>
</thead>
<tbody>
<tr>
<td>Proton Number</td>
<td>0.65</td>
<td>0.39</td>
</tr>
<tr>
<td>IBD Selection</td>
<td>0.33\textsuperscript{FD} / 0.12\textsuperscript{ND}</td>
<td>0.27</td>
</tr>
<tr>
<td>Boundary Effect</td>
<td>0.20</td>
<td>–</td>
</tr>
<tr>
<td>Vetoes Efficiency</td>
<td>$\leq$0.05</td>
<td>$\leq$0.05</td>
</tr>
</tbody>
</table>

The central column shows the uncertainties on the signal normalisation for the single detector (SD) case. The multi detector (MD) case in the column on the right shows the uncertainty on the ratio of the signal rates (FD/ND) assuming simultaneous operation (not fully representative for DC fit configuration). The total systematics is dominated by the uncertainty of the number of protons for IBD interactions (mainly the GC). This is to be re-measured upon future detector dismantling. The TnC reduces selection systematics as compared to the element dependent detection, since it is not sensitive to the knowledge of the Gd/H fraction of neutron captures. Boundary systematics relying on the modeling of spill-in/out events at volume interfaces in the MC are assumed fully correlated between detectors. The selection systematics rely on an IBD data-driven method, thus inclusively accounting and averaging over selection and energy scale (stability, uniformity and linearity) variations. The vetoes play a negligible role as they were optimized to maximize the selection efficiency while adding a negligible systematic.

### Table 3: IBD Candidates Background

<table>
<thead>
<tr>
<th>Rate (day\textsuperscript{−1})</th>
<th>FD</th>
<th>ND</th>
</tr>
</thead>
<tbody>
<tr>
<td>IBD Candidates</td>
<td>112</td>
<td>816</td>
</tr>
<tr>
<td><strong>Breakdown</strong></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Accidental</td>
<td>4.13 ± 0.02</td>
<td>3.110 ± 0.004</td>
</tr>
<tr>
<td>Fast-Neutron</td>
<td>2.50 ± 0.05</td>
<td>20.85 ± 0.31</td>
</tr>
<tr>
<td>$^9$Li Isotope [(\mu\text{-tag})]</td>
<td>2.62 ± 0.27</td>
<td>14.52 ± 1.48</td>
</tr>
<tr>
<td>Stopped-$\mu$</td>
<td>$&lt;0.19 @ 98%\text{CL}$</td>
<td>$&lt;0.21 @ 98%\text{CL}$</td>
</tr>
<tr>
<td>Others ($^{12}$B, BiPo)</td>
<td>$&lt;0.01$</td>
<td>$0.04 ± 0.01$</td>
</tr>
<tr>
<td><strong>Total</strong></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\Sigma$-Exclusive</td>
<td>9.3 ± 0.3</td>
<td>38.5 ± 1.5</td>
</tr>
<tr>
<td>Inclusive (17 days)</td>
<td>9.8 ± 0.9</td>
<td>39.6 ± 2.5</td>
</tr>
<tr>
<td><strong>Signal to BG</strong></td>
<td>11.0</td>
<td>20.2</td>
</tr>
</tbody>
</table>

The BG consists of all physical events mimicking the IBD time-space coincidence implied by the TnC selection. This
using ID (inner detector), IV (inner-veto) and OV (outer-veto) detectors, as shown in Fig. 1.

A fraction of $^{9}\text{Li}$ and most $^{12}\text{B}$ are tagged by identifying spallation activity via correlated neutrons upon each tracked $\mu$. Fast-neutron rejection exploits direct neutron tagging in the IV or the primary $\mu$ using the OV. The ANN discards most accidents with a rejection factor $>300$. Since IBD's extend to the GC, about 25% of external $\gamma$'s are tagged in the IV acting as anti-Compton veto. Stopped $\mu$'s are fully suppressed based on Michel $e^{\pm}$ discrimination using information of the ID pulse-shapes and goodness of fit information from the position reconstruction. Any veto using the ID only affects delayed triggers to prevent any prompt spectral distortion affecting $\theta_{13}$. The IBD inefficiency for all vetoes combined is 4.5% (FD) and 5.7% (ND) with negligible (<0.05%) systematics. See Appendix for further complementary selection and veto details.

Table 3 summarises the remaining BG estimates including several independent measurements employed to validate the accuracy. The BG rates, larger compared to Gd-only [13], have negligible impact on the determination of $\theta_{13}$. The BG subtraction impact is small for signal-to-background $>10$ while the dominant statistical BG uncertainty is reduced. The vetoed BG samples provide copious data-driven spectra, which contain BG information independent of MC simulations. The shape of cosmogenic BG is found to be identical across ND and FD within statistical uncertainties. The non-flat energy spectrum of the fast neutrons was carefully evaluated as it could mimic the $\theta_{13}$ signature. Its spectra was characterised over an extended window up to 100 MeV. The overall impact of BG on $\theta_{13}$ is marginal. The dominant BG systematic is the $^{9}\text{Li}$ uncertainty. The BG model accuracy was scrutinised independently with $\sim17$ days of inclusive 0-reactor data samples in both ND and FD-II. Thus, these data are not used in the $\theta_{13}$ fit. No non-statistical bias or tension (<1$\sigma$) is found on the measured BG-model, rates and/or spectral shapes.

The $\theta_{13}$ Measurement

The $\theta_{13}$ measurement is obtained by contrasting the observed IBD rate+shape spectral distortion against the specific neutrino oscillation model prediction, similar to Eq.3 in which the rate reduces following the flux modulation given by

$$P(\bar{\nu}_e \to \bar{\nu}_e) \approx 1 - \sin^2 2\theta_{13} \sin^2(1.267\Delta m^2_{\odot} L/E_{\nu_e})$$

where $\sin^2 2\theta_{13}$ is the unknown. $L(m)$ is the baseline distance between each reactor-detector pair, $E_{\nu_e}$ (MeV) is the neutrino energy obtained from the prompt energy deposition or Visible Energy ($E_{\nu_e} \approx E_{\nu_e} + 0.78$ MeV). $\Delta m^2_{\odot}$ is the two-neutrino weighted average of $\Delta m^2_{21}$ and $\Delta m^2_{32}$ [23], where $|\Delta m^2_{\odot}| = (2.484 \pm 0.036) \times 10^{-5}$ eV$^2$ [26] is used as input to the fit.

The $\theta_{13}$ rate+shape fit measurement uses all detectors data simultaneously. The nominal fit considers the input from each SD fit (data to its MC) including pertinent constraints and correlations. The SD fit is shown in Fig.3 (bottom). In our MD analysis, all SD fits (FD-I, FD-II and ND) are simultaneously performed, constrained by the inter-detector correlations such as BG (shape), detection (rate), energy (shape) and flux (rate+shape). Thus, the common ND provides direct and almost un-oscillated rate+shape reference spectrum. Systematic uncertainties cancel due to correlations with both FD-I and FD-II. The iso-flux FD-II benefits from the maximum error cancellation. The $\theta_{13}$ measurement is, in principle, independent from any common or correlated contributions across the MC and detectors. Fig.4 (left) illustrates the inter-detector ratio fit exhibiting the expected $\theta_{13}$ flux modulation and demonstrating the suppression of the spectral distortion against the MC. The common MC serves both as link to the neutrino energy ($E_{\nu_e}$) and an inter-detector comparison mediator. The non-trivial role of the reactor model is scrutinised later on. This is a delicate point since the data to prediction comparison exhibits clear distortions uncovered by the uncertainties such as the 5 MeV excess shown in Fig.8.

The BG constraints benefit from ~8 days of FD-I 0-reactor data (taken in 2011 and 2012) and the 20 MeV range extension. The systematics are treated both via covariance matrices (energy and reactor flux) and nuisance pull terms. In the covariance treatment data points and uncertainty bands in figures do not fully represent the fit constraints upon minimisation. Pull terms are also used in the $\chi^2$ minimisation and give access to physical observables (BG rates, inter-detector normalisation, etc) as well as insight to the fit consistency. There is negligible (<1$\sigma$) tension in all fit output values. The fit strategy follows an unbiasing scheme where performance and robustness to $\theta_{13}$ measurements are scrutinised using MC and fixed prior to the final data fit.

The best fit value is $\sin^2 2\theta_{13} = 0.105 \pm 0.014$ (13.3% precision) with $\chi^2$/DoF = 182/112 with a $p$-value of $3.2 \times 10^{-5}$. The expected total uncertainty (i.e. sensitivity) was 0.014. The statistical precision is 0.005, so systematics largely dominate. The large $\chi^2$/DoF is caused by the mismatch between data and prediction, which is not covered by the model uncertainties. This topic is addressed later on. The $^{9}\text{Li}$ rate is unconstrained in the fit. This is because the fit $^{9}\text{Li}$ sample is up to 50% statistically correlated to the one used for $^{9}\text{Li}$ estimation ($\mu$-to-IBD time correlation), as summarised in Table 3. The rate+shape spectral distortion is shown in Fig.4 (right). An empirical fit to the spectral distortion residual appears to resolve a structure consistent with a slope and one or two Gaussian peaks. Their origin remains unknown. The common normalisation across all detectors can be measured as output of the fit. The value is $1.004 \pm 0.008$, while the input value was constraint by the uncertainty of the Bugay+1 measurement (1.4%) [23]. The sizeable smaller output uncertainty indicates that DC holds valuable independent information about rate normalisation.

Table 4 summarises the contributions to the total uncertainty in the oscillation fit. The FD drives the overall statistical precision (~90k IBD’s). Because the FD-I data set still represents a large fraction of the data analysed here, the flux systematics are the largest contribution to the uncertainty on $\theta_{13}$. However, the operation of the two detectors in iso-flux configuration during FD-II lead to a large reduction of the systematics and this phase drives the overall sensitivity. The detection uncertainties are dominated by the large GC proton-number uncertainty. The BG has a small role, thanks to the statistics and since the $^{9}\text{Li}$ rate is constrained in the
Figure 3: **ND and FD Spectra & SD Ratios.** Both ND (~210k IBD’s) and FD (~90k IBD’s) spectra are shown (top) within the fit [1.0,20.0] MeV range, including the un-oscillated MC prediction (red) and the BG model: accidentals (clear grey), $^7$Li (grey) and fast-neutrons (dark grey). Cosmogenic BGs are estimated during the fit since $^7$Li (unconstraint) dominates in the [7.0,12.0] MeV region and fast-neutrons above 12 MeV. The impact of accidentals to the $\theta_{13}$ measurement is negligible. The data (BG subtracted) to prediction ratio is shown (bottom). The best fit solution (blue) contrasts with the no-oscillation hypothesis (red). Two dominant spectral distortions can be appreciated: the $\theta_{13}$ signature (mainly FD) and a common 5 MeV excess, leading to a large $\chi^2$/DoF of 182/112. Bugey4 constrains the prediction rate. The normalisation with this constraint is lower as compared to the prediction rate not using the Bugey4 information. The cancellation of both common distortions and correlated uncertainties takes place from the MD (yellow) to the ND (green) configurations. The covariances used (not shown) play an important role during the fit.

**Discussion & Implications**

Our reported MD $\theta_{13}$ exhibits an up to 48% higher central value whose significance is $<2.0 \sigma$’s compared to all other measurements. The latest published values of $\theta_{13}$ are shown in Fig. 5. NOvA [19] and MINOS [20] are also sensitive to $\theta_{13}$. Since the statistical uncertainties in reactor experiments are small, a simple statistical fluctuation is unlikely to be the sole cause of the difference. Differences are however today consistent within the context of the dominant systematics uncertainties. The consistency among the DC, DYB and RENO reactor measurements remains critical check for the $\theta_{13}$ final global value used everywhere else.

**Systematic Uncertainty Scrutiny**

The reported $\theta_{13}$ result deserves thorough scrutiny to ensure that the accuracy (i.e. any bias) is controlled well within the quoted uncertainties. In order to do this, DC has performed
Table 4: $\sin^22\theta_{13}$ Measurement Uncertainties Breakdown. The match between the $\theta_{13}$ uncertainty from data (0.0139) and the predicted sensitivity (0.0141) allows for a MC uncertainty breakdown. In the central column the fractional uncertainties of the different systematics ($\Delta$) are given. They were calculated from the sensitivity assuming just one systematic contribution in addition to the statistical uncertainty. The statistical part was then subtracted in quadrature. The total is larger than the square root of the sum of the individual squared uncertainties because of correlations. The difference corresponds to a (0.0065)$^2$ term. The column on the right shows the total uncertainty when the corresponding single systematics is removed. The impact of background and in particular of the energy scale on the sensitivity is higher than one might expect from the values given in the central column. Again, this is due to the correlations.

<table>
<thead>
<tr>
<th>Uncertainty</th>
<th>Single Syst.</th>
<th>Total - $\Delta$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Reactor Flux</td>
<td>0.0081</td>
<td>0.0112</td>
</tr>
<tr>
<td>Detection</td>
<td>0.0073</td>
<td>0.0113</td>
</tr>
<tr>
<td>Energy</td>
<td>0.0018</td>
<td>0.0121</td>
</tr>
<tr>
<td>Background</td>
<td>0.0018</td>
<td>0.0134</td>
</tr>
<tr>
<td>$[\Delta m_{ee}^2]^{-1}$</td>
<td>0.0018</td>
<td>0.0140</td>
</tr>
<tr>
<td>Statistics</td>
<td>0.0054</td>
<td>-</td>
</tr>
<tr>
<td><strong>Total</strong></td>
<td><strong>0.0141</strong></td>
<td><strong>(13.3%)</strong></td>
</tr>
</tbody>
</table>

Figure 4: The Spectral Ratios. The FD to ND ratio (left) represents a clean $\theta_{13}$ rate+shape disappearance evidence used by the fit for parameter extraction. No traces of any remaining distortion are found, demonstrating the expected inter-detector cancellation key to ensure the $\theta_{13}$ accuracy. Instead, ND data to MC prediction ratio (right) allows for precise extraction of the spectral distortion, which is common in FD and ND as seen in Fig. 3. Both the rate and shape effects are visible so rate+shape feature extraction is possible. An empirical structure is examined by fitting with two models: one and two empirical Gaussian peaks with a common slope. Both models reproduced data. The origin of those empirical features remains unknown. Shape-only analysis is described in Appendix.

Impact of Reactor Model on $\theta_{13}$. Since DC data has a limited ability to test the validity of the model, the stability of the $\theta_{13}$ measurement is scrutinised and demonstrated against the behaviour of the reactor model for both the SD and MD configurations as explained in Fig. 6. The SD case is more illustrative as a stronger dependence on the model biases is expected due to the lack of a ND. Also, past DC results with data from the FD only can be directly validated. Indeed, we demonstrate that a measurement of $\theta_{13}$ is compromised using the standard rate+shape model prescription due to the large data to model mismatch. Despite the constraint on the rate from Bugey4, the shape distortion biases the fit via the shape-only term. This effect grows with statistics. The new empirical prescription for FD-I+FD-II data explained in Fig. 6 yields $\sin^22\theta_{13} = 0.108 \pm 0.028$ ($\chi^2$/DoF = 53/74) matching the MD result. This is the best SD $\theta_{13}$ measurement to date. Conversely, the standard reactor uncertainty leads to $\sin^22\theta_{13} = 0.122 \pm 0.022$ ($\chi^2$/DoF: 105/74). Thus, SD is proved a fragile measurement framework due to the unavoidable dependence on the rate+shape reactor model deviations and systematics. The MD case is demonstrated a robust $\theta_{13}$ measurement. With the ND, the inter-detector cancellation protects $\theta_{13}$ largely from any common or correlated rate+shape bias. The

4 The risk to compromise SD accuracy was already suspected in [13], so a rate-only $\theta_{13}$ measurement was conservatively adopted as baseline.
model uncertainty underestimation mainly manifests as the larger $\chi^2$/DoF tension due to the large ND statistical precision. The increase on the uncertainty of the reference spectrum causes both a more robust $\theta_{13}$ value (<1% effect) and the alleviation of the $\chi^2$ tension. The latter was corroborated with data, in which $\chi^2$ went from 182 to 93 (DoF: 112). Instead, the increase of the model uncertainty has almost no impact on the $\theta_{13}$ precision, as shown in Fig. 6.

Our studies allow a few empirical observations linked to today’s model limitations whose origins remain unknown. i) The 1σ envelope for today’s prediction appears insufficient to accommodate the mismatch between data and model for both rate and shape. A better understanding of the origin of model deviations remains critical. In the meantime, the adoption of IBD data driven methods is the only way to bypass the model limitations. Here DC demonstrates that Bugey4 (or alike) can be used to bypass the rate model bias with few per mille accuracy. However, the same is less evident for the spectral shape bias or distortion due to unresolved remaining differences among experiments at the few % level today – see Appendix. ii) The DC IBD data prescription favours the increase of today’s shape-only uncertainty to extend the empirical model, as described in Fig. 6. Unless new physics proves otherwise, significant reactor model progress is needed to attain SD precision below ~6%. iii) The spectral-based bias is expected to depend on the relative position between the dominant features, such as the ~5 MeV excess, and the energy range where the $\theta_{13}$ oscillations affect the spectrum.

**Rate and Shape Decomposition.** Since the $\theta_{13}$ measurement exploits both the rate and the shape, further insight comes from splitting the measurement into the rate-only (16% precision) and shape-only (43% precision) contributions. The main $\theta_{13}$ constraint is due to rate-only (systematics dominated). The shape-only information has enhanced significantly due to the higher statistics in the FD as compared to previous DC results [13]. No tension is found between rate and shape measurements (<0.5σ). The shape-only $\theta_{13}$ fit value is about 20% lower as compared to our main result whereas the rate-only fit is higher. This indicates shape effects as the observed spectral distortions do not introduce a bias towards a higher $\theta_{13}$. The shape stability of the central value is seen by freeing the $\Delta m^2_{32}$ marginalisation: $\sin^22\theta_{13}=0.104^{+0.032}_{-0.025}$ and $[\Delta m^2_{32}] = (2.49^{+0.49}_{-0.40}) \times 10^{-3}$ eV$^2$ are obtained. A loss in precision is expected due to the DC non-optimal baseline. See Appendix for further cross-checks.

**FD-I and FD-II Decomposition.** Thanks to the direct iso-flux monitoring, the reactor flux prediction of FD-II is largely correlated with the ND ($\sim 0.1$% uncorrelated normalisation uncertainty), hence the precision on $\theta_{13}$ is much better as compared to FD-I. It is interesting to decompose the $\theta_{13}$ measurement into the statistically independent FD-I (22% precision) and FD-II (15% precision) samples. Indeed, FD-II drives the reported central value of $\theta_{13}$. Using the SD, the FD-I and FD-II samples are demonstrated to be statistically consistent ($\sim 0.7\sigma$). FD-II (46% of data) is expected to dominate future DC results.
The ND Rate Dependence: $\langle \sigma_f \rangle$. Past DC $\theta_{13}$ results have relied on Bugey4 normalisation constraint. Beyond the better precision achieved, this implies that those measurements depended on the Bugey4 rate normalisation. In the MD case, the impact of Bugey4 cancels across both the ND and FD. The consistency of the ND normalisation can be experimentally tested via the measurement of $\langle \sigma_f \rangle$: i.e. the mean cross-section per fission. The value of $\langle \sigma_f \rangle$ is proportional to the reactor mean flux. So, the reactor luminosity can be used as an inter-experiment reference allowing comparison across experiments. This way, the ND normalisation can be validated against Bugey4 – the most precise measurement to date and reference to all past DC measurements. Fig. 7 shows the data to MC model ratio ($R_{\sigma_f}$) for the most relevant published results so far and their corresponding $\langle \sigma_f \rangle$. The $R_{\sigma_f}(\text{ND})$ is in agreement within $\sigma_3$ dominated by prediction uncertainty, with all other experiments, including the 2017 world average. The prediction normalisation used [35] allows easier comparability across experiments – see further details in the Appendix. Other normalisation estimations [39] are consistent within order 1%. The ND measurement of $\langle \sigma_f \rangle = (5.71 \pm 0.06) \times 10^{-43} \text{cm}^2/\text{fission}$ is the most precise measurement to date. The $\langle \sigma_f \rangle$ central values of Bugey4 and DC(ND) agree $<0.5\%$, once corrected by the relative differences in average fuel composition. This was quantified by the overall DC $\theta_{13}$ fit normalisation output when using the Bugey4 constraint.

Conclusions

The first DC $\theta_{13}$ multi-detector (MD) measurement is presented with a best value $\sin^22\theta_{13} = 0.105 \pm 0.014$. This measurement pioneers the IBD Total Neutron Capture (TnC) detection technique with a major reduction of both SD (single-detector) and MD selection systematics and by a boost of statistics by a factor $\sim 2.5 \times$ relative to Gd-n selection. DC demonstrates the robustness of the quoted background, detection and energy systematics via the articulation of several independent measurements for each systematic. The external reactor flux model systematics are studied and shown to have negligible impact for the extraction of $\theta_{13}$ with the MD configuration. However, the SD $\theta_{13}$ extraction is more sensitive to the systematics of the predicted spectrum. An empirical model extension is introduced here with increased uncertainties for the reactor systematics. In this way an accurate SD $\theta_{13}$ measurement in very good agreement with the MD analysis is demonstrated. DC also reports here for the first time an empirical model of the distortion found between the measured and the predicted spectra. The observed structure of the empirical fit might shed light on the origin behind these deviations. DC also reports here the most precise mean cross-section per fission $\langle \sigma_f \rangle = (5.71 \pm 0.06) \times 10^{-43} \text{cm}^2$ to date, in good agreement with Bugey4 and others experiments. All the results presented here are expected to provide access to the best reactor-$\theta_{13}$ knowledge and related physics. The DC $\theta_{13}$ precision is expected to improve with more statistics and better systematics such as the proton-number.

References

14. DC Collab. (Abe Y et al.) JHEP 01 163 (2016)
Acknowledgments

This publication is dedicated to our colleague Hervé de Kerret. We thank the EDF (“Electricity of France”) company; the European fund FEDER; the Région Grand Est (formerly known as the Région Champagne-Ardenne); the Département des Ardennes; and the Communauté de Communes Ardenne Rives de Meuse. We acknowledge the support of the CEA, CNRS/IN2P3, the computer centre CC-IN2P3 and LabEx UnivEarthS in France; the Max Planck Gesellschaft, the Deutsche Forschungsgemeinschaft DFG, the Transregional Collaborative Research Center TR27, the excellence cluster “Origin and Structure of the Universe” and the Maier-Leibnitz-Laboratorium Garching in Germany; the Ministry of Education, Culture, Sports, Science and Technology of Japan (MEXT) and the Japan Society for the Promotion of Science (JSPS) in Japan; the Ministerio de Economía, Industria y Competitividad (SEIDI-MINECO) under grants FPA2016-77347-C2-1-P and MdM-2015-0509 in Spain; the Department of Energy and the National Science Foundation; the Russian Academy of Science, the Kurchatov Institute and the Russian Foundation for Basic Research (RFBR) in Russia; the Brazilian Ministry of Science, Technology and Innovation (MCTI), the Financiadora de Estudos e Projetos (FINEP), the Conselho Nacional de Desenvolvimento Científico e Tecnológico (CNPq), the São Paulo Research Foundation (FAPESP) and the Brazilian Network for High Energy Physics (RENAFAE) in Brazil.
Appendix

In this section, we shall provide a complementary description about the following topics: a) neutrino oscillation mechanism, b) the DC detectors, c) the TnC IBD selection and BG rejection strategy, d) the cross-check and validation of the θ13 measurement, e) the shape-only characterisation of the spectral distortion and f) the mean cross-section per fission measurement.

Neutrino Oscillation

Neutrino oscillation changes the neutrino flavour periodically while it travels in space-time. Its basic mechanism is the same as the Pauli equation extended to three basic states. The oscillation probability, shown in Eq. 2, is the same as the spin precession in a magnetic field. The neutrino states νe and νμ correspond to the spin-up and spin-down states. θ corresponds to the polar angle of the applied magnetic field and the $\Delta m^2/4E$ term corresponds to the angular velocity of the precession slowed down due to the time dilation of the relativistic effect. There is no parameter in the oscillation form that corresponds to a azimuthal angle of the magnetic field, φ, which is included in the Pauli equation as an imaginary phase. However, actually there are at least 3 neutrinos and for three-state oscillations, an imaginary phase of the transition amplitudes manifests itself as a physical effect, which may violate CP symmetry.

The Double Chooz Near & Far Detectors

The DC detectors are identically designed and are expected to provide identical responses after calibration. This was a key requirement to achieve cancellation of most detection and energy systematics when measuring $\theta_{13}$ with both detectors. During the first phase of the experiment with only the far detector taking data (FD-I with no ND), the DC main effort was devoted to control of the SD systematics thus yielding unprecedented rate & shape uncertainties. This precision has led to among the most precise SD $\theta_{13}$ and $\langle \sigma f \rangle$ measurements to date. A key element for a high SD physics performance is the detector simulation accuracy. This is given by a detailed description of the optical interface, material and geometry, based on Geant4 [40], and the electronics response including PMT (photo-multiplier tube) electronics signals and sampling. Both data and simulation followed identical calibration methods, full event reconstruction and energy response definition. In this way, the simulation is treated as an independent detector able to yield an accurate representation of the data needed for SD systematics reduction.

The DC cylindrical detector design consists of several internal volumes and an external muon detector. From inside out, the main detector (labelled Inner-Detector or ID) is subdivided into three optically coupled volumes: a Neutrino-Target (GdT: 10.3 m$^3$ of liquid scintillator loaded with 1g/l Gd), a Gamma-Catcher (GC: 22.6 m$^3$ of liquid scintillator with no loading), and a Buffer (100 m$^3$ of non-scintillating oil). The ID is fully surrounded by the Inner Veto (IV) detector. The ID and IV are topped by the Outer Veto (OV) muon detector. The IV is a ~0.5 m thick liquid scintillator detector except for a small fraction in the chimney region for the 4π tagging of cosmic μ’s. However, the IV has been able to tag also external rock γ’s (i.e. as an anti-Compton veto) and cosmic neutrons caused by cosmic μ’s going through the nearby rock. The OV is a tracking plastic scintillator for cosmic μ’s with a few cm positioning resolution [41]. Both the ID and IV share one common architecture, i.e. a detector chimney (centre-top) allowing access to the sensitive volumes mainly for calibration deployment. The chimney is an acceptance hole to the IV (only a small fraction of the top), thus allowing for stopped μ’s to reach the ID undetected by the IV. An extension of the OV was placed to cover much – but not all – of the IV where it was missing acceptance due to the chimney.

The detector readout employs PMTs and an 8-bit Flash-ADC electronics sampling at 500 MHz [42]. The ID and IV are instrumented with 390 10” [43] and 78 8” PMTs, respectively operated at a nominal gain of $10^7$. Custom-made front-end electronics ensure the pulse dynamics, including pre-amplification (gain ∼10) and match the Flash-ADC specification for accurate sampling, thus minimising digitisation artefacts. A dedicated global self-trigger system [44] was used to identify IBD interactions using a combined and tuneable energy and multiplicity criterion. The readout energy threshold was kept low (∼0.3 MeV) thus having negligible impact on the relevant physics which starts well above 0.5 MeV. After each trigger, the deadtime-less DAQ causes 256 ns of Flash-ADC sampling to be read out, thus allowing for >150 ns of light scintillating pulse sampling. The Flash-ADC waveforms are reconstructed offline to infer time and charge information per channel yielding <1 ns time resolution and efficient pulse identification from signals with less than 1/5 of a photo-electron (PE). This information seeds the subsequent high-level event reconstruction stages such as vertex position, pulse shape, energy, etc. Further details on event reconstruction were covered in [13]. The Flash-ADC information, once reconstructed, can also provide unique pulse-shape event classification [15][16]. A stringent event-by-event background rejection was possible with negligible detection systematics.

The detector configuration was slightly modified for the FD-II running by implementing a few DAQ optimisations identified during the FD-I period as well as by the increase of the gain per channel by ∼2× to reduce the effect of the Flash-ADC-induced non-linearity [17]. Upon the discovery of the spontaneous light emission [48] effect in the FD, the ND PMTs were covered with black films. The ND exhibits barely no light noise effect, as opposed to the FD. However, careful analysis of the FD data has demonstrated that almost full rejection was possible. Indeed, the PMT light noise has demonstrated to have negligible effect to IBD selection, as also corroborated with the ND data with negligible effect. Hence, both FD-II and ND have all PMTs switched on$^5$ to maximise response linearity and uniformity across both detectors. So, the ND configuration is identical to that of FD-II.

The main energy estimator relies on the PE sum detected

$^5$During the commissioning of the FD-I 14 PMT’s with the highest spontaneous light emission were left off since the experimental rejection of those events have not been fully proved.
The main rationale behind the TnC selection is the wider aperture of the delayed energy window to accept neutron captures from all elements present in the detector; i.e. the H-C elements and the loaded Gd. The TnC selection can be regarded as an effective combination of the well understood Gd-n and H-n selections [13, 14]. Table 5 summarises the TnC selection criteria categorised into $\mu$-tags, singles, IBD coincidences and BG vetoes. The description follows.

The selection starts by identifying $\mu$’s using the $\mu$-tagging criterion such that the sample of singles excludes events just after $\mu$’s. The after $\mu$’s events – rich in cosmic neutron captures – are not used for IBD selection. However, those events are used for self-calibration and regular detector response and...
The IBD selection starts by imposing the IBD coincidence, in both time and space, between the prompt and the delayed neutron capture candidates. The ANN imposes the multi-dimensional prompt-delayed correlation trained to reject random coincidences; i.e. most of the accidental BG. Since the Gd sample has negligible accidental BG contribution (a few per months) [13], the ANN impact the contribution mainly from the GC volume affecting mainly the H sample [13]. The input variables are $\Delta t_{\text{prompt-delay}}$, $\Delta r_{\text{prompt-delay}}$, $\Delta E_{\text{prompt}}$, and $\Delta E_{\text{delay}}$. The first two variables exploit the fact that random coincidences exhibit, by definition, aleatory distributions. For example, the $\Delta t_{\text{prompt-delay}}$ is flat for random coincidences as opposed to correlated events which exhibit the characteristic neutron-capture time distribution. Instead, $\Delta E_{\text{prompt-delay}}$ grows with $t^3$ (saturating within the detector acceptance) for random coincidences, while correlated events are contained within about 1 m. Thus, the ANN mainly exploits these very different measured patterns between correlated and random coincidences to select correlated events, such as IBD’s. The $\Delta E_{\text{delay}}$ has a minor impact since the TnC has a somewhat larger volume. So, the overall training of the ANN makes the rejection insensitive to events with a neutron in the final state, which includes IBD’s, all cosmogenic BG and some remaining accidental BG. This is consistent with the observation, again in Fig. 2 that the remaining accidental BG (dark grey) exhibits a clear irreducible H-n peak contribution. An important intrinsic feature of the ANN is that it internally constructs the best combined selection from all the input variables to achieve an optimal signal identification. This yields a superior background rejection as compared to the combination of simple background vetoes. The overall rejection factors are $\sim 193$ (FD) and $\sim 34$ (ND) relative to the loose coincidence.

Table 5: TnC IBD Selection Criteria & Background Rejection. The complete TnC selection definition is here detailed, including selection criteria and BG vetoes. The type of background rejected by each cut is also highlighted.

<table>
<thead>
<tr>
<th>Observable</th>
<th>Condition</th>
<th>Background</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\mu$ Tagging</td>
<td>$\text{IV:} &gt; 15\text{MeV}$, $\text{ID:} &gt; 100\text{MeV}$</td>
<td>through-$\mu$, $\text{stop-}\mu$</td>
</tr>
<tr>
<td>Energy Deposited</td>
<td>$&gt; 0.3\text{ MeV}$, $&gt; 1.25\text{ ms}$</td>
<td>accidental, after $\mu$ activity</td>
</tr>
<tr>
<td>Light Noise veto</td>
<td>anomalous trigger</td>
<td>no “light noise”</td>
</tr>
<tr>
<td>IBD Coincidence</td>
<td>$\Delta t_{\text{unicity}} &gt; 800,900\text{ ms}$, multi-coincidence</td>
<td></td>
</tr>
<tr>
<td>$\Delta E_{\text{prompt}}$</td>
<td>$[1.0,20.0]\text{ MeV}$</td>
<td>keep IBD</td>
</tr>
<tr>
<td>$\Delta E_{\text{delay}}$</td>
<td>$&gt; 0.85\text{ MeV}$ (FD), $&gt; 0.86\text{ MeV}$ (ND)</td>
<td>accidental</td>
</tr>
<tr>
<td>$\Delta t_{\text{prompt-delay}}$</td>
<td>$[1.3,10.0]\text{ MeV}$</td>
<td>keep n-captures</td>
</tr>
<tr>
<td>$\Delta t_{\text{prompt-delay}}$</td>
<td>$&lt; 1.2\text{ m}$</td>
<td>accidental</td>
</tr>
<tr>
<td>$\Delta t_{\text{unicity}}$</td>
<td>$[1.2,900]\text{ s}$</td>
<td>multi-coincidence</td>
</tr>
<tr>
<td>BG Vetoes</td>
<td></td>
<td></td>
</tr>
<tr>
<td>IV$_\text{prompt}$</td>
<td>IV activity</td>
<td>fast-n, stop-$\mu$</td>
</tr>
<tr>
<td>IV$_\text{delay}$</td>
<td>IV activity</td>
<td>anti-Compton $\gamma$</td>
</tr>
<tr>
<td>OV$_\text{prompt}$</td>
<td>OV activity</td>
<td>fast-n, stop-$\mu$</td>
</tr>
<tr>
<td>OV$_\text{delay}$</td>
<td></td>
<td>Michel-$e^+\text{ stop-}\mu$</td>
</tr>
<tr>
<td>Stop-$\mu$$_\text{decay}$</td>
<td>chimney</td>
<td>$^9\text{Li}$, $^{12}\text{B}$</td>
</tr>
<tr>
<td>Spallation$_\text{prompt}$</td>
<td>n tagged $\mu$</td>
<td></td>
</tr>
</tbody>
</table>

Figure 9: TnC Efficiency & Background Rejection. The evolution of the TnC selection is illustrated in terms of IBD selection efficiency (solid lines), total BG rejection (dotted lines) and the accidental BG rejection (dashed lines). The estimation of the total BG rejection uses 17 days of 0-reactor data. The average single rate per detector is $\sim 10\text{ s}^{-1}$. The first criterion corresponds to a time of $[0.5,800]\text{ ms}$ as a “loose” coincidence with a $[1.0,20.0]\text{ MeV}$ prompt and the $[1.3,10.0]\text{ MeV}$ delayed triggers. The rates are 2291 day$^{-1}$ (FD) and 2375 day$^{-1}$ (ND), which imply a rejection factor of $\sim 357$ relative to singles. These numbers provide an absolute scale to the all other shown below. The $\Delta t_{\text{unicity}} < 1.2\text{ m}$ condition yields some important reduction. However, major accidental BG rejection is only obtained by the ANN with a $\sim 400$ rejection factor. After the ANN, the challenging correlated cosmogenic BG dominates the total BG rate, as expected due to the shallow overburden. The FD is better shielded. Extra rejection uses the cosmogenic vetoes. The overall rejection factors are $\sim 193$ (FD) and $\sim 34$ (ND) relative to the loose coincidence. The ANN leads to relatively small rejection fraction of fast-neutrons via the spatial coincidence condition since those events could extend over a somewhat larger volume.

The ANN optimisation criteria discussion is illustrated in Fig. 8. Here, the challenge is matching data to MC for both the $\Delta t_{\text{unicity}}$ variable and energy resolution, both considered by DC in previous publications. The BG relies on high statistics samples of accidental BG obtained using the standard “off-time window” technique, as described in [13]. So, the overall training of the ANN makes the rejection insensitive to events with a neutron in the final state, which includes IBD’s, all cosmogenic BG and some remaining accidental BG. This is consistent with the observation, again in Fig. 2 that the remaining accidental BG (dark grey) exhibits a clear irreducible H-n peak contribution. An important intrinsic feature of the ANN is that it internally constructs the best combined selection from all the input variables to achieve an optimal signal identification. This yields a superior background rejection as compared to the combination of simple cosmogenic BGs, specially the $^9\text{Li BG}$.

6The ANN leads to relatively small rejection fraction of fast-neutrons via the spatial coincidence condition since those events could extend over a somewhat larger volume.
individual cuts, as done in the Gd-only selection – irrelevant then since the accidental BG rate was negligible. The overall accidental BG rejection factor is ≥400× using the ANN, as illustrated in Fig. 9. The input ANN selection is so widely open that the selection efficiency in the GdT is close to 100% in the detector center due to the lower average capture time caused by the presence of Gd. The ANN mainly affects events in the GC which are most sensitive to the rejection of external accidental BG. Much of the ANN performance was explored, tuned and demonstrated during the H-n capture θ13 measurement [14]. Once the ANN-based IBD coincidence is defined, the unicity condition imposes that only two-fold coincidences are considered (consistent with IBD’s). After the ANN, cosmogenic BG dominates, especially in the ND due to the lower overburden.

The last stage is the application of vetoes targeting the cosmogenic BGs. The vetoes exploit the fact that BG events deposit energy in different detector sensitive layers used to tag large fractions (typically about ~50%) of the cosmogenic BG. Further details on the definition of all vetoes is found in [13, 14] where they were first used and fully described. The performance of the IBD TnC selection is illustrated in Fig. 9 including the overall BG rejection factors and selection efficiency 0.0164 and χ²/DoF: 28/37). Overall consistency is found within uncertainties.

Further θ13 Measurement Cross-Checks

As part of the internal validation of the nominal θ13 measurement, we have investigated several fits. Fig. 10 summarises the most relevant results. Most of them have already been mentioned in past sections when demonstrating the minimal dependence of the θ13 measurement to the observed spectral distortion between the data and the prediction model. However, we shall highlight two other θ13 measurements with lower dependence on the model.

First is the so called Data-to-Data (D2D) rate+shape θ13 fit. This fit determines θ13 by the comparison of the observed FD spectra to the prediction extracted from the ND spectrum. The model calculation is used only as the ratio of near versus far spectra, reducing the sensitivity to potential common model biases. The FD-I can be compared to the ND data with a correction using MC ratios, whose uncertainty is significantly smaller than the statistical precision. Likewise, common (or correlated) inter-detector effects are expected to fully cancel. The value obtained is sin²2θ13 = 0.103 ± 0.017 (sensitivity 0.0164 and χ²/DoF: 28/37). The output inter-detector ratio exhibits no traces of spectral mismatch and looks identical to Fig. 4 (left). The excellent agreement found to the main result (deviation δ ≤ 0.002 units) further supports the negligible impact of the reactor model in the MD case. Second is the so called Reactor Rate Modulation (RRM) rate-only θ13 fit that has been articulated in past publications for the SD only [49]. The rate-only implementation makes the RRM spectral distortion independent, by construction, even though the reactor rate prediction is indeed used. This time the RRM prompt energy window has been confined to the [1.0,8.5] MeV range with a slightly better IBD signal to BG. This way, the [8.5,12.0] MeV window data provide an independent cosmogenic BG counter (no shape information) used in the fit for a higher precision BG constraint. The RRM fit remains very sensitive to the BG input, unlike the nominal rate+shape implementation reassessing the BG using the spectral shapes. Despite this limitation, the novel RRM method yields a more precise BG with similar precision to the nominal θ13 measurement. The best value obtained is sin²2θ13 = 0.095 ± 0.016 (χ²/DoF = 12/14). Overall consistency is found within uncertainties.

The Reactor Model Structure

Since 2011, there has been much debate about the rate deficit emerging with the revision of the reactor flux prediction [20, 27], including possible new particle physics [50, 51]. Before, the predicted shape appeared to be well reproduced by the world best data sample provided by the Bugey3 ex-
with the Bugey4 constraint. The RN data to MC ratios (R = \text{normalisation}, etc.) the ND number of events a rate+shape analysis was addressed in Fig. 4.

experiments and factorises out any rate normalisation effect. The shape-only basis allows for some simplified comparison among and there could be rate-to-shape correlations. Regardless, the characterisation of those features while their specific origin and mechanism remains unknown. We exploit the shape-only basis even though this basis might not fully describe the data by ignoring the rate information. In fact, rate effects are large and there could be rate-to-shape correlations. Regardless, the shape-only basis allows for some simplified comparison among experiments and factorises out any rate normalisation effect. A rate+shape analysis was addressed in Fig. 4.

Upon the oscillation fit (including corrections for \(\theta_{13}\), BG, normalisation, etc.) the ND number of events data to MC ratios (R\(_N\)) are R\(_N\)(ND) = 0.943 ± 0.022 and R\(_N\)(ND\text{\(\oplus\)}Bugey4) = 0.995 ± 0.019, respectively, without and with the Bugey4 constraint. The R\(_N\)(ND) exhibits a 2.6\(\sigma\) rate deficit, where the uncertainty is largely dominated by the prediction normalisation of the reactor model, as illustrated in Fig. 4. The quoted significance could be lower in case the reactor model uncertainty would be underestimated. The R(Bugey4) illustrates the excellent agreement found between Bugey4 and DC data. Most experiments are consistent with these observations.

The significant observation of a spectral distortion by the reactor-\(\theta_{13}\) experiments, dominated by an excess of events around 5 MeV, has shed new light on the prediction capability of the reactor antineutrino spectrum. Its existence suggests the presence of sizeable shape inaccuracies in addition to the aforementioned rate-only deficit. The 5 MeV excess was first reported \cite{52} and published by DC in 2014 \cite{13}. Confirmations by RENO and DYB were reported shortly after. Today, there is a significant DC ND data to model disagreement when considering the full energy range. The ND and FD spectra, once corrected for \(\theta_{13}\), match within \(\sim 1\sigma\text{(stat)}\) averaged over the full spectra. This is expected as effective cancellation across detectors has been demonstrated in Fig. 4 as a part of the \(\theta_{13}\) measurement consistency. In addition, this distortion scales with reactor power: 6.4\(\sigma\) (ND) and 7.1\(\sigma\) (FD). So, an unknown new BG hypothesis is ruled out. The reactor model prediction and/or a residual non-linearity in the detector energy response \cite{12} are both possible hypotheses. More exotic hypotheses have also been suggested \cite{51}.

Further insight on the origin of the spectral distortion is gained when comparing the data to prediction ratios provided by the DC, DYB, RENO and NEOS \cite{55} experiments, as summarised in Fig. 11. From the point of view of shape distortions, the leading order pattern is well reproduced by all latest experiments. Other past experiments are not conclusive but CHOOZ and Goesgen \cite{56} data show a similar pattern. However, the Bugey3 experiment (not shown) exhibits a rather flat and featureless ratio spectrum, thus inconsistent with the

Figure 11: Shape-Only Reactor Spectral Distortion. The data to prediction spectral ratio for the latest DC-ND (black), DYB \cite{38} (blue), RENO \cite{11} (red), NEOS \cite{55} (green) are shown, exhibiting a common dominant pattern predominantly characterised by the 5 MeV excess. Small differences across experiments are still possible but unresolved so far. The Bugey3 \cite{51} (not shown) is the only experiment known not to reproduce this structure. This remains an issue. The RENO and NEOS normalisation has been modified relative to publications to ensure the shape-only condition (average R = 1) is met. The reactor model prediction shape-only uncertainty is shown in grey, which is significantly smaller than the dominant rate-only uncertainties. Since the same reactor model prediction is used, this uncertainty is expected to remain a representative guideline to all experiments. The 5 MeV excess is compensated by a deficit region [1.5,4.0] MeV for all experiments due to the shape-only condition. A good agreement is found between DC and Daya Baya data throughout the entire energy range. The non-trivial match among different experiments suggests that most detector and part of the reactor effects are accurately reproduced by the MC, thus cancelling across in R. This implies that the common reactor prediction model inaccuracies are expected to dominate the observed distortion. This is consistent with the fact that all other experiments use the same prediction strategy.
shown experiments. The data to prediction ratio is expected to cancel (or suppress) the following contributions: a) the known and calibrated energy non-linearities, b) the detection effects (selection, vetoes, background subtraction bias), c) the $\theta_{13}$ spectral distortion, d) the reactor fuel burn-up and e) the overall shape of the reactor prediction, common to all experiments. However, energy response could be more complex since different detectors might have different capabilities to resolve spectral features. Nevertheless, most detector specific systematics are expected to be suppressed as a consequence of the MC tuning campaign to match the calibration data. So, while detector effects cannot be a priori ruled out, they are expected to be sub-dominant and not a priori identical across different experimental setups. Instead, the external reactor model, common to all experiments here considered, cannot be tuned to yield a data to prediction match.

A good agreement has been found between DC ND and the DYB ND’s. Their similar energy resolutions allow a simpler direct comparison. While all experiments plotted in Fig. 11 show similar behaviour, the RENO data seem to hint at a larger 5 MeV excess. NEOS, located in the same RENO reactor plant, also appears consistent with DC data. By re-binning the DC data to the DYB binning, the compatibility of both data sets can be studied. The DC data has been corrected to account for the non-linearity constrained during the $\theta_{13}$ fit and $\theta_{13}$ itself. The DC uncertainties consider statistics and systematics on both BG and energy. DC and DYB data are found consistent within uncertainties over the entire IBD energy spectrum. The best agreement manifests below 6 MeV. This observation disfavours detector driven effects as the main cause of the spectral distortion. Hence, the dominant cause appears to be the prediction inaccuracies in the reactor model. More data should further resolve the issue.

Future efforts should scrutinise the data to prediction ratio in the rate+shape basis for completeness, as the shape-only basis remains incomplete. So far, most experimental data appear consistent with a reactor model origin for the observed spectral distortion as the leading order effect. Sub-dominant detector and/or reactor effects might still exist, possibly explaining some of the residual differences across experiments. So, today the undistorted Bugey3 spectrum remains an issue to be understood to be able to yield a coherent experimental vision. The need for further nuclear physics effort to improve the reactor model predictions remains a critical topic for high precision reactor neutrino physics. This is particularly important for any SD experiments. Last, a debate remains whether the main discrepancy feature is only the 5 MeV excess region since that part exhibits R-1, as illustrated in Fig. 11 in a rate+shape treatment. A preliminary more conservative reactor model error budget remains a pending issue with important impact to past and present experiments, so that their data and uncertainties are properly treated.

The Mean Cross-Section per Fission

The mean cross-section per fission, $\langle \sigma_f \rangle$, is defined as:

$$\langle \sigma_f \rangle = \frac{N(\bar{\nu}_e)}{N_p \times \epsilon} \times \left( \sum_{r=B1,B2} \frac{(P_{th})_r}{4\pi L_r^2 \times \langle E_f \rangle_r} \right)^{-1} [\text{cm}^2/\text{fission}]$$

where $\langle \sigma_f \rangle$ provides a reactor neutrino interaction probability measured over the integrated reactor spectrum. $N(\bar{\nu}_e)$ stands for the measured neutrino rate (after BG subtraction and correction for the $\theta_{13}$ driven oscillation) in the considered detector, $N_p$ is the proton-number in the target volume, $\epsilon$ is the average absolute detection efficiency, $\langle P_{th} \rangle_r$ is the average thermal power per reactor, $L_r$ is the baseline per reactor and $\langle E_f \rangle_r$ is the average energy per fission per reactor [57]. The main physical value and meaning of the $\langle \sigma_f \rangle$ is that it provides a measure of the total reactor neutrino integrated flux measurement per reactor which can be used as a common reference across experiments. By correcting $\langle \sigma_f \rangle$ for the IBD cross-section a flux could be obtained (i.e. neutrino/fission).

DC provides the most precise $\langle \sigma_f \rangle$ thus superseding the Bugey4 measurement, whose precision is 1.4%. Both measurements are compatible with one another as illustrated in Fig. 7. The systematics breakdown of the ND is summarised in Table 6 yielding just below 1% precision for the first time. Other experiments might use the DC $\langle \sigma_f \rangle$ in the same way DC has been using the $\langle \sigma_f \rangle$ of Bugey4; i.e. to factorise out the rate (or normalisation) bias of the reactor prediction model. As demonstrated in the $\theta_{13}$ measurement discussion, $\langle \sigma_f \rangle$ continues to be a critical value for past and future SD reactor experiments. The $\langle \sigma_f \rangle$ measured by experiments using commercial nuclear reactors can slightly differ because of the different fuel composition. Differences typically at the 0.5% level are expected when several reactor cycles are considered. For comparison with other experiments, two pieces of information are necessary. The ND average fission fractions 52.0% ($^{235}U$), 8.7% ($^{238}U$), 33.3% ($^{239}Pu$) and 6.0% ($^{241}Pu$). And, also the $\langle \sigma_f \rangle$ per isotope 6.69 ($^{235}U$), 10.10 ($^{238}U$), 4.36 ($^{239}Pu$) and 6.05 ($^{241}Pu$) in units of $10^{-43} \text{cm}^2$ per fission, as obtained from [58].

<table>
<thead>
<tr>
<th>Uncertainty (%)</th>
<th>ND</th>
</tr>
</thead>
<tbody>
<tr>
<td>Proton Number</td>
<td>0.66</td>
</tr>
<tr>
<td>Thermal Power</td>
<td>0.47</td>
</tr>
<tr>
<td>ThC Selection</td>
<td>0.24</td>
</tr>
<tr>
<td>Background</td>
<td>0.18</td>
</tr>
<tr>
<td>Energy per Fission</td>
<td>0.16</td>
</tr>
<tr>
<td>$\theta_{13}$ Correction</td>
<td>0.16</td>
</tr>
<tr>
<td>Statistics</td>
<td>0.22</td>
</tr>
<tr>
<td>Total</td>
<td>0.97</td>
</tr>
</tbody>
</table>

The FD $\langle \sigma_f \rangle$ is, by definition, consistent with the ND upon $\theta_{13}$ correction. The FD alone yields a precision close 1.1%. Further precision improvements are expected in the future for $\langle \sigma_f \rangle$, especially if the dominant proton precision is improved upon detector dismantling.
References

Bauer C. et al. JINST 6 P06008 (2011)
[44] Beissel F. et al. JINST 8 01 T01003 (2013)
[47] Abe Y. et al. JINST 8 08 P08015 (2013)
    [indico.lal.in2p3.fr/event/2454].
[54] Berryman J. et al. [arXiv:1803.08506
[56] Zacek V. et al. [arXiv:1807.01810